# A Note on KAM Theorem for Symplectic Mappings 

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# A Note on KAM Theorem for Symplectic Mappings 

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## 1. Introduction and main result

a. In this note we prove the existence of differentiable foliation structures of invariant tori for nearly integrable symplectic mappings in the sense of Whitney as J. Pöschel did in the case of Hamiltonian systems([10]). The formulation of the main theorem is almost a copy of that given in [10] and the proof also follows from essentially the same arguments only some different technical details should be pointed out particularly. The main aim of the note is to give two relevant estimates: one is for perturbations under which a majority of invariant tori of the integrable mapping persists and another is for invariant Liouville measure of the set filled by the invariant tori of the perturbed mappings in phase space. The estimates are both given explicitly in terms of three parameters: the first one is $\gamma$, which appears in the well-known diophantine condition satisfied by the frequencies, say $\omega=\left(\omega_{1}, \cdots, \omega_{n}\right)$, of invariant tori

$$
\begin{equation*}
\left|e^{i(k, \omega\rangle}-1\right| \geq \frac{\gamma}{|k|^{\tau}}, \quad \text { for } 0 \neq k=\left(k_{1}, \cdots, k_{n}\right) \in \mathbf{Z}^{n} \tag{1.1}
\end{equation*}
$$

with another fixed constant $\tau>0$, where $\langle k, \omega\rangle=\sum_{j=1}^{n} k_{j} \omega^{(j)}$ and $|k|=\sum_{j=1}^{n}\left|k_{j}\right|$ for integers $k \in \mathbf{Z}^{n}$. The other two, say $\theta$ and $\Theta$, are used to describe the nondegeneracy of the frequency map $\omega: I \rightarrow \Omega$, of the unperturbed integrable mapping, and its inverse, for example, of the form

$$
\begin{equation*}
\theta\left|p_{1}-p_{2}\right| \leq\left|\omega\left(p_{1}\right)-\omega\left(p_{2}\right)\right| \leq \Theta\left|p_{1}-p_{2}\right| \tag{1.2}
\end{equation*}
$$

Here $I$ and $\Omega$ are the domains of action variables and the corresponding frequency values respectively. We always assume that they are open in $\mathbf{R}^{n}$. In fact, in our case we require $\omega: I \rightarrow \Omega$ to be analytic so $\omega$ is assumed to be defined and analytic in some complex extension, say $I+r$, with radius $r$ of real domain $I \subset \mathbf{R}^{n}$ of action variables (cf. [10]) and (1.2) is assumed to be true for $p_{1}, p_{2} \in I+r$ with $\left|p_{1}-p_{2}\right| \leq r$. Note that this nondegeneracy condition implies that the frequency map $\omega$ is invertible in any ball with radius $r$ and centered in $I$ and so it is stronger
than the standard Kolmogorov's nondegeneracy assumption but equivalent to the Pöschel's assumption ([10]).

It should be noted that there are already some optimal results on the estimates in terms of $\gamma$. For example, H. Rüssmann([11], $n=2$ ) and N. V. Svanidze([14]) in mapping case with $\omega(p)=p$ (and so $\theta=\Theta=1$ ) and J. Pöschel( $([10])$ and A. I. Neishtadt $([8])$ in Hamiltonian system case proved that if a perturbation has a norm, say $\varepsilon$, in some relevant function space, proportional to $\gamma^{2}$ with coefficient, say $\Delta$, not dependent of $\gamma$ and small enough, then the perturbed mapping has invariant tori, deformed slightly from those of unperturbed one, whose complement in the phase space has invariant Liouville measure proportional to $\gamma$ with coefficient, say $C$, not dependent of $\gamma$. Therefore, the invariant tori of the perturbed mapping form a set with large measure in the phase space if $\gamma$ is sufficiently small (in fact, $\gamma$ may be chosen reasonably to be proportional to $\sqrt{\varepsilon}$ from the fact that $\varepsilon$ is admitted to be proportional to $\gamma^{2}$ (cf. [10], [8])). In this note we also obtain the same results.

We note that in most general cases, the bounds of both $\Delta$ and $C$ in various versions of KAM theorem depends on the "nondegenerate" behavior, in some sense, of the frequency map $\omega: I \rightarrow \Omega$ (cf. $[1,2,6,7]$ and $[3,12]$ ). Here we work simply under the "ordinary" nondegeneracy assumption (1.2) (this is the strongest nondegeneracy assumption) and investigate the dependence of the bounds of $\Delta$ and $C$, say $\Delta_{0}$ and $C_{0}$ respectively, on the nondegeneracy (or twist) parameters $\theta$ and $\Theta$. We obtain the dependence to be $\Delta_{0}=\delta_{0} \theta \Theta^{-2}$ and $C_{0}=c_{0}\left(\theta \Theta^{-1}\right)^{-n}$, where $\delta_{0}$ and $c_{0}$ are constants not dependent of $\theta, \Theta$ or $\gamma$. If $\omega$ only satisfies the Kolmogorov's nondegeneracy condition, say,

$$
\begin{equation*}
\theta|d p| \leq|d \omega(p)| \leq \Theta|d p|, \quad \text { for } p \in I+r \tag{1.2}
\end{equation*}
$$

with constants $0<\theta \leq \Theta$, then we may get the estimate of the form $\Delta_{0}=\delta_{0} \theta^{2} \Theta^{-3}$ for perturbation part. We remark that this estimate is better than that given in [5], in the case of analytic perturbations, by directly estimating the convergence of the Lindstedt series for individual KAM tori, which is essentially of the form $\varepsilon=\delta_{0} \gamma^{4} \theta^{2} \Theta^{-4}$ - this estimate can not be applied to small twist problem.

We are interested in such estimates because they relate to some significant problems such as the small twist problem and therefore also relate to the stability of symplectic integrators in solving integrable and nearly integrable Hamiltonian
systems. In the small twist problem, a small parameter, say $s$, enters into the frequency map and so the parameters $\theta, \Theta$ and even $\gamma$ are all $s$-dependent, say, they turn into $s \theta, s \Theta$ and $s \gamma$ respectively. For the case of $n=1$, Moser obtained the existence of invariant curves ([7]). But for $n>1$, I have not seen an available reference. Our estimates naturally lead to a proper answer to the problem.
b. We consider an exact symplectic mapping $S:(p, q) \rightarrow(\hat{p}, \hat{q})$ to be defined in phase space $I \times \mathbf{T}^{n}$ by its generating function $H: I \times \mathrm{T}^{n} \rightarrow \mathbf{R}$ through relation

$$
\left\{\begin{array}{l}
\hat{p}=p-\frac{\partial H}{\partial q}(\hat{p}, q)  \tag{1.3}\\
\hat{q}=q+\frac{\partial H}{\partial p}(\hat{p}, q),
\end{array}\right.
$$

where $I$ is an open and usually bounded set of $\mathbf{R}^{n}$ in which the action variables $p$ vary and $\mathbf{T}^{n}$ is the standard $n$-torus which is identified as $\mathbf{R}^{n} / 2 \pi \mathbf{Z}^{n}$ and the angle variables $q$ vary on. When $H(p, q)=H_{0}(p)$ does not depend on angle variables, then (1.3) represents an integrable mapping $S=S_{0}:(p, q) \rightarrow(\hat{p}, \hat{q})=$ $(p, q+\omega(p))$, which is well defined on $I \times \mathbf{T}^{n}$, with frequency map

$$
\begin{equation*}
\omega(p)=\frac{\partial H_{0}}{\partial p}(p), p \in I \tag{1.4}
\end{equation*}
$$

Under $S_{0}$, the phase space $I \times \mathbf{T}^{n}$ is completely foliated into invariant $n$-tori $\{p\} \times \mathbf{T}^{n}, p \in I$, on each of which the iterations of $S_{0}$ are linear with frequencies $\omega=\omega(p)$. When a perturbation, say $h(p, q)$, is added to $H_{0}$, i.e., when $H(p, q)=H_{0}(p)+h(p, q)$, then (1.3) does not define an integrable mapping generally. However, KAM theorem shows that the perturbed mapping $S$, defined by the perturbed generating function $H$ according to (1.3), still exibits in a large extent the integrable behavior even not in any open domain but at least in some Cantor set of the phase space if the frequency map $\omega$ is nondegenerate in some sense (see $[1,2,6,7]$ for strong (ordinary) nondegeneracy and [3,12] for weak nondegeneracy) and if the perturbation $h$ belongs to the class $C^{a}\left(I \times \mathrm{T}^{n}\right)$ for a suitably large number $a>0$ with a sufficiently small associated norm (cf. [10]). We note that for $h$ small in $C^{2}\left(I \times \mathbf{T}^{n}\right)$, (1.3) really defines an exact symplectic mapping on $I^{\prime} \times \mathrm{T}^{n}$ with some open subset $I^{\prime}$, of $I$, whose boundary is a little away from the boundary of $I$ with distance equal to the $C^{2}$-norm of $h$.

As a Banach space, the class $C^{a}\left(I \times \mathbf{T}^{n}\right)$, with the corresponding norm denoted as $\|\cdot\|_{a, I \times \mathbf{T}^{n}}$, is always understood according to [10]. In this note, we also get the anisotropic differentiability of the foliations of invariant tori. So we also need the
class $C^{\nu_{1}, \nu_{2}}\left(I \times \mathrm{T}^{n}\right)$ of anisotropic differentiable functions, whose exact definition was also given in [10]. The norm of a function in $C^{\nu_{1} \nu_{2}}\left(I \times \mathrm{T}^{n}\right)$ is denoted by $\|\cdot\|_{\nu_{1}, \nu_{2} ; I \times \mathbf{T}^{n}}$. We also use another norm $\|\cdot\|_{\nu_{1}, \nu_{2} ; I \times \mathbf{T}^{n}, \rho}$ for $\rho>0$ defined by

$$
\begin{equation*}
\|u\|_{\nu_{1}, \nu_{2} ; I \times \mathbf{T}^{n}, \rho}=\left\|u \circ \sigma_{\rho}\right\|_{\nu_{1}, \nu_{2} ; \sigma_{\rho}^{-1}\left(I \times \mathbf{T}^{n}\right)} \tag{1.5}
\end{equation*}
$$

for $u \in C^{\nu_{1}, \nu_{2}}\left(I \times \mathrm{T}^{n}\right)$, where $\sigma_{\rho}$ denotes the partial stretching $(x, y) \rightarrow(\rho x, y)$ for $(x, y) \in I \times \mathbf{T}^{n}$. Note that the following relation between these two norms is valid for $0<\rho \leq 1$ :

$$
\begin{equation*}
\|u\|_{\nu_{1}, \nu_{2} ; \rho} \leq\|u\|_{\nu_{1}, \nu_{2}} \leq \rho^{-\nu_{1}}\|u\|_{\nu_{1}, \nu_{2} ; \rho} \tag{1.6}
\end{equation*}
$$

where we dropped the domains to simplify the notations.
Write $\Omega=\omega(I)$ and denote by $\Omega_{\gamma}$ the set of those frequencies, in $\Omega$, which satisfy the diophantine condition (1.1) for the given $\gamma$ and whose distance to the boundary of $\Omega$ is at least equal to $2 \gamma$. The difference $\Omega \backslash \bigcup_{\gamma>0} \Omega_{\gamma}$ is a zero set if $\tau>n+1$ and so $\Omega_{\gamma}$ is large for small $\gamma$.

Now we formulate our main result as follows.
Theorem 1.1. With given positive integer $n$ and given $\tau>n+1$, consider mapping $S$ to be defined in phase space $I \times \mathrm{T}^{n}$ by generating function $H(p, q)=$ $H_{0}(p)+h(p, q)$ in the form of (1.3), where $H_{0}(p)$ is analytic in $p \in I+r$ with $r>0$ and $h(p, q)$ belongs to the class $C^{\alpha \lambda+\lambda+\tau}\left(I \times \mathrm{T}^{n}\right)$ with fixed $\lambda>\tau+1$ and $\alpha>1$ not in the discrete set

$$
\Lambda=\{i / \lambda+j: i, j \geq 0 \text { integer }\}
$$

Suppose that the frequency map $\omega: I \rightarrow \Omega$, defined by (1.4) from $H_{0}$, satisfies the nondegeneracy condition (1.2) for $p_{1}, p_{2} \in I+r$ with $\left|p_{1}-p_{2}\right| \leq r$ where the constants $\theta$ and $\Theta$ satisfy $0<\theta \leq \Theta$. Then there exists a positive constant $\delta_{0}$, depending only on $n, \tau, \lambda$ and $\alpha$, such that for any $0<\gamma \leq \min \left(1, \frac{1}{2} r \Theta\right)$, if

$$
\begin{equation*}
\|h\|_{\alpha \lambda+\lambda+\tau, I \times \mathbf{T}^{n_{;}} \gamma \Theta^{-1}} \leq \delta_{0} \gamma^{2} \theta \Theta^{-2} \tag{1.7}
\end{equation*}
$$

then there exist a closed set $I_{\gamma} \subset I$, a surjective map $\omega_{\gamma}: I_{\gamma} \rightarrow \Omega_{\gamma}$ of $C^{\alpha+1}$ class and an injection $\Phi: I_{\gamma} \times \mathbf{T}^{n} \rightarrow \mathbf{R}^{n} \times \mathbf{T}^{n}$ of $C^{\alpha, \alpha \lambda}$ class, in the Whitney's sense, such that
(1) $S \circ \Phi=\Phi \circ R$. Moreover, this equation may be differentiated as often as $\Phi$ allows, where $R$ is the integrable rotation on $I_{\gamma} \times \mathrm{T}^{\mathrm{n}}$ with frequency map $\omega_{\gamma}$.
(2) If $\Omega$ is a bounded open set of type $D$ in the Arnold's sense [1], then we have measure estimate

$$
\begin{equation*}
m \mathcal{E}_{\gamma} \geq\left(1-c_{4}\left(\theta \Theta^{-1}\right)^{-n} \gamma\right) m \mathcal{E} \tag{1.9}
\end{equation*}
$$

where $m$ denotes the invariant Liouville measure of the phase space $\mathcal{E}=I \times \mathbf{T}^{n}$ and $\mathcal{E}_{\gamma}=\Phi\left(I_{\gamma} \times \mathbf{T}^{n}\right) ; c_{4}$ is a positive constant depending on $n, \tau, a$ and the geometry of the domain $\Omega$.
(3) If $h$ is of $C^{\beta \lambda+\lambda+\tau}$ class with $\alpha \leq \beta$ not in $\Lambda$, then we have further that $\omega_{\gamma} \in C^{\beta+1}\left(I_{\gamma}\right)$ and $\Phi \in C^{\beta, \beta \lambda}\left(I_{\gamma} \times \mathrm{T}^{n}\right)$. Moreover,

$$
\begin{equation*}
\left\|\sigma_{\gamma \Theta^{-1}}^{-1} \circ(\Phi-I)\right\|_{\beta, \beta \lambda ; \gamma \Theta^{-1}}, \gamma^{-1}\left\|\omega_{\gamma}-\omega\right\|_{\beta+1 ; \gamma \Theta^{-1}} \leq c_{5} \gamma^{-2} \Theta\|h\|_{\beta \lambda+\lambda+\tau ; \gamma \Theta^{-1}} \tag{1.10}
\end{equation*}
$$

with constant $c_{5}$ depending on $n, \tau, \lambda$ and $\beta$, here we have droped the domains in the notation of norms.

Theorem 1.2. In Theorem 1.1, if the frequency map $\omega$ satisfies the nondegeneracy condition (1.2)', then the required smallness condition for $h$ is, instead of (1.7),

$$
\begin{equation*}
\|h\|_{\alpha \lambda+\lambda+\tau, I \times \mathbf{T}^{n} ; \gamma \Theta^{-1}} \leq \delta_{0} \gamma^{2} \theta^{2} \Theta^{-3} \tag{1.11}
\end{equation*}
$$

with sufficiently small $\delta_{0}$, depending only on $n, \tau, \lambda$ and $\alpha$, under which the conclusions of Theorem 1.1 are also true with the same estimates.

Remark: We have following two further conclusions in the above theorems:

1. If $h \in C^{\infty}\left(I \times \mathrm{T}^{n}\right)$, then $\omega_{\gamma} \in C^{\infty}\left(I_{\gamma}\right)$ and $\Phi \in C^{\infty}\left(I_{\gamma} \times \mathrm{T}^{n}\right)$ with the estimates (1.10) for any $\beta \geq \alpha$.
2. If $h \in C^{\omega}\left(I \times \mathrm{T}^{n}\right)$, then we have further $\Phi \in C^{\infty, \omega}\left(I_{\gamma}, \mathrm{T}^{n}\right)$ under a further smallness condition for $\delta$ which also depends on the radius of analyticity of $h$ with respect to angle variables.

## 2. Outline of the proof of Theorem 1.1

In this section, we outline the proof of Theorem 1.1. The detailed arguments will be given in the latter sections.
a. First we transform the mapping $S$ by the partial coordinates stretching $\sigma_{\rho}:(x, y) \rightarrow(p, q)=(\rho x, y)$ of phase space $I \times \mathrm{T}^{n}$ and obtain a new mapping
$T=\sigma_{\rho}^{-1} \circ S \circ \sigma_{\rho}:(x, y) \rightarrow(\hat{x}, \hat{y})$ to be defined in the new phase space $I_{\rho} \times \mathbf{T}^{n}$ by

$$
\left\{\begin{array}{l}
\hat{x}=x-\frac{\partial F}{\partial y}(\hat{x}, y)  \tag{2.1}\\
\hat{y}=y+\frac{\partial F}{\partial x}(\hat{x}, y)
\end{array}\right.
$$

where

$$
\begin{equation*}
F(x, y)=F_{0}(x)+f(x, y) \tag{2.2}
\end{equation*}
$$

is well defined on $I_{\rho} \times \mathrm{T}^{n}$ with

$$
\begin{equation*}
F_{0}(x)=\rho^{-1} H_{0}(\rho x), \quad f(x, y)=\rho^{-1} h(\rho x, y) \tag{2.3}
\end{equation*}
$$

and

$$
\begin{equation*}
I_{\rho}=\rho^{-1} I=\left\{x \in \mathbf{R}^{n} \mid \rho x \in I\right\} \tag{2.4}
\end{equation*}
$$

For the time being, $\rho$ is considered as a free parameter. $F_{0}(x)$ is real analytic for $x \in I_{\rho}+r_{\rho}$ with $r_{\rho}=\rho^{-1} r$ and $f$ belongs to the class $C^{a}\left(I_{\rho} \times \mathrm{T}^{n}\right)$ with $a=\alpha \lambda+\lambda+$ $\tau$. So the new mapping $T$ satisfies the same assumptions of Theorem 1.1 in which only $I, r, H, H_{0}, h$ are replaced by $I_{\rho}, r_{\rho}, F, F_{0}, f$ respectively. Correspondingly, the frequency map of integrable mapping associated to the generating function $F_{0}$ turns into

$$
\tilde{\omega}(x)=\frac{\partial F_{0}}{\partial x}(x), \quad x \in I_{\rho}
$$

and the nondegenerate condition for the mapping turns out to be

$$
\begin{equation*}
\rho \theta\left|x_{1}-x_{2}\right| \leq\left|\widetilde{\omega}\left(x_{1}\right)-\widetilde{\omega}\left(x_{2}\right)\right| \leq \rho \Theta\left|x_{1}-x_{2}\right|, \quad x_{1}, x_{2} \in I_{\rho}+r_{\rho},\left|x_{1}-x_{2}\right| \leq r_{\rho} . \tag{2.5}
\end{equation*}
$$

In addition, from (2.3), we have

$$
\|f\|_{a, I_{\rho} \times \mathbf{T}^{n}}=\rho^{-1}\|h\|_{a, I \times \mathbf{T}^{n} ; \rho} .
$$

From now on we fix $\rho=\gamma \Theta^{-1}$. Then the assumption $0<\gamma \leq \frac{1}{2} r \Theta$ in Theorem 1.1 implies that $0<\rho \leq \frac{1}{2} r$ and so $r_{\rho} \geq 2$. Let $I_{\rho}^{*}$ be the set of points in $I_{\rho}$ with distance at least one to its boundary and let

$$
\begin{equation*}
I_{\rho ; \gamma}=\widetilde{\omega}^{-1}\left(\Omega_{\gamma}\right) \cap I_{\rho} . \tag{2.6}
\end{equation*}
$$

Then, from (2.5) and the definition of $\Omega_{\gamma}$, we have

$$
\begin{equation*}
\left(I_{\rho ; \gamma}+1\right) \cap \mathbf{R}^{n} \subset I_{\rho}^{*} \subset\left(I_{\rho}^{*}+1\right) \cap \mathbf{R}^{n} \subset I_{\rho}, \tag{2.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\gamma \mu\left|x_{1}-x_{2}\right| \leq\left|\tilde{\omega}\left(x_{1}\right)-\tilde{\omega}\left(x_{2}\right)\right| \leq \gamma\left|x_{1}-x_{2}\right|, \quad x_{1}, x_{2} \in I_{\rho}+2,\left|x_{1}-x_{2}\right| \leq 2 \tag{2.8}
\end{equation*}
$$

with $\mu=\theta \Theta^{-1}$.
b. We approximate $f$ by real analytic functions. Let

$$
\begin{equation*}
s_{j}=s_{0} 4^{-j}, \quad r_{j}=s_{j}^{\lambda}, \quad j=0,1,2, \cdots \tag{2.9}
\end{equation*}
$$

with fixed $\lambda>\tau+1$ and $s_{0}>0$ to be determined later on and let

$$
\mathcal{U}_{j}=I_{\rho} \times \mathbf{T}^{n}+\left(4 s_{j}, 4 s_{j}\right)
$$

Then by Approximation Lemma of [10], p. 665, there exist real analytic functions $f_{j}$ defined on $\mathcal{U}_{0}$ with $f_{0}=0$ such that, for $f \in C^{b}\left(I \times \mathbf{T}^{n}\right)$ with $b \geq a$, we have

$$
\begin{align*}
& \left|f_{j}-f_{j-1}\right|_{u_{j}} \leq s_{j}^{b} c_{b}\|f\|_{b ; I_{\rho} \times \mathrm{T}^{n}} \quad j=1,2, \cdots \\
& \left\|f-f_{j}\right\|_{\psi, I_{\rho} \times \mathrm{T}^{n}} \rightarrow 0 \quad(j \rightarrow \infty) \quad \text { for } 0<b^{\prime}<b \tag{2.10}
\end{align*}
$$

where $c_{b}$ is a positive constant only depending on $b, n$ and $s_{0}$ but not depending on the domain $I_{\rho}$ and hence not depending on the parameter $\rho$. Moreover, we may require $f_{j}$ to be $2 \pi$-periodic in the last $n$ variables. In (2.10), $|\cdot| u_{j}$ denotes the maximum norm of analytic functions on the complex domains $\mathcal{U}_{j}$. Almost all the notations in this note come from [10], which we will not specify particularly.
c. We give the KAM iteration process which is essentially the same as that given by Pöschel in Hamiltonian system case ([10]). Associating to each $f_{j}$, we will define a mapping $T_{j}:(x, y) \rightarrow(\hat{x}, \hat{y})$ by

$$
\left\{\begin{array}{l}
\hat{x}=x-\frac{\partial F_{j}}{\partial y}(\hat{x}, y)  \tag{2.11}\\
\hat{y}=y+\frac{\partial F_{j}}{\partial x}(\hat{x}, y)
\end{array}\right.
$$

with $F_{j}(x, y)=F_{0}(x)+f_{j}(x, y)$, which is well-defined and real analytic on $\mathcal{U}_{j}$ if $4 s_{j} \leq r_{\rho}=\rho^{-1} r$ which is always true for $j=0,1, \cdots$ if we choose $0<s_{0} \leq$ $4^{-1}$ (noting that $0<\gamma<\frac{1}{2} r \Theta$ ). We will show that each of $T_{j}$ for $j \geq 0$ is really well-defined on a complex domain, of the phase space $I_{\rho} \times \mathrm{T}^{n}$, which is appropriate for the KAM iteration if $h$ is bounded by (1.7) with $\delta_{0}$ chosen to be small enough, independently of $\gamma, \theta$ and $\Theta$. In fact, as $j$ approaches infinity, $T_{j}$
will converge to the $T$ in $C^{a^{\prime}}$ norm with $0<a^{\prime}<a-1$ on some sub-domain $I_{\beta}^{\prime} \times \mathrm{T}^{n}$ of the phase space $I_{\rho} \times \mathrm{T}^{n}$ with $I_{\beta}^{\prime}$ being a sufficiently large open subset of $I_{\rho}$. The central problem is to find transformations $\Phi_{j}$ and integrable rotations $R_{j}$ with frequency maps, say $\omega^{(j)}$, such that, as $j \rightarrow \infty$,

$$
\begin{equation*}
C_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ T_{j} \circ \Phi_{j} \rightarrow \text { identity }, \Phi_{j} \rightarrow \tilde{\Phi}, R_{j} \rightarrow \tilde{R}, \omega^{(j)} \rightarrow \tilde{\omega}_{\gamma} \tag{2.12}
\end{equation*}
$$

in $C^{\tilde{a}}$ norm, with a suitable positive number $\tilde{a}$, on $\tilde{I}_{p ; \gamma} \times \mathrm{T}^{n}$ with some closed set $\tilde{I}_{\rho ; \gamma} \subset I_{\rho}$ such that $\tilde{\omega}_{\gamma}\left(\tilde{I}_{\rho ; \gamma}\right)=\Omega_{\gamma}$; here $\tilde{\Phi}$ and $\tilde{R}$ are well-defined on $\tilde{I}_{\rho ; \gamma} \times \mathbf{T}^{n}$ and $\tilde{\omega}_{\gamma}$, as the frequency map of the rotation $\tilde{R}$, is well-defined on $\tilde{I}_{p ; \gamma}$. So in the limit we have

$$
\begin{equation*}
T \circ \tilde{\Phi}=\tilde{\Phi} \circ \tilde{R} \quad \text { on } \tilde{I}_{\rho ; \gamma} \times \mathbf{T}^{n} \tag{2.13}
\end{equation*}
$$

Reversing the mapping $T$ to $S$ by the partial coordinates stretching $\sigma_{\rho}$ and, at the same time, reversing $\widetilde{\Phi}$ to $\Phi$ and $\tilde{R}$ to $R$ with the frequency $\operatorname{map} \omega_{\gamma}(p)=\widetilde{\omega}_{\gamma}\left(\rho^{-1} p\right)$, then we have

$$
S \circ \Phi=\Phi \circ R \quad \text { on } I_{\gamma} \times \mathbf{T}^{n}
$$

with

$$
I_{\gamma}=\rho \widetilde{I}_{\rho ; \gamma}=\left\{p \in \mathbf{R}^{n} \mid \rho^{-1} p \in \tilde{I}_{p ; \gamma}\right\}
$$

a closed subset of $I$ and such that $\omega_{\gamma}\left(I_{\gamma}\right)=\Omega_{\gamma}$. This is just the conclusion (1) of Theorem 1.1.

The construction of $\Phi_{j}$ and $R_{j}$ uses the KAM iteration which we describe as follows.

Assume

$$
\begin{equation*}
\left|f_{j}-f_{j-1}\right| \sim \varepsilon_{j}, \quad j=1,2, \cdots \tag{2.14}
\end{equation*}
$$

with a decreasing sequence of positive numbers $\left\{\varepsilon_{j}\right\}_{1}^{\infty}$ and suppose we have already found a transformation $\Phi_{j}$ and a rotation $R_{j}$ with frequency map $\omega^{(j)}$ such that

$$
\begin{equation*}
C_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ T_{j} \circ \Phi_{j} \tag{2.15}
\end{equation*}
$$

satisfies

$$
\begin{equation*}
\left|C_{j}-I\right| \sim \varepsilon_{j+1} . \tag{2.16}
\end{equation*}
$$

We then construct a transformation $\Psi_{j}$ and a new rotation $R_{j+1}$ with frequency map $\omega^{(j+1)}$ such that

$$
\begin{equation*}
\Phi_{j+1}=\Phi_{j} \circ \Psi_{j} \tag{2.17}
\end{equation*}
$$

and (2.16) is also true for the next index $j+1$ with $C_{j+1}$ defined by (2.15) where $j$ is replaced by $j+1$. As was remarked in [10], "for this procedure to be successful it is essential to have precise control over the various domains of definition".

We define the transformation $\Psi_{j}:(\xi, \eta) \rightarrow(x, y)$ implicitly with the help of a generating function $\psi_{j}$ by

$$
\begin{equation*}
x=\xi+\frac{\partial \psi_{j}}{\partial y}(\xi, y), \quad y=\eta-\frac{\partial \psi_{j}}{\partial x}(\xi, y) \tag{2.18}
\end{equation*}
$$

To define the function $\psi_{j}$, we consider the mapping

$$
\begin{equation*}
B_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ T_{j+1} \circ \Phi_{j} \tag{2.19}
\end{equation*}
$$

which is also near identity and assumed to be given implicitly from its generating function $b_{j}$ by

$$
\begin{equation*}
\hat{x}=x-\frac{\partial b_{j}}{\partial y}(\hat{x}, y), \quad \hat{y}=y+\frac{\partial b_{j}}{\partial x}(\hat{x}, y) \tag{2.20}
\end{equation*}
$$

The function $\psi_{j}$ is then determined from $b_{j}$ by the following holomorphic equation

$$
\begin{equation*}
\psi_{j}\left(x, y+w^{(j)}(x)\right)-\psi_{j}(x, y)+\tilde{b}_{j}(x, y)=0 \tag{2.21}
\end{equation*}
$$

where $\widetilde{b}_{j}(x, y)=b_{j}(x, y)-\left[b_{j}\right](x)$ with $\left[b_{j}\right]$ being the mean value of $b_{j}$ over $\mathbf{T}^{n}$. Define

$$
\begin{equation*}
\omega^{(j+1)}(x)=\omega^{(j)}(x)+\frac{\partial\left[b_{j}\right]}{\partial x}(x) \tag{2.22}
\end{equation*}
$$

then $R_{j+1}:(x, y) \rightarrow(\hat{x}, \hat{y})$ is just given by

$$
\begin{equation*}
\hat{x}=x, \quad \hat{y}=y+\omega^{(j+1)}(x) \tag{2.23}
\end{equation*}
$$

With just defined $\Psi_{j}$ and $R_{j+1}$, we easily show that, formally,

$$
\Psi_{j}^{-1} \circ R_{j} \circ B_{j} \circ \Psi_{j}=R_{j+1} \circ C_{j+1}
$$

Similar formal calculation to that in [10] shows that (2.16) is valid if we replace $j$ by $j+1$.

We will not solve the equation (2.21) exactly indeed. Instead, we will solve an equation with a replacement of $\tilde{b}_{j}$ by a finite truncation of its fourier expansion with respect to angle variables so that "only finitely many resonances remain, and we obtain a real analytic solution $\psi_{j}$ defined on an open set" ([10], pp. 677). The idea goes back to Arnol'd [1,2].

In the following sections, we will make the above arguments more precise by carefully controlling the domains of definition of functions and mappings and giving the relevant estimates. For the convenience of the later statements, we follow [10] and abbreviate the differentiability orders by

$$
\begin{equation*}
a=\alpha \lambda+\lambda+\tau \leq b=\beta \lambda+\lambda+\tau, \tag{2.24}
\end{equation*}
$$

and set, for $f \in C^{b}\left(I_{\rho} \times \mathbf{T}^{n}\right)$,

$$
\begin{equation*}
\|f\|_{b, I_{\rho} \times \mathrm{T}^{n}}=\gamma \mu \delta_{b} . \tag{2.25}
\end{equation*}
$$

In addition, we denote by $\pi_{k}$ the translation $(x, y) \rightarrow\left(x, y+2 \pi e_{k}\right)$ of $\mathbf{R}^{n}$ for $k=1, \cdots, n$, where $e_{k}$ is the vector of $\mathbf{R}^{n}$ with the $k$-th entry equal to one and others equal to zero. With this notation, a well-defined transformation on $I \times \mathbf{T}^{n}$ may be viewed as a transformation on $I \times \mathbf{R}^{n}$ which commutes with $\pi_{k}$, $k=1, \cdots, n$.

## 3. Iterative lemma and proof of Theorem 1.1

Now we state a so-called iterative lemma which is a quantitative formulation of the KAM iteration process outlined in the above section.

Lemma 3.1. Assume that $F_{0}$ is real analytic on $I_{\rho}+2$ with the gradient map $\widetilde{\omega}=\frac{\partial F_{0}}{\partial x}$ satisfying the condition (2.8) in which $\gamma>0$ and $0<\mu \leq 1$ are given in advance. For fixed $\lambda>\tau+1$ and $\alpha>1$, let $f \in C^{a}\left(I_{\rho} \times \mathbf{T}^{n}\right)$ with $a=\alpha \lambda+\lambda+\tau$ and let $f_{j}$ be real analytic approximants to the $f$ as given in the preceeding section. Assume $\delta_{a} \leq \delta$. If $\delta$ is small enough, then for each $j \geq 0$, there exist a closed set $I_{\rho ; \gamma}^{(j)} \subset I_{\rho}$, a real analytic function $F_{0}^{(j)}$ on $V_{j}=I_{\rho ; \gamma}^{(j)} \times \mathrm{T}^{n}+\left(r_{j}, s_{j}\right)$ which is independent of the angle variables, and a real analytic exact symplectic map $\Phi_{j}: V_{j} \rightarrow \mathcal{U}_{j}$ which commutes with $\pi_{k}, k=1, \cdots, n$ such that, for $f \in C^{b}\left(I_{\rho} \times \mathbf{T}^{n}\right)$ with $b=\beta \lambda+\lambda+\tau \geq a$, the following hold.
(i) For $j \geq 1, I_{\rho ; \gamma}^{(j)}+r_{j} \subset I_{\rho ; \gamma}^{(j-1)}+\frac{5}{4} r_{j} \subset I_{\rho ; \gamma}^{(j-1)}+r_{j-1}$.
 condition: for $x_{1}, x_{2} \in I_{p ; \gamma}^{(j)}+r_{j}$ with $\left|x_{1}-x_{2}\right| \leq r_{j}$,

$$
\begin{equation*}
\left(1-\sum_{k=1}^{j} s_{k}\right) \gamma \mu\left|x_{1}-x_{2}\right| \leq\left|\omega^{(j)}\left(x_{1}\right)-\omega^{(j)}\left(x_{2}\right)\right| \leq\left(1+\sum_{k=1}^{j} s_{k}\right) \gamma\left|x_{1}-x_{2}\right| . \tag{3.1}
\end{equation*}
$$

(iii) $\Phi_{0}$ is the identity and for $j \geq 1, \Phi_{j}$ is well-defined and real analytic on

$$
V_{j-1}^{2}=I_{\rho ; \gamma}^{(j-1)} \times \mathrm{T}^{n}+\left(2 r_{j}, 2 s_{j}\right)
$$

and commutes with $\pi_{k}, k=1, \cdots, n$ and

$$
\begin{equation*}
\left|\Phi_{j}-\Phi_{j-1}\right|_{V_{j-1}^{2}} \leq r_{j}^{\beta} \cdot c_{1} \mu \delta_{b} \tag{3.2}
\end{equation*}
$$

(iv) $C_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ T_{j} \circ \Phi_{j}$ is well-defined and real analytic on $V_{j}$ and commutes with $\pi_{k}, k=1, \cdots, n$ and

$$
\begin{equation*}
\left|C_{j}-I\right|_{V_{j}} \leq s_{j+1}^{b-1} \cdot c_{2} \gamma \mu \delta_{b}, \quad \text { with } c_{2} \geq 16 c_{b} \tag{3.3}
\end{equation*}
$$

where $R_{j}$ is the rotation map $(x, y) \rightarrow\left(x, y+\omega^{(j)}(x)\right)$.
(v) $\omega^{(0)}=\tilde{\omega}$ and for $j \geq 1$,

$$
\begin{equation*}
\left|\omega^{(j)}-\omega^{(j-1)}\right|_{V_{j}} \leq r_{j}^{\beta+1} \cdot c_{3} \gamma \mu \delta_{b} . \tag{3.4}
\end{equation*}
$$

The proof of Lemma 3.1 is postponed to the next section. Now we prove Theorem 1.1, taking this lemma for granted. We let

$$
\tilde{I}_{\rho ; \gamma}=\bigcap_{j=0}^{\infty} \overline{I_{\rho ; \gamma}^{(j)}+r_{j}} .
$$

By (i) of Lemma 3.1, $\tilde{I}_{\rho ; \gamma}$ is a nonempty closed subset of $I_{\rho}$. Moreover, $\widetilde{I}_{\rho ; \gamma} \subset$ $I_{\rho ; \gamma}^{(j+1)}+r_{j+1}$ and hence $\tilde{I}_{p ; \gamma}+r_{j+1} \subset I_{\rho ; \gamma}^{(j)}+r_{j}$ for each $j$. Therefore, $\omega^{(j)}, \Phi_{j}$ and $C_{j}$ are well defined on $\tilde{V}_{j+1}=\tilde{I}_{\rho ; \gamma} \times \mathrm{T}^{n}+\left(r_{j+1}, s_{j+1}\right)$. From the estimates (3.2) - (3.4) and the Inverse Approximation Lemma of [10], P. 665, the sequences $\left(\omega^{(j)}-\widetilde{\omega}\right)$ and $\left(\Phi_{j}-I\right)$ have real limits

$$
\begin{equation*}
\tilde{\omega}_{\gamma}-\tilde{\omega} \in C^{\beta+1}\left(\tilde{I}_{\rho ; \gamma}\right), \quad\left\|\tilde{\omega}_{\gamma}-\tilde{\omega}\right\|_{\beta+1 ; \tilde{I}_{p ; \gamma}} \leq c_{5} \gamma \mu \delta_{b} \tag{3.5}
\end{equation*}
$$

and

$$
\begin{equation*}
\tilde{\Phi}-I \in C^{\beta, \beta \lambda}\left(\tilde{I}_{\rho ; \gamma} \times \mathbf{T}^{n}\right), \quad\|\widetilde{\Phi}-I\|_{\beta, \beta \lambda ; \tilde{I}_{p ; \gamma} \times \mathbf{T}^{n}} \leq c_{5} \mu \delta_{b} \tag{3.6}
\end{equation*}
$$

respectively, such that

$$
\begin{equation*}
\left\|\omega^{(j)}-\tilde{\omega}_{\gamma}\right\|_{\beta+1-\kappa, \tilde{I}_{\rho ; \gamma}} \rightarrow 0, \quad\left\|\Phi_{j}-\widetilde{\Phi}\right\|_{\beta-\iota,(\beta-\iota) \lambda ; \tilde{I}_{p ; \gamma} \times \mathbf{T}^{n}} \rightarrow 0 \tag{3.7}
\end{equation*}
$$

for any $\kappa>0$ and $\iota>0$, here $c_{5}$ is a constant depending only on $n, \tau, \lambda, \beta$. From (3.3), it follows that

$$
\begin{equation*}
\left|T_{j} \circ \Phi_{j}-\Phi_{j} \circ R_{j}\right|_{\tilde{v}_{j+1}} \leq s_{j+1}^{b-1} \cdot c_{6} \gamma \mu \delta_{b} \tag{3.8}
\end{equation*}
$$

since $R_{j}$ and $\Phi_{j}$ are uniformly bounded with their Jocobian if $\delta$ is small as requied - this fact follows from (2.8), (3.2) and (3.4) where we take $b=a$ and we will point it out again in next section. As argued in [10] and noting that $\left\|T_{j}-T\right\|_{b-1-\kappa ; I_{\rho} \times \mathbf{T}^{\mathbf{n}}} \rightarrow 0$ for any $\kappa>0$ with $\delta$ small enough, we see that (3.8) implies the following equation

$$
\begin{equation*}
T \circ \tilde{\Phi}=\tilde{\Phi} \circ \tilde{R} \tag{3.9}
\end{equation*}
$$

on $\tilde{I}_{\rho ; \gamma} \times \mathrm{T}^{n}$, which may be differentiated as often as $\tilde{\Phi}$ allows, where $\tilde{R}$, with $\tilde{\omega}_{\gamma}$ as its frequency map, is the limit of $R_{j}$. Transforming $T$ back to $S$ by the coordinates stretching $\sigma_{\rho}$ with $\rho=\gamma \Theta^{-1}$, and accordingly, transforming $\tilde{I}_{p ; \gamma}$ to $I_{\gamma}, \widetilde{\Phi}$ to $\Phi, \widetilde{R}$ to $R, \widetilde{\omega}_{\gamma}$ to $\omega_{\gamma}$, we then get the conclusions (1) and (3) of Theorem 1.1 and the estimates (1.10). From the fact that $\omega^{(j)}\left(I_{\rho ; \gamma}^{(j)}\right)=\Omega_{\gamma}$ for any $j$ we easily show that $\omega_{\gamma}$ maps $I_{\gamma}$ onto $\Omega_{\gamma}$. The measure estimate (1.9) follows, by Arnold's argument (cf. [1]), from the uniform boundedness of the Jacobian determinant of $\omega_{\gamma}$ from above and below by $(2 \Theta)^{n}$ and $\left(\frac{1}{2} \theta\right)^{n}$ respectively, which is easily observed from (3.1) and the fact that $d \omega_{\gamma}=\left(\gamma \Theta^{-1}\right)^{-1} d \tilde{\omega}_{\gamma}$. The proof of Theorem 1.1 is completed.

## 4. Proof of the iterative lemma

We set $I_{\rho ; \gamma}^{(0)}=\widetilde{\omega}^{-1}\left(\Omega_{\gamma}\right)$. Then Lemma 3.1 is valid for $j=0$. Assume the lemma is proven for indices $0, \cdots, j$. Then $C_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ T_{j} \circ \Phi_{j}$ is well-defined and real analytic on $V_{j}$ and commutes with $\pi_{k}, k=1, \cdots, n$ and

$$
\begin{equation*}
\left|C_{j}-I\right|_{V_{j}} \leq \gamma \varepsilon_{b}^{j}=s_{j+1}^{b-1} \cdot c_{2} \gamma \mu \delta_{b} \tag{4.1}
\end{equation*}
$$

by (iv). To prove Lemma 3.1 for $j+1$, we consider the map $B_{j}=R_{j}^{-1} \circ \Phi_{j}^{-1} \circ$ $T_{j+1} \circ \Phi_{j}$, as suggested in Section 2. First, we have
Lemma 4.1. If $\delta$ is small enough, then $B_{j}$ is well-defined and real analytic on

$$
V_{j}^{3}=I_{\rho ; \gamma}^{(j)} \times \mathbf{T}^{n}+\left(3 r_{j+1}, 3 s_{j+1}\right)
$$

and commutes with $\pi_{k}, k=1, \cdots, n$. Moreover, we have

$$
\begin{equation*}
\left|B_{j}-I\right|_{V_{j}^{3}} \leq 2 \gamma \varepsilon_{b}^{j} \tag{4.2}
\end{equation*}
$$

To prove this lemma, we may write, formally, $B_{j}=C_{j} \circ D_{j}$, where

$$
D_{j}=\Phi_{j}^{-1} \circ T_{j}^{-1} \circ T_{j+1} \circ \Phi_{j} .
$$

So far the mappings $T_{j+1}$ and $T_{j}^{-1}$ are defined only implicitly. First we need to determine the domains of definition of them.

Lemma 4.2. If $\delta$ is small, then $T_{j+1}$ is well defined on

$$
U_{j+1}^{*}=I_{\rho}^{*} \times \mathrm{T}^{n}+\left(\frac{13}{4} s_{j+1}, \frac{13}{4} s_{j+1}\right)
$$

and maps this domain into

$$
U_{j}^{* *}=I_{\rho}^{* *} \times \mathrm{T}^{n}+\left(\frac{14}{4} s_{j+1}, 7 s_{j+1}\right)
$$

on which $T_{j}^{-1}$ is well-defined, where

$$
I_{\rho}^{* *}=\left(I_{\rho}^{*}+\frac{1}{4}\right) \cap \mathbf{R}^{n} .
$$

Moreover, $T_{j+1}$ and $T_{j}^{-1}$ commute with $\pi_{k}, k=1, \cdots, n$ and,

$$
\begin{equation*}
\left|T_{j}^{-1} \circ T_{j+1}-I\right|_{U_{j+1}} \leq s_{j+1}^{b-1} \cdot 8 c_{b} \gamma \mu \delta_{b} . \tag{4.3}
\end{equation*}
$$

Proof of Lemma 4.2. It is observed that Lemma 4.2 of [10] does not apply to the proof of this lemma. But the contraction argument in an appropriate Banach space still works in our case. To express $T_{j+1}$ in explicit form, let us first solve $\hat{x}$ in terms of $x$ and $y$ from

$$
\begin{equation*}
\hat{x}=x-\frac{\partial f_{j+1}}{\partial y}(\hat{x}, y) \tag{4.4}
\end{equation*}
$$

the first equation of (2.11) where $j$ is replaced by $j+1$. Let $\mathcal{M}_{j+1}$ be the set of all real analytic $n$-valued functions $\eta(x, y)$, on $U_{j+1}^{*}$, which are $2 \pi$-periodic in the last $n$ variables and satisfy

$$
\begin{equation*}
|\eta|_{U_{j+1}^{*}} \leq \frac{1}{4}, \quad|\Im \eta|_{U_{j+1}} \leq \frac{1}{4} s_{j+1} \tag{4.5}
\end{equation*}
$$

Then $\mathcal{M}_{j+1}$ is a Banach space. Now consider the map

$$
\begin{equation*}
\mathcal{F}(\eta)(x, y)=-\frac{\partial f_{j+1}}{\partial y}(x+\eta(x, y), y) \tag{4.6}
\end{equation*}
$$

which is well-defined for $\eta \in \mathcal{M}_{j+1}$ and maps the space into itself because $f_{j+1}$ is real analytic on $\mathcal{U}_{j+1}$ and $2 \pi$-periodic in the last $n$ variables and, from (2.10) and Cauchy's estimates, for $\eta \in \mathcal{M}_{j+1}$ and $(x, y) \in U_{j+1}^{*}$,

$$
\begin{align*}
|\mathcal{F}(\eta)(x, y)| & \left.=\left|\frac{\partial f_{j+1}}{\partial y}(x+\eta(x, y), y)\right| \leq \sum_{k=0}^{j}\left|\frac{\partial}{\partial y}\left(f_{k+1}-f_{k}\right)\right|_{L_{\dot{p}} \times \mathbf{T}^{n}+\left(\frac{14}{4} s_{k+1} \frac{14}{4} s_{k+1}\right)}\right) \\
& \leq \sum_{k=0}^{j}\left(\frac{1}{2} s_{k+1}\right)^{-1}\left|f_{k+1}-f_{k}\right|_{u_{k+1}} \leq\left(\sum_{k=0}^{j} s_{k+1}^{a-1}\right) \cdot 2 c_{a} \gamma / \Lambda \delta_{a} \\
& \leq 2^{-(a-2)} s_{0}^{a-1} c_{a} \delta \leq \frac{1}{4}, \quad \text { if } \delta \leq \delta_{1}=4^{a-3} c_{a}^{-1} \tag{4.7}
\end{align*}
$$

and

$$
\begin{align*}
|\Im \mathcal{F}(\eta)(x, y)| & \leq\left|\frac{\partial f_{j+1}}{\partial y}(x+\eta(x, y), y)-\frac{\partial f_{j+1}}{\partial y}(\Re x+\Re \eta(x, y), \Re y)\right| \\
& \leq\left|\frac{\partial^{2} f_{j+1}}{\partial y \partial x}(\bar{x}, \bar{y})(\Im x+\Im \eta(x, y))+\frac{\partial^{2} f_{j+1}}{\partial y \partial y}(\bar{x}, \bar{y})(\Im y)\right| \\
& \leq 4^{-(a-2)} n c_{a} \delta \cdot 8 s_{j+1} \leq \frac{1}{4} s_{j+1}, \quad \text { if } \delta \leq \delta_{2}=4^{a-5}\left(n c_{a}\right)^{-1} \tag{4.8}
\end{align*}
$$

In the above, $(\bar{x}, \bar{y}) \in I_{\rho}^{* *} \times \mathbf{T}^{n}+\left(\frac{14}{4} s_{j+1}, \frac{14}{4} s_{j+1}\right)$ and, we have used the estimates, for example,

$$
\begin{align*}
\left|\frac{\partial^{2} f_{j+1}}{\partial y \partial x}(\bar{x}, \bar{y})\right| & \left.\leq \sum_{k=0}^{j}\left|\frac{\partial^{2}}{\partial y \partial x}\left(f_{k+1}-f_{k}\right)\right|_{I_{\vec{p}} \times \mathbf{T}^{n}+\left(\frac{14}{4} s_{k+1}, \frac{14}{4} s_{k+1}\right.}\right) \\
& \leq \sum_{k=0}^{j} \frac{2!}{\left(\frac{1}{2} s_{k+1}\right)^{2}}\left|f_{k+1}-f_{k}\right|_{u_{k+1}} \leq \sum_{k=0}^{j} s_{k+1}^{a-2} \cdot 8 c_{a} \delta_{a} \\
& \leq 2^{-(a-5)} s_{0}^{a-2} c_{a} \delta \leq 4^{-(a-2)} c_{a} \delta, \quad \text { if } 0<s_{0} \leq \frac{1}{4} \tag{4.9}
\end{align*}
$$

with the notice that $a=\alpha \lambda+\lambda+\tau>3 \tau+2>3 n+2 \geq 5$. By (4.9), we see that, for $\eta_{1}, \eta_{2} \in \mathcal{M}_{j+1},(x, y) \in U_{j+1}^{*}$,

$$
\begin{aligned}
\left|\mathcal{F}\left(\eta_{1}\right)(x, y)-\mathcal{F}\left(\eta_{2}\right)(x, y)\right| & =\left|\frac{\partial f_{j+1}}{\partial y}\left(x+\eta_{1}(x, y), y\right)-\frac{\partial f_{j+1}}{\partial y}\left(x+\eta_{2}(x, y), y\right)\right| \\
& \leq\left|\sum_{k=0}^{j} \frac{\partial^{2} f_{k+1}}{\partial y \partial x}\left(x_{k}^{*}, y\right)\left(\eta_{1}(x, y)-\eta_{2}(x, y)\right)\right| \\
& \leq \frac{1}{2}\left|\eta_{1}-\eta_{2}\right|_{U_{j+1}^{*}}, \quad \text { if } \delta \leq \delta_{2} .
\end{aligned}
$$

In the above equation, for each $k, x_{k}^{*}$ is a point in $I_{\rho}^{* *}+\frac{14}{4} s_{j+1}$. This shows that the map $\mathcal{F}: \mathcal{M}_{j+1} \rightarrow \mathcal{M}_{j+1}$ is contractive. Therefore, there exists a unique $y^{*} \in \mathcal{M}_{j+1}$ such that

$$
\begin{equation*}
\mathcal{F}\left(y^{*}\right)=y^{*}, \tag{4.10}
\end{equation*}
$$

and $T_{j+1}:(x, y) \rightarrow(\hat{x}, \hat{y})$ is expressed explicitly in the form

$$
\left\{\begin{array}{l}
\hat{x}=x+y^{*}(x, y)  \tag{4.11}\\
\hat{y}=y+\widetilde{\omega}\left(x+y^{*}(x, y)\right)+\frac{\partial f_{+1}}{\partial x}\left(x+y^{*}(x, y), y\right)
\end{array}\right.
$$

which is clearly well-defined and real analytic for $(x, y) \in U_{j+1}^{*}$ and commutes with $\pi_{k}, k=1, \cdots, n$. It is easy to check that $T_{j+1}$ maps $U_{j+1}^{*}$ into $U_{j}^{* *}$ by using the condition (2.8) for $\tilde{\omega}$ and the similar arguments to the above. Also, the contraction arguments may be applied to show that $T_{j}^{-1}$ is well-defined and real analytic on $U_{j}^{* *}$ and commutes with $\pi_{k}, k=1, \cdots, n$ only but, instead of (4.6), here we need to consider the following map

$$
\begin{equation*}
\mathcal{G}(\eta)(x, y)=-\frac{\partial f_{j}}{\partial x}(x, y-\widetilde{\omega}(x)+\eta(x, y)) \tag{4.12}
\end{equation*}
$$

which maps the Banach space, say $\mathcal{M}_{j}^{*}$, of real analytic $n$-valued functions $\eta(x, y)$ with $2 \pi$-periodic in the last $n$ variables $y$ for $(x, y)$ in $U_{j}^{* *}$ and with

$$
\begin{equation*}
|\eta|_{U_{j}^{* *}} \leq \frac{1}{4}, \quad|\Im \eta|_{U_{j}^{* *}} \leq \frac{1}{4} s_{j} \tag{4.13}
\end{equation*}
$$

into itself contractively if $\delta$ is small enough (say, $\delta \leq \delta_{2}$ ). Therefore we have a unique fixed point $y^{* *} \in \mathcal{M}_{j}^{*}$ of the $\operatorname{map} \mathcal{G}$ and $T_{j}^{-1}$ is explicitly given by

$$
\left\{\begin{array}{l}
\hat{x}=x+\frac{\partial f_{j}}{\partial y}\left(x, y-\tilde{\omega}(x)+y^{* *}(x, y)\right)  \tag{4.14}\\
\hat{y}=y-\tilde{\omega}(x)+y^{* *}(x, y)
\end{array}\right.
$$

Simple calculations will show that $T_{j}^{-1}$ maps $U_{j}^{* *}$ into

$$
U_{j}^{* *}=I_{\rho}^{* *} \times \mathrm{T}^{n}+\left(2 s_{j}, 3 s_{j}\right)
$$

with $I_{\rho}^{* * *}=\left(I_{\rho}^{*}+\frac{1}{2}\right) \cap \mathrm{R}^{n}$. It remains to verify (4.3). Note that $T_{j}^{-1} \circ T_{j+1}$ : $(x, y) \rightarrow(\hat{x}, \hat{y})$ can be written in the form

$$
\left\{\begin{array}{l}
\hat{x}=x-\left(\frac{\partial f_{j+1}}{\partial y}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial y}(\widetilde{x}, \hat{y})\right)  \tag{4.15}\\
\hat{y}=y+\left(\frac{\partial f_{j+1}}{\partial x}(\widetilde{x}, y)-\frac{\partial f_{1}}{\partial x}(\widetilde{x}, \hat{y})\right)
\end{array}\right.
$$

where $(\widetilde{x}, \tilde{y})=T_{j+1}(x, y)$. Solve $\hat{y}$ in terms of $x$ and $y$ from the second equation of (4.15), we get

$$
\begin{equation*}
\hat{y}=y+y^{* * *}(\widetilde{x}, y) \tag{4.16}
\end{equation*}
$$

where $y^{* * *}(\widetilde{x}, y)$ is the fixed point of the contractive map

$$
\begin{equation*}
\mathcal{H}(\eta)(\widetilde{x}, y)=\frac{\partial f_{j+1}}{\partial x}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial x}(\tilde{x}, y+\eta(\tilde{x}, y)) \tag{4.17}
\end{equation*}
$$

on the Banach space, say $\mathcal{M}_{j}^{* *}$, of real analytic $n$-valued functions $\eta(\tilde{x}, y)$ with $2 \pi$-periodic in variables $y$ for $(\widetilde{x}, y) \in \tilde{U}_{j}^{* *}=I_{\rho}^{* *} \times \mathbf{T}^{n}+\left(\frac{14}{4} s_{j+1}, \frac{13}{4} s_{j+1}\right)$ and with

$$
\begin{equation*}
|\eta|_{\tilde{U}_{j}^{*}} \leq \frac{1}{4} s_{j+1} \tag{4.18}
\end{equation*}
$$

The well-definedness and the contractivity of the map $\mathcal{H}$ on $\mathcal{M}_{j}^{* *}$ are easily proved by the previous arguments with the notice of the inequality (2.10) for $f_{j+1}-f_{j}$. With the well-defined $y^{* * *}(\tilde{x}, y)$ and $(\widetilde{x}, \tilde{y})=T_{j+1}(x, y)$, the map $T_{j}^{-1} \circ T_{j+1}$ : $(x, y) \rightarrow(\hat{x}, \hat{y})$ is given explicitly by (4.15) the second equation of which is in fact equivalent to (4.16). Direct verification shows that, if $\delta \leq \delta_{2}$,

$$
\begin{aligned}
\left|y^{* * *}\right| \widetilde{U}_{j}^{*} & =\sup _{(\widetilde{x}, y) \in \widetilde{U}_{j}^{*}}\left|\frac{\partial f_{j+1}}{\partial x}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial x}\left(\widetilde{x}, y+y^{* * *}(\widetilde{x}, y)\right)\right| \\
& \leq \sup _{(\widetilde{x}, y) \in \widetilde{U}_{j} \cdot}\left|\frac{\partial f_{j+1}}{\partial x}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial x}(\widetilde{x}, y)\right| \\
& +\sup _{(\widetilde{x}, y) \in \widetilde{U}_{j}^{*}}\left|\frac{\partial f_{j}}{\partial x}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial x}\left(\widetilde{x}, y+y^{* * *}(\widetilde{x}, y)\right)\right| \\
& \leq s_{j+1}^{b-1} \cdot 4 c_{b} \gamma \mu \delta_{b}+\frac{1}{2}\left|y^{* * *}\right|_{\tilde{U}_{j}} .
\end{aligned}
$$

and therefore

$$
\begin{equation*}
\left|y^{* * *}\right|_{\tilde{u}_{j}^{*}} \leq s_{j+1}^{b-1} \cdot 8 c_{b} \gamma \mu \delta_{b} . \tag{4.19}
\end{equation*}
$$

In a similar way, we get

$$
\begin{equation*}
\sup _{(\widetilde{x}, y) \in \widetilde{U}_{j}^{*}}\left|\frac{\partial f_{j+1}}{\partial y}(\widetilde{x}, y)-\frac{\partial f_{j}}{\partial y}\left(\widetilde{x}, y+y^{* *}(\widetilde{x}, y)\right)\right| \leq s_{j+1}^{b-1} \cdot 8 c_{b} \gamma \mu \delta_{b} \tag{4.20}
\end{equation*}
$$

This verifies (4.3) and completes the proof of Lemma 4.2.

Now we return to the proof of Lemma 4.1. By inequality (3.2) for indices $0,1, \cdots, j$ and Cauchy's estimates, we easily prove that, if $\delta \leq \delta_{3}$ with

$$
\delta_{3}=\left(8 n c_{1, a} \sum_{k=1}^{\infty} r_{k}^{\alpha-1}\right)^{-1}=\frac{4^{(\alpha-1) \lambda}-1}{8 n c_{1, a} s_{0}^{(\alpha-1) \lambda}}
$$

then

$$
\begin{equation*}
\frac{1}{2}\left|z_{1}-z_{2}\right| \leq\left|\Phi_{j}\left(z_{1}\right)-\Phi_{j}\left(z_{2}\right)\right| \leq 2\left|z_{1}-z_{2}\right| \tag{4.21}
\end{equation*}
$$

for $z_{1} \in V_{j}$ and $\left|z_{1}-z_{2}\right| \leq \frac{1}{4} r_{j}$, which implies that $\left|D \Phi_{j}\right|_{V_{j}},\left|D \Phi_{j}^{-1}\right|_{V_{j}} \leq 2$. This, together with (3.2), implies that $\Phi_{j}\left(V_{j}\right) \subset U_{j+1}^{*}$. From (4.21), we also have

$$
\begin{equation*}
\Phi_{j}\left(V_{j}^{3}\right)+\frac{1}{8} r_{j} \subset \Phi_{j}\left(V_{j}\right) \tag{4.22}
\end{equation*}
$$

By (4.3) and the fact that $s_{j+1}^{a-1} \cdot 8 c_{a} \gamma \mu \delta \leq \frac{1}{8} r_{j}$ if $\delta \leq \delta_{4}=4^{\lambda-3} c_{a}^{-1}$, we see that $D_{j}$ is well-defined on $V_{j}^{3}$ and maps this domain into $V_{j}$. Therefore, $B_{j}$ is welldefined on $V_{j}^{3}$. On the other hand, from (4.3) and the fact that $\left|D \Phi_{j}^{-1}\right|_{V_{j}} \leq 2$ and $c_{2} \geq 16 c_{b}$, we get

$$
\left|D_{j}-1\right|_{V_{j}^{3}} \leq 2\left|T_{j}^{-1} \circ T_{j+1} \circ \Phi_{j}-\Phi_{j}\right|_{V_{j}^{3}} \leq 2\left|T_{j}^{-1} \circ T_{j+1}-I\right|_{U_{j+1}^{*}} \leq \gamma \varepsilon_{b}^{j} .
$$

Therefore,

$$
\left|B_{j}-I\right|_{V_{j}^{3}} \leq\left|C_{j} \circ D_{j}-D_{j}\right|_{V_{j}^{3}}+\left|D_{j}-I\right|_{V_{j}^{3}} \leq\left|C_{j}-I\right|_{V_{j}}+\left|D_{j}-I\right|_{V_{j}^{3}} \leq 2 \gamma \varepsilon_{b}^{j}
$$

It is clear that $B_{j}$ commutes with $\pi_{k}, k=1, \cdots, n$. Lemma 4.1 is then proved.
$B_{j}$ is an exact symplectic map. So there exists a generating function, say $b_{j}$, such that $B_{j}:(x, y) \rightarrow(\hat{x}, \hat{y})$ is generated from $b_{j}$ by (2.20). By the exact symplecticity of $B_{j}$ and the previous contraction arguments, we have
Lemma 4.3. If $\delta \leq \delta_{5}=\left(8 n c_{2, a}\right)^{-1}$, then $b_{j}$ is well-defined and real analytic on

$$
\tilde{V}_{j}^{*}=I_{\rho ; \gamma}^{(j)} \times \mathrm{T}^{n}+\left(\frac{5}{2} r_{j+1}, \frac{5}{2} s_{j+1}\right)
$$

and is $2 \pi$-periodic in the last $n$ variables. Moreover, we have

$$
\begin{equation*}
\left|\frac{\partial b_{j}}{\partial x}\right|_{\widetilde{v}_{j}^{*}},\left|\frac{\partial b_{j}}{\partial y}\right|_{\tilde{v}_{j}^{*}} \leq 2 \gamma \varepsilon_{b}^{j} . \tag{4.23}
\end{equation*}
$$

Now we may define

$$
\begin{equation*}
F_{0}^{(j+1)}(x)=F_{0}^{(j)}(x)+\left[b_{j}\right](x), \quad \omega^{(j+1)}(x)=\frac{\partial F_{0}^{(j+1)}}{\partial x}(x), \tag{4.24}
\end{equation*}
$$

where $\left[b_{j}\right](x)=\frac{1}{(2 \pi)^{n}} \int_{T^{n}} b_{j}(x, y) d^{n} y$. Then $F_{0}^{(j+1)}$ and $\omega^{(j+1)}$ are well-defined and real analytic on $\tilde{V}_{j}^{*}$ and independent of the angle variables. Below we first show that there exists a closed subset $I_{\rho ; \gamma}^{(j+1)}$ of $I_{\rho}$ such that (i), (ii) anf (v) of Lemma 3.1 are true for index $j+1$. For this, we state the following lemma cited from [9] in which the proof can be found.

Lemma 4.4. Assume that $F: D \rightarrow \mathbf{R}^{n}, G: D \rightarrow \mathbf{R}^{n}$ are two continuously differentiable mappings where $D$ is an open subset of $\mathbf{R}^{n}$. Let $C$ be a bounded open set such that $\bar{C} \subset D$ where $\bar{C}$ denotes the closure of $C$. Let $y \in \mathbf{R}^{n}$ satisfy

$$
L=\min _{x \in \partial C}\left\{\|F(x)-y\|_{2}\right\}>0 .
$$

If

$$
\sup _{x \in ट}\left\{\|F(x)-G(x)\|_{2}\right\} \leq \frac{L}{8}
$$

then $\operatorname{deg}(F, C, y)=\operatorname{deg}(G, C, y)$. Where $\operatorname{deg}(F, C, y)$ denotes the degree of the mapping $F$ associated to the point $y$ and the region $C$, and $\|\cdot\|_{2}$ denotes the usual Euclidean metric.

We continue our proof. To apply Lemma 4.4, we fix a point $x^{*} \in I_{p ; \gamma}^{(j)}$ and let $F=\omega^{(j)}, G=\omega^{(j+1)}, D=\left\{x^{*}\right\}+\frac{1}{2} r_{j}, C=\left\{x^{*}\right\}+\frac{1}{4} r_{j}, y=\omega^{*}=\omega^{(j)}\left(x^{*}\right)$. Since $0<s_{0} \leq 1$, we have from (3.1)

$$
\begin{equation*}
\frac{2}{3} \gamma \mu\left|x_{1}-x_{2}\right| \leq\left|F\left(x_{1}\right)-F\left(x_{2}\right)\right| \leq \frac{4}{3} \gamma\left|x_{1}-x_{2}\right|, \quad x_{1}, x_{2} \in D \tag{4.25}
\end{equation*}
$$

Therefore,

$$
L=\min _{x \in \partial C}\left\{\|F(x)-y\|_{2}\right\} \geq \min _{x \in \partial C}\left\{\left\{\omega^{(j)}(x)-\omega^{(j)}\left(x^{*}\right) \mid\right\} \geq \frac{1}{6} \gamma \mu r_{j}\right.
$$

and

$$
\sup _{x \in C}\left\{\|F(p)-G(p)\|_{2}\right\} \leq \sqrt{n}\left|\frac{\partial b_{j}}{\partial x}\right|_{\tilde{V}_{j}^{*}} \leq \frac{L}{8}, \quad \text { if } \delta \leq \delta_{5} .
$$

By Lemma 4.4, we have $\operatorname{deg}\left(\omega^{(j+1)}, C, \omega^{*}\right)=\operatorname{deg}\left(\omega^{(j)}, C, \omega^{*}\right)=1$. From the theorem of topology degree of a mapping, there exists a point $x_{1}^{*} \in C$ such that $\omega^{(j+1)}\left(x_{1}^{*}\right)=\omega^{*}$. Moreover, we have

$$
\begin{aligned}
\left|x_{1}^{*}-x^{*}\right| & \leq\left(\frac{2}{3} \gamma \mu\right)^{-1}\left|\omega^{(j)}\left(x_{1}^{*}\right)-\omega^{(j)}\left(x^{*}\right)\right| \\
& \leq\left(\frac{2}{3} \gamma \mu\right)^{-1}\left|\frac{\partial\left[b_{j}\right]}{\partial x}\left(x_{1}^{*}\right)\right| \leq \frac{1}{4} r_{j+1}
\end{aligned}
$$

Let $I_{\rho ; \gamma}^{(j+1)}$ be the set of points $x_{1}^{*} \in I_{p}$ such that $\omega^{(j+1)}\left(x_{1}^{*}\right)=\omega^{(j)}\left(x^{*}\right)$ for $x^{*} \in I_{\rho ; \gamma}^{(j)}$. Then the conclusion (i) of Lemma 3.1 is obviously true for $j+1$ and $\omega^{(j+1)}$ maps $I_{\rho ; \gamma}^{(j+1)}$ onto $\Omega_{\gamma}$. So $\omega^{(j+1)}$ is well-defined on $V_{j+1}$ and (3.4) is also valid for $j+1$ with $c_{3}=2 c_{2}$. Seeing that, for $x_{1}, x_{2} \in I_{p ; \gamma}^{(j+1)}+r_{j+1}$ with $\left|x_{1}-x_{2}\right| \leq r_{j+1}$,

$$
\begin{align*}
\left|\frac{\partial\left[b_{j}\right]}{\partial x}\left(x_{1}\right)-\frac{\partial\left[b_{j}\right]}{\partial x}\left(x_{2}\right)\right| & \leq \frac{1}{(2 \pi)^{n}} \int_{\mathrm{T}^{n}}\left|\frac{\partial^{2} b_{j}}{\partial x \partial x}(\widetilde{x}, y)\left(x_{1}-x_{2}\right)\right| d^{n^{n} y} \\
& \leq n \cdot r_{j+1}^{-1} \cdot\left|\frac{\partial b_{j}}{\partial x}\right|_{\tilde{v}_{j}^{*}} \cdot\left|x_{1}-x_{2}\right| \leq s_{j+1} \cdot \gamma \mu\left|x_{1}-x_{2}\right|( \tag{4.26}
\end{align*}
$$

if $\delta \leq \delta_{5}$ with $\widetilde{x} \in \overline{x_{1} x_{2}}$, the straight line connecting $x_{1}$ and $x_{2}$ and completely contained in $I_{\rho ; \gamma}^{(j)}+\frac{5}{2} r_{j+1}$, the equation (3.1) is then verified for $j+1$. In (4.26) we have applied the Cauchy's estimate to analytic function $\frac{\partial b_{j}}{\partial x}$. It remains to verify (iii) and (iv) for $j+1$. For this, we define $\Psi_{j}$.

Note that the Fourier series expansion

$$
b_{j}(x, y)=\left[b_{j}\right](x)+\sum_{0 \neq k \in \mathbf{Z}^{n}} b_{j ; k}(x) e^{i(k, y\rangle}
$$

is defined on $\tilde{V}_{j}^{*}$. Let $\tilde{b}_{j}(x, y)=\sum_{0 \neq k \in \mathbf{Z}^{n}} b_{j ; k}(x) e^{i(k, y\rangle}$. We take generating function $\psi_{j}(x, y)$, of the symplectic transformation $\Psi_{j}$, as the solution of the equation

$$
\begin{equation*}
\psi_{j}\left(x, y+\omega^{(j)}(x)\right)-\psi_{j}(x, y)+T_{m_{j}} \widetilde{b}_{j}(x, y)=0 \tag{4.27}
\end{equation*}
$$

with $\left[\psi_{j}\right](x)=0$, where

$$
T_{m_{j}} \tilde{b}_{j}(x, y)=\sum_{0<|k| \leq m_{j}} b_{j ; k}(x) e^{i\langle k, y\rangle}
$$

with $m_{j}$ given by

$$
\begin{equation*}
m_{j}^{\tau+1}=\frac{1}{r_{j}} \tag{4.28}
\end{equation*}
$$

We find

$$
\begin{equation*}
\psi_{j}(x, y)=-\sum_{0<|k| \leq m_{j}} \frac{b_{j ; k}(x)}{e^{i\left(k, \omega^{j j}(x)\right\rangle}-1} e^{i(k, y\rangle} \tag{4.29}
\end{equation*}
$$

Lemma 4.5. $\psi_{j}(x, y)$ is well-defined and real analytic on $\tilde{V}_{j}^{*}$ and is $2 \pi$-periodic in each variables of $y$. Moreover, if $f \in C^{b}\left(I_{\rho} \times \mathrm{T}^{n}\right)$, then we have the following estimates for its gradient:

$$
\begin{equation*}
\left|\frac{\partial \psi_{j}}{\partial x}(x, y)\right| \leq \sigma_{1} s_{j+1}^{-2 \tau} \varepsilon_{b}^{j}, \quad\left|\frac{\partial \psi_{j}}{\partial y}(x, y)\right| \leq \sigma_{1} s_{j+1}^{-\tau} \varepsilon_{b}^{j} \tag{4.30}
\end{equation*}
$$

for $(x, y)$ in

$$
V_{j}^{*}=I_{\rho ; \gamma}^{(j)}+\left(\frac{9}{4} r_{j+1}, \frac{9}{4} s_{j+1}\right)
$$

where $\sigma_{1}$ is a positive constant depending only on $n$ and $\tau$.
Proof: The well-definedness of $\psi_{j}$ is proved by simply verifying the nonvanishing of $\left(e^{i\left\langle k, \omega^{(j)}(x)\right\rangle}-1\right)$ for $x \in I_{\rho ; \gamma}^{(j)}+\frac{5}{2} r_{j+1}$ and $0<|k| \leq m_{j}$. By the definitions of $I_{\rho ; \gamma}^{(j)}$ and $m_{j}$ with the notice that $r_{j+1}=4^{-\lambda} r_{j}$ and that $\lambda \geq 2$, we get for $x^{*} \in I_{\rho ; \gamma}^{(j)}$, $\left|x-x^{*}\right| \leq \frac{5}{2} r_{j+1}$ and $0<|k| \leq m_{j}$,

$$
\begin{aligned}
\left|e^{i\left\langle k, \omega^{(j)}(x)\right\rangle}-1\right| & \geq\left|e^{i\left\langle k, \omega^{(j)}\left(x^{*}\right)\right\rangle}-1\right|-\left|e^{i\left\langle k, \omega^{(j)}(x)\right\rangle}-e^{i\left\langle k, \omega^{(j)}\left(x^{*}\right)\right\rangle}\right| \\
& \geq \frac{\gamma}{|k|^{\tau}}-2\left|\left\langle k, \omega^{(j)}(x)-\omega^{(j)}\left(x^{*}\right)\right\rangle\right| \\
& \geq \frac{\gamma}{|k|^{\tau}}-2 m_{j} \cdot 2 \gamma \cdot \frac{5}{2} r_{j+1} \geq \frac{1}{4} \cdot \frac{\gamma}{|k|^{\tau}}
\end{aligned}
$$

which does not vanish. The analyticity and the periodicity in the last $n$ variables of this function are clear. The remainder of this lemma can be proved by differentiating the equation (4.27) with respect to $x$ and $y$, and then estimating $\left|\frac{\partial \psi_{j}}{\partial x}\right|$ and $\left|\frac{\partial \psi_{j}}{\partial y}\right|$ over $V_{j}^{*}$ in terms of the norms $\left|\frac{\partial b_{j}}{\partial x}\right|_{\widetilde{V}_{j}}$ and $\left|\frac{\partial b_{j}}{\partial y}\right|_{\widetilde{V}_{j}}$ respectively, by using the standard arguments. The details are referred to [10] and [13] for a similar problem.

Therefore, $\Psi_{j}:(\xi, \eta) \rightarrow(x, y)$ is well-defined from $\psi_{j}$ by (2.18). More precisely, we have the following result which is easily proved by the standard contraction argument.
Lemma 4.6. If $\delta \leq \delta_{6}=\left(8 n \sigma_{1} c_{2, a}\right)^{-1}$, then $\Psi_{j}$ and its inverse $\Psi_{j}^{-1}$ are welldefined on $V_{j}^{2}$ and map $V_{j}^{2}$ into $V_{j}^{*}$ with

$$
\begin{equation*}
\left|\Psi_{j}-I\right|_{V_{j}^{2}} \leq s_{j+1}^{-2 \tau} \cdot \sigma_{1} \varepsilon_{b}^{j} \tag{4.31}
\end{equation*}
$$

and $\Psi_{j}^{-1} \circ R_{j} \circ B_{j} \circ \Psi_{j}$ is well defined and real analytic on $V_{j+1}$ and maps $V_{j+1}$ into $\tilde{V}_{j}^{*}$.

Let $\Phi_{j+1}=\Phi_{j} \circ \Psi_{j}$. Then, $\Phi_{j+1}$ is well-defined and real analytic on $V_{j}^{2}$ and commutes with $\pi_{k}, k=1, \cdots, n$ and, from (4.21),

$$
\left|\Phi_{j+1}-\Phi_{j}\right|_{V_{j}^{2}}=\left|\Phi_{j} \circ \Psi_{j}-\Phi_{j}\right|_{V_{j}^{2}} \leq 2\left|\Psi_{j}-I\right|_{V_{j}^{2}} \leq r_{j+1}^{\beta} \cdot c_{1} \mu \delta_{b}
$$

with $c_{1}=2 \sigma_{1} c_{2}$. (3.2) is then verified. Next we verify (3.3) for $j+1$. Note that $C_{j+1}=R_{j+1}^{-1} \circ \Psi_{j}^{-1} \circ R_{j} \circ B_{j} \circ \Psi_{j}$, which is analytically defined on $V_{j+1}$ and commutes with $\pi_{k}, k=1, \cdots, n$ and, moreover, it may be written as $(\xi, \eta) \rightarrow(\hat{\xi}, \hat{\eta})$ with

$$
\begin{aligned}
& \hat{\xi}=\xi-\frac{\partial \psi_{j}}{\partial y}(\hat{\xi}, \hat{y})+\frac{\partial \psi_{j}}{\partial y}(\xi, y)-\frac{\partial b_{j}}{\partial y}(\hat{x}, y) \\
& \hat{\eta}=\eta+\frac{\partial \psi_{j}}{\partial \xi}(\hat{\xi}, \hat{y})-\frac{\partial \psi_{j}}{\partial \xi}(\xi, y)+\frac{\partial b_{j}}{\partial x}(\hat{x}, y)-\omega^{(j+1)}(\hat{\xi})+\omega^{(j)}(\hat{x})
\end{aligned}
$$

where $(x, y)=\Psi_{j}(\xi, \eta) \in T_{\beta ; \gamma}^{(j)} \times \mathbf{T}^{n}+\left(\frac{6}{4} r_{j+1}, \frac{6}{4} s_{j+1}\right),(\hat{x}, \hat{y})=R_{j} \circ B_{j}(x, y) \in$ $I_{\rho ; \gamma}^{(j)} \times \mathrm{T}^{n}+\left(\frac{7}{4} r_{j+1}, 2 s_{j+1}\right),(\hat{\xi}, \hat{\eta})=R_{j+1}^{-1} \circ \Psi_{j}^{-1}(\hat{x}, \hat{y}) \in I_{\rho ; \gamma}^{(j)} \times \mathrm{T}^{n}+\left(2 r_{j+1}, \frac{10}{4} s_{j+1}\right)$ for $(\xi, \eta) \in V_{j+1}$ if $\delta \leq \delta_{6}$. So,

$$
\begin{aligned}
|\hat{\xi}-\xi| & =\left|\frac{\partial \psi_{j}}{\partial y}(\hat{\xi}, \hat{y})-\frac{\partial \psi_{j}}{\partial y}(\xi, y)+\frac{\partial b_{j}}{\partial y}(\hat{x}, y)\right| \\
& \leq I_{1}+I_{2}+I_{3}+I_{4} \\
|\hat{\eta}-\eta| & =\left|\frac{\partial \psi_{j}}{\partial \xi}(\hat{\xi}, \hat{y})-\frac{\partial \psi_{j}}{\partial \xi}(\xi, y)+\frac{\partial b_{j}}{\partial x}(\hat{x}, y)-\omega^{(j+1}(\hat{\xi})+\omega^{(j)}(\hat{x})\right| \\
& \leq J_{1}+J_{2}+J_{3}+J_{4}+J_{5}+J_{6}+J_{7}
\end{aligned}
$$

where

$$
\begin{aligned}
& I_{1}=\left|\frac{\partial \psi_{j}}{\partial y}(\hat{\xi}, \hat{y})-\frac{\partial \psi_{j}}{\partial y}(\xi, \hat{y})\right|, \\
& I_{2}=\left|\frac{\partial \psi_{j}}{\partial y}(\xi, \hat{y})-\frac{\partial \psi_{j}}{\partial y}\left(\xi, y+\omega^{(j)}(\xi)\right)\right|, \\
& I_{3}=\left|\frac{\partial \psi_{j}}{\partial y}\left(\xi, y+\omega^{(j)}(\xi)\right)-\frac{\partial \psi_{j}}{\partial y}(\xi, y)+\frac{\partial b_{j}}{\partial y}(\xi, y)\right|, \\
& I_{4}=\left|\frac{\partial b_{j}}{\partial y}(\hat{x}, y)-\frac{\partial b_{j}}{\partial y}(\xi, y)\right|, \\
& J_{1}=\left|\frac{\partial \psi_{j}}{\partial \xi}(\hat{\xi}, \hat{y})-\frac{\partial \psi_{j}}{\partial \xi}(\xi, \hat{y})\right|, \\
& J_{2}=\left|\frac{\partial \psi_{j}}{\partial \xi}(\xi, \hat{y})-\frac{\partial \psi_{j}}{\partial \xi}\left(\xi, y+\omega^{(j)}(\xi)\right)\right|, \\
& J_{3}=\left|\frac{\partial \psi_{j}}{\partial \xi}\left(\xi, y+\omega^{(j)}(\xi)\right)-\frac{\partial \psi_{j}}{\partial \xi}(\xi, y)-\frac{\partial \omega^{(j)}}{\partial \xi}(\xi) \frac{\partial \psi_{j}}{\partial y}\left(\xi, y+\omega^{(j)}(\xi)\right)+\frac{\partial \tilde{b}_{j}}{\partial x}(\xi, y)\right|,
\end{aligned}
$$

$$
\begin{aligned}
& J_{4}=\left|\frac{\partial b_{j}}{\partial x}(\hat{x}, y)-\frac{\partial b_{j}}{\partial x}(\xi, y)\right| \\
& J_{5}=\left|\omega^{(j)}(\hat{x})-\omega^{(j)}(\xi)-\frac{\partial \omega^{(j)}}{\partial \xi}(\xi)(\hat{x}-\xi)\right|, \\
& J_{6}=\left|\frac{\partial \omega^{(j)}}{\partial \xi}(\xi)\left(\frac{\partial \psi_{j}}{\partial y}\left(\xi, y+\omega^{(j)}(\xi)\right)-\frac{\partial \psi_{j}}{\partial y}(\xi, y)+\frac{\partial \tilde{b}_{j}}{\partial y}(\hat{x}, y)\right)\right|, \\
& J_{7}=\left|\frac{\partial\left[b_{j}\right]}{\partial x}(\hat{\xi})-\frac{\partial\left[b_{j}\right]}{\partial x}(\xi)\right|+\left|\omega^{(j)}(\hat{\xi})-\omega^{(j)}(\xi)\right| .
\end{aligned}
$$

By Taylor formula and Cauchy's estimates with the notice that all the concerned variables are in the corresponding shrunken domains with shrunken width, say, $\left(\frac{1}{4} r_{j+1}, \frac{1}{4} s_{j+1}\right)$, we obtain

$$
\begin{aligned}
& I_{1} \leq n \cdot\left(\frac{1}{4} r_{j+1}\right)^{-1}\left|\frac{\partial \psi_{j}}{\partial y}\right|_{V_{j}^{*}} \cdot|\hat{\xi}-\xi| \leq s_{j+1}^{-\lambda-\tau} \cdot 4 n \sigma_{1} \varepsilon_{a}^{j} \cdot|\hat{\xi}-\xi|, \\
& I_{2} \leq n \cdot\left(\frac{1}{4} s_{j+1}\right)^{-1}\left|\frac{\partial \psi_{j}}{\partial y}\right|_{V_{j}^{*}} \cdot\left|\hat{y}-y-\omega^{(j)}(\xi)\right| \leq s_{j+1}^{-2 \tau-1} \cdot \sigma_{2} \varepsilon_{a}^{j} \cdot \gamma \varepsilon_{b}^{j}
\end{aligned}
$$

with $\sigma_{2}=8 n \sigma_{1}\left(\sigma_{1}+3\right)$, here the following estimate has been used

$$
\begin{align*}
\left|\hat{y}-y-\omega^{(j)}(\xi)\right| & =\left|\omega^{(j)}(\hat{x})-\omega^{(j)}(\xi)+\frac{\partial b_{j}}{\partial x}(\hat{x}, y)\right| \\
& \leq 2 \gamma|\hat{x}-\xi|+\left|\frac{\partial b_{j}}{\partial x}\right|_{\tilde{v}_{j}} \\
& \leq 2 \gamma\left|\frac{\partial b_{j}}{\partial y}\right|_{\tilde{v}_{j}^{*}}+2 \gamma\left|\frac{\partial \psi_{j}}{\partial y}\right|_{V_{j}^{*}}+\left|\frac{\partial b_{j}}{\partial x}\right|_{\tilde{V}_{j}^{*}} \\
& \leq s_{j+1}^{-\tau} \cdot 2\left(\sigma_{1}+3\right) \cdot \gamma \varepsilon_{b}^{j} . \tag{4.32}
\end{align*}
$$

$I_{3}$ is bounded by (cf. [10], p. 684)

$$
I_{3} \leq\left|\frac{\partial b_{j}}{\partial y}-T_{m_{j}} \frac{\partial b_{j}}{\partial y}\right|_{I_{i ;}^{(j)} \times T^{n}+\left(\frac{6}{4} r_{j+1}, \frac{\epsilon_{4}^{4}}{4} s_{j+1}\right)} \leq s_{j+1}^{a-\lambda-\tau-1} \cdot \gamma \varepsilon_{b}^{j},
$$

if $s_{0}$ is small enough but depends only on $n, \tau, \lambda, \alpha$. From (4.32), we get the estimate $|\hat{x}-\xi| \leq s_{j+1}^{-\tau} \cdot\left(\sigma_{1}+2\right) \cdot \varepsilon_{b}^{j}$, which implies

$$
I_{4} \leq n \cdot\left(\frac{1}{4} r_{j+1}\right)^{-1}\left|\frac{\partial b_{j}}{\partial y}\right|_{\widetilde{v}_{j}} \cdot|\hat{x}-\xi| \leq s_{j+1}^{-\lambda-\tau} \cdot 8 n\left(\sigma_{1}+2\right) \varepsilon_{a}^{j} \cdot \gamma \varepsilon_{b}^{j}
$$

Noting that $\varepsilon_{a}^{j}=s_{j+1}^{a-1} \cdot c_{2, a} \mu \delta_{a} \leq s_{j+1}^{a-1} \cdot \sigma_{2}^{-1}$ if $\delta$ is small enough, say,

$$
\delta \leq \delta_{7}=\left(\sigma_{2} c_{2, a}\right)^{-1}
$$

we get, by combining the above estimates,

$$
\begin{equation*}
|\hat{\xi}-\xi| \leq s_{j+1}^{a-\lambda-\tau-1} \cdot 6 \gamma \varepsilon_{b}^{j} . \tag{4.33}
\end{equation*}
$$

By estimating $J_{k}, k=1, \cdots, 7$ in a similar way to the above with making use of the previous estimates, we obtain

$$
\begin{equation*}
|\hat{\eta}-\eta| \leq s_{j+1}^{a-\lambda-2 \tau-1} \cdot 6 \gamma \varepsilon_{b}^{j} \tag{4.34}
\end{equation*}
$$

if $\delta$ is smaller, say,

$$
\delta \leq \delta_{8}=\left(\sigma_{2}^{2} c_{2, a}\right)^{-1}
$$

and $s_{0}$ is chosen to be small as before. To summarize, if $\delta \leq \delta_{0}=\min _{1 \leq l \leq 8} \delta_{l}$ with $s_{0}$ sufficiently small and depending on $n, \tau, \lambda, \alpha$, then

$$
\begin{equation*}
\left|C_{j+1}-I\right|_{V_{j+1}} \leq 6 s_{j+1}^{(\alpha-1) \lambda} \cdot \gamma \varepsilon_{b}^{j} \tag{4.35}
\end{equation*}
$$

due to the fact that $(\alpha-1) \lambda<a-\lambda-2 \tau-1<a-\lambda-\tau-1$. As argued in [10], p. 688, for $b=a,(4.35)$ is bounded by $\gamma \varepsilon_{a}^{j+1}$ with an appropriate choice of $s_{0}$, say, $s_{0}^{(\alpha-1) \lambda} \leq 4^{-(a-1)} / 6$. And for any finite $b \geq a$, we can also bound (4.35) by $\gamma \varepsilon_{b}^{j+1}$ in only finitely many iteration steps - the number of steps needed is, say, $N_{\beta, \alpha}=(\beta-1) /(\alpha-1)$ and, of course, the constant $c_{2, b}$ involved in $\varepsilon_{b}^{j}$ has to be adjusted to a larger one, say, $4^{(\beta-\alpha) \lambda N_{\beta, a}} c_{2, b}$, and the other related constants $c_{1, b}$, $c_{3, b}$ also change accordingly, but all of these constants do not change for $b=a$, which implies that the smallness of $\delta$ required by the induction does not change and therefore, does not depend on $\beta$. This shows that the iteration from $j$-th step to $(j+1)$-th step may be carried out and therefore, Lemma 3.1 is proved.

## 5. Proof of Theorem 1.2

To prove Theorem 1.2, we only need to reexamine Lemma 3.1 and its proof. Under the assumption (1.2) for $\omega$, the corresponding nondegeneracy condition for $\tilde{\omega}$ in Lemma 3.1 is

$$
\begin{equation*}
\gamma \mu|d x| \leq|d \widetilde{\omega}(x)| \leq \gamma|d x|, \quad x \in I_{p}+2 \tag{5.1}
\end{equation*}
$$

with $\mu=\theta \Theta^{-1}$ and $\rho=\gamma \Theta^{-1}$. Accordingly, the equation (3.1) of Lemma 3.1 turns out to be

$$
\begin{equation*}
\left(1-\sum_{k=1}^{j} s_{k}\right) \gamma \mu|d x| \leq|d \tilde{\omega}(x)| \leq\left(1+\sum_{k=1}^{j} s_{k}\right) \gamma|d x|, \quad x \in I_{\rho ; \gamma}^{(j)}+r_{j} \tag{5.2}
\end{equation*}
$$

and all other assumptions and conclusions in Lemma 3.1 remain unchanged. The smallness condition for $h$, of the form (1.11), and therefore, the smallness condition $0<\delta_{a} \leq \delta$ with the corresponding notation

$$
\begin{equation*}
\|f\|_{\beta \lambda+\lambda+\tau ; I_{\rho} \times \mathbf{T}^{n}}=\gamma \mu^{2} \delta_{b}, \quad b=\beta \lambda+\lambda+\tau \geq a=\alpha \lambda+\lambda+\tau \tag{5.3}
\end{equation*}
$$

instead of (2.25), and with $f$ defined from $h$ by (2.3), is needed only for proving the existence of $I_{\rho ; \gamma}^{(j+1)}$ such that $\omega^{(j+1)}\left(I_{\rho ; \gamma}^{(j+1)}\right)=\Omega_{\gamma}$ from $j$-th step to $(j+1)$-th step in induction. But this is no problem because in this case, to apply Lemma 4.4, we need only to let $C=\left\{x^{*}\right\}+\frac{1}{8 n^{2}} \mu r_{j}$ without any other change. From (5.2) and the induction assumption for the first $j$ steps, we have

$$
\begin{equation*}
\frac{2}{3} \gamma \mu|d x| \leq|d F(x)| \leq \frac{4}{3} \gamma|d x|, \quad x \in D \tag{5.4}
\end{equation*}
$$

which implies that $\left|\frac{\partial F}{\partial x}\right|_{D} \leq \frac{4}{3} \gamma$. By Taylor's formula for $F\left(x_{1}\right)-F\left(x_{2}\right)$ up to second order and Cauchy's estimate with the notice that $C+4^{-1} r_{j} \subset D$ and $\left|\frac{\partial^{2} F}{\partial x \partial x}\right|_{C} \leq\left(4^{-1} r_{j}\right)^{-1} \cdot \frac{4}{3} \gamma$, we easily show that, for $x_{1}, x_{2} \in C$,

$$
\begin{equation*}
\left|F\left(x_{1}\right)-F\left(x_{2}\right)\right| \geq \frac{1}{3} \gamma \mu\left|x_{1}-x_{2}\right| . \tag{5.5}
\end{equation*}
$$

As a result, we get $L=\min _{x \in \partial C}\left\{\|F(x)-y\|_{2}\right\} \geq \frac{1}{24 n^{2}} \gamma \mu^{2} r_{j}$ with $y=\omega^{(j)}\left(x^{*}\right)$ as assumed tacitly. Therefore, $\sup _{x \in C}\left\{\|F(x)-G(x)\|_{2}\right\} \leq \sqrt{n}\left|\frac{\partial b_{j}}{\partial x}\right|_{\tilde{\nu}_{j}} \leq \frac{1}{8} \cdot \frac{1}{24 n^{2}} \gamma \mu^{2} r_{j}$, if $\delta$ is small enough but only depending on $n, \tau, \lambda$ and $\alpha$ because the assumption (5.3) will certainly lead to such an estimate for $b_{j}$, which guarantees the existence of $I_{\rho ; \gamma}^{(j+1)}$ as required. The proof of the remainder is standard.

## 6. Application to small twist problem

A direct application of the above theorems gives the existence of invariant tori with a smooth foliation structure (differentiable, $C^{\infty}$-smooth or analytic according to the corresponding smoothness of the considered mapping respectively) of
a nearly integrable symplectic mapping with a small twist. The result may be formulated as follows.

Theorem 6.1. Under the assumptions of Theorem 1.1, consider one parameter family of mappings $S_{t}:(p, q) \rightarrow(\hat{p}, \hat{q})$ with $S_{0}=I$ and $S_{1}=S$, to be defined in phase space $I \times \mathrm{T}^{n}$ by

$$
\left\{\begin{array}{l}
\hat{p}=p-t \frac{\partial H}{\partial q}(\hat{p}, q)=p-t \frac{\partial h}{\partial q}(\hat{p}, q)  \tag{6.1}\\
\hat{q}=q+t \frac{\partial H}{\partial p}(\hat{p}, q)=q+t \omega(\hat{p})+t \frac{\partial h}{\partial p}(\hat{p}, q) .
\end{array}\right.
$$

Under the smallness conditions for $h$ of Theorems 1.1 and 1.2 (in the case when only nondegeneracy condition $(1.2)^{\prime}$ is satisfied by $\omega$ ), the corresponding conclusions of them are still valid for $S_{t}, 0<t \leq 1$, only with the following remarks:

1. $\Omega_{\gamma}$ is replaced by

$$
\begin{equation*}
\Omega_{t, \gamma}=\left\{\omega \in \Omega_{*}:\left|e^{i(k, t \omega)}-1\right| \geq \frac{t \gamma}{|k|^{\tau}} \quad \text { for } k \in \mathbf{Z}^{n} \backslash\{0\}\right\} \tag{6.2}
\end{equation*}
$$

which means to depend on both $\gamma$ and $t$, where $\Omega_{*}$ denotes the set of points in $\Omega$ with distance to its boundary at least equal to $2 \gamma$; and accordingly, $I_{\gamma}$ is replaced by $I_{t, \gamma}$, a closed subset of $I ; \omega_{\gamma}$ replaced by $\omega_{t, \gamma}: I_{t, \gamma} \rightarrow \Omega_{t, \gamma}$, an onto map; $\Phi$ replaced by $\Phi_{t}: I_{t, \gamma} \times \mathbf{T}^{n} \rightarrow \mathbf{R}^{n} \times \mathbf{T}^{n}$ and $R$ replaced by $R_{t}:(\xi, \eta) \rightarrow\left(\xi, \eta+t \omega_{t, \gamma}(\xi)\right)$.
2. If $\Omega$ is a bounded open set of type $D$ in Arnold's sense [1], then we have the following Lebesgue measure estimate

$$
\begin{equation*}
m\left(\Omega \backslash \Omega_{t, \gamma}\right) \leq D \gamma m \Omega \tag{6.3}
\end{equation*}
$$

for $t \in(0,1]$, with constant $D$ only depending on $n, \tau$ and the geometry of $\Omega$. So in this case, $\Omega_{t, \gamma}$ is still a large Cantor set in $\Omega$ if $\gamma$ is small enough.

We conclude the note by remarking that Theorem 6.1 implies the existence of invariant tori with smooth foliation structure and therefore, also implies the existence of $n$ independent smooth invariant functions which are in involution and well-defined on the set filled by the invariant tori in the Whitney's sense, of a symplectic numerical integrator applied to an integrable or a nearly integrable Hamiltonian system if the system is nondegenerate and the time-step size of the integrator is small enough. The invariant tori are just those level sets of the $n$ invariant functions. The nondegeneracy of a nearly integrable system means the nondegeneracy of the integrable part of the system. Symplectic integrators or, in
other words, symplectic algorithms with their computer performance are referred to [4].

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