# Residues of Closed Differential Forms 

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# Residues of Closed Differential Forms 

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November 2, 1996

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#### Abstract

Let $E \rightarrow \bar{M}$ be an oriented real vector bundle of even rank over a smooth compact manifold $\bar{M}$ of dimension $N$, with or without boundary. We show that each section $f$ of $E$ gives rise to a primitive of the Euler form of $E$ away from the zero set of $f$. This leads to the generalized Hopf formula for vector fields on a compact manifold. A very particular case of the Hopf formula is the Gauss-Bonnet theorem on a compact manifold with boundary. The proof is similar in spirit to that of Bott and Chern (Acta Math., 114 (1965)). We indicate how these techniques may be used to highlight the de Rham cohomology of $\bar{M}$ in relation to the homotopy class of the mapping $\Delta: M \times \partial \bar{M} \rightarrow \mathbf{R}^{N} \backslash\{0\}$, where $\Delta$ is a defining mapping for the diagonal of $\bar{M} \times \bar{M}$ and $M$ the "interior" of $\bar{M}$. As a consequence, we derive an explicit formula for residues of closed differential forms in a shell.


AMS subject classification: primary: 55 M 20 ; secondary: $57 \mathrm{R} 20,58 \mathrm{~A} 10$.
Key words and phrases: differential forms, Euler characteristic, Hopf formula.

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## 1 Introduction

The classical Hopf formula says that if $f: M \rightarrow T(M)$ is a vector field on an oriented compact closed manifold $M$ and each critical point of $f$ is isolated, then

$$
\sum_{p \in f^{-1}(0)} \operatorname{deg}(f, p)=\chi(M)
$$

where $\operatorname{deg}(f, p)$ is the local degree of $f$ at $p$ and $\chi(M)$ the Euler characteristic of the manifold $M$ (see for instance Milnor [M, p.37]). Thus, the sum on the left is actually independent of the particular choice of the vector field $f$. We give a short proof of the Hopf formula in the appendix.

For compact manifolds with boundary this formula is no longer true, even if $f$ has no critical points on the boundary.

Example 1.1 Let $M=B_{2} \backslash \overline{B_{1}}$, where $B_{2}$ is a ball with center 0 in $\mathbb{R}^{N}$ and $B_{1} \subset \subset B_{2} \backslash\{0\}$ is a smaller ball. Consider the vector field

$$
\begin{aligned}
f(x) & =\left(\begin{array}{c}
x_{1} \\
\ldots \\
x_{N}
\end{array}\right) \\
& =\frac{1}{2} \nabla\left(x_{1}^{2}+\ldots+x_{N}^{2}\right)
\end{aligned}
$$

in $\bar{M}$. Then, the only critical point of $f$ in $\bar{M}$ is $p=0$ and $\operatorname{deg}(f, 0)=1$, while $\chi(\bar{M})=1+(-1)^{N-1}$.

The aim of this paper is to extend the Hopf formula to smooth vector fields on a compact manifold with boundary.

Our generalization of the Hopf formula relies on the observation that the Euler characteristic of a compact closed manifold is equal to the integral of the Euler form of $M$. We then show that a smooth vector field $f$ on a compact manifold $\bar{M}$ of even dimension gives rise to a primitive of the Euler form of $\bar{M}$ away from the set of critical points of $f$. This enables us to write the integral of the Euler form of $\bar{M}$ as the sum of two terms, the first of the two is the integral of the primitive over the boundary of $\bar{M}$, the second being the sum of local degrees of $f$ at the critical points.

In particular, we may consider as $f$ the gradient of a defining function $\rho$ of the boundary of $\bar{M}$. Then, by the Morse Inequalities, the sum of local degrees of $f$ at the critical points is equal to the Euler characteristic of $\bar{M}$. In this way we obtain what Bott and Chern [BCh] called the relative Gauss-Bonnet Theorem. Our solution falls short of providing an explicit formula for the primitive.

If $M$ is a bounded domain in $\mathbb{R}^{N}$ with $C^{2}$ boundary, then the Euler form of $M$ is zero. The integral over the boundary in the Gauss-Bonnet formula is none other
than the Poincaré index of the unit outward normal vector to the surface $\partial \bar{M}$. With the help of this formula, we prove that the rotation on the boundary of $\bar{M}$ of the unit outward normal vector to $\partial \bar{M}$ is equal to the Euler characteristic of $\bar{M}$. Hence the same remains valid for all smooth vector fields $f$ on $\bar{M}$ whose restrictions to the boundary of $\bar{M}$ are homotopic to the unit outward normal vector of $\partial \bar{M}$.

Starting with the Gauss-Bonnet formula, we introduce a family of double differential forms on $M \times \partial \bar{M}$ which gives rise to an explicit homotopy formula for differential forms on $M$.

If $\bar{M}$ is of the form of "Swiss cheese" $\bar{M}_{0} \backslash\left(M_{1} \cup \ldots \cup M_{N}\right)$, where $\bar{M}_{0} \subset M^{\prime}$ and $\bar{M}_{i} \subset M_{0}$ satisfy appropriate convexity conditions, the homotopy formula yields a residue formula for closed differential forms in $\bar{M}_{0}$ with singularities in the holes $M_{i}, i=1, \ldots, I$.

We briefly sketch also the case of complex-valued vector fields indicating how these techniques result in the Rouchet principle for holomorphic mappings.

## 2 Integration of the Euler form

Let $E \rightarrow M$ be an oriented real vector bundle of rank $2 n$ over a smooth manifold $M$, with or without boundary.

Fix a Euclidean metric $(\cdot, \cdot)$ on $E$ and a connection $\partial: C_{l o c}^{\infty}(E) \rightarrow C_{l o c}^{\infty}\left(E \otimes \Lambda^{1}\right)$ on $E$ preserving the metric, i.e., such that $d(u, v)_{x}=(\partial u, v)_{x}+(u, \partial v)_{x}, x \in M$, for all $u, v \in C_{\text {loc }}^{\infty}(E)$ (see Fedosov [F, p.27] ${ }^{1}$ ).

Denote by $\Omega=\partial \circ \partial$ the curvature of the connection $\partial$; this is a global section of $C_{\text {loc }}^{\infty}\left(\operatorname{Hom}(E, E) \otimes \Lambda^{2}\right)$.

Given any local orthonormal frame $e_{1}, \ldots, e_{2 n}$ for $E$, the curvature is represented by a skew-symmetric matrix

$$
\Omega=\left(\Omega_{i j}\right)_{\substack{i=1, \ldots, 2 n \\ j=1, \ldots, 2 n}}
$$

whose entries are 2-forms (ibid., p. 28). The Euler form of $E$ is the differential form of degree $2 n$ on $M$ defined by

$$
\begin{aligned}
x(E) & =\text { Pfaffian }\left(\frac{1}{2 \pi} \Omega\right) \\
& =\frac{1}{2^{n} n!(2 \pi)^{n}} \sum_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}(-1)^{e_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)} \Omega_{i_{1} j_{1}} \ldots \Omega_{i_{n} j_{n}}},
\end{aligned}
$$

where the sum is taken over all rearrangements $\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)$ of the numbers $(1,2, \ldots, 2 n-1,2 n)$ and $\varepsilon_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}$ is the parity of the rearrangement. In other words, $\chi(E)$ is the coefficient of proportionality in the relation

$$
\frac{1}{n!}\left(\frac{1}{2} \sum_{i, j} \frac{1}{2 \pi} \Omega_{i j} e_{i} \wedge e_{j}\right)^{\wedge n}=\chi(E) e_{1} \wedge \ldots \wedge e_{2 n}
$$

[^1]it is known that it is independent of the choice of the local oriented orthonormal frame for $E$ (ibid., p. 29).

We are interested in evaluating the integral

$$
\begin{equation*}
\int_{M} \chi(E) \tag{2.1}
\end{equation*}
$$

for compact manifolds $M$, with or without boundary (the dimension of $M$ must be equal to $2 n$, since otherwise the integral is zero).

The cohomology class of the Euler form $\chi(E)$ is known to be independent of the particular choice of the Euclidean connection $\partial$ on $E$ (ibid., p. 29). Therefore, if $M$ is closed (i.e., compact and without boundary), then integral (2.1) is a topological invariant of $M$ (i.e., it does not depend on the choice of the connection). It is called the Euler number of the bundle $E$. In particular, if $E=T(M)$, then integral (2.1) is equal to the Euler characteristic of $M$ (the Gauss-Bonnet-Chern Theorem).

We now restrict our attention to the case where $M$ is a compact manifold with boundary. If such is the case, it is convenient to write $\bar{M}$ rather than $M$, thus letting $M$ stand for the "interior" of $\bar{M}$.

Theorem 2.1 Suppose that $E$ possesses a section $f$ which vanishes at no point of $\bar{M}$. Then the Euler form of $E$ is exact, i.e., there is a differential form $\phi(f) \in$ $C^{\infty}\left(\Lambda^{2 n-1} T^{*}(\bar{M})\right)$ such that $d \phi(f)=\chi(E)$ on $\bar{M}$.

Proof. Let

$$
e(x)=\frac{f(x)}{|f(x)|}, \quad x \in \bar{M},
$$

be the section, lying on the unit sphere of $E$, which corresponds to $f$.
Consider the connection $\partial_{1}$ in the bundle $E$, whose value at a section $u \in C_{\text {loc }}^{\infty}(E)$ is

$$
\partial_{1} u(x)=\partial u(x)+e(x)(u, \partial e)_{x}-\partial e(x)(u, e)_{x}, \quad x \in \bar{M}
$$

(it is worth mentioning that any two connections in $E$ differ by a global one-form $\Delta \Gamma \in C_{\text {loc }}^{\infty}\left(\operatorname{Hom}(E, E) \otimes \Lambda^{1}\right)$.)

For each $u, v \in C_{\text {loc }}^{\infty}(E)$, we have

$$
\begin{aligned}
\left(\partial_{1} u, v\right)_{x}+\left(u, \partial_{1} v\right)_{x}= & (\partial u, v)_{x}+(u, \partial e)_{x}(e, v)_{x}-(u, e)_{x}(\partial e, v)_{x} \\
& +(u, \partial v)_{x}+(\partial e, v)_{x}(u, e)_{x}-(e, v)_{x}(u, \partial e)_{x} \\
= & (\partial u, v)_{x}+(u, \partial v)_{x} \\
= & d(u, v)_{x},
\end{aligned}
$$

i.e., $\partial_{1}$ is a Euclidean connection. Moreover,

$$
\begin{aligned}
\partial_{1} e(x) & =\partial e(x)+e(x)(e, \partial e)_{x}-\partial e(x)(e, e)_{x} \\
& =0
\end{aligned}
$$

for $(e, e)_{x} \equiv 1$ and $(e, \partial e)_{x}=\frac{1}{2} d(e, e)_{x}=0$. Hence $\partial_{1}$ vanishes at the section $e$.

Now, we consider the homotopy between $\partial$ and $\partial_{1}$, i.e., the family of connections in the bundle $E$

$$
\begin{align*}
\partial_{t} & =(1-t) \partial+t \partial_{1} \\
& =\partial+t(e(\cdot, \partial e)-\partial e(\cdot, e)) \tag{2.2}
\end{align*}
$$

depending on the parameter $t \in[0,1]$. Analysis similar to the above shows that $\partial_{t}$ is a Euclidean connection in $E$.

Letting $\Omega_{t}=\partial_{t} \circ \partial_{t}$ denote the curvature of the connection, we get, for each $u \in C_{l o c}^{\infty}(E)$,

$$
\begin{align*}
\Omega_{t} u= & \partial \partial u+t\left(\partial e(u, \partial e)+e d(u, \partial e)_{x}-\partial \partial e(u, e)_{x}+\partial e \wedge d(u, e)_{x}\right) \\
& -t e(\partial u, \partial e)_{x}+t^{2} e\left((e, \partial e)_{x} \wedge(u, \partial e)_{x}+(\partial e, \partial e)_{x}(u, e)_{x}\right) \\
& -t \partial e \wedge(\partial u, e)_{x}-t^{2} \partial e \wedge\left((e, e)_{x}(u, \partial e)_{x}+(\partial e, e)(u, e)_{x}\right) \\
= & \Omega u+\left(2 t-t^{2}\right) \partial e \wedge(u, \partial e)_{x}+t\left(e(\Omega e, u)_{x}-\Omega e(e, u)_{x}\right), \tag{2.3}
\end{align*}
$$

since $(e, \partial e)_{x}=0$ and $(\partial e, \partial e)_{x}=0$.
Denote by $\chi_{t}(E)=\operatorname{Pfaffian}\left(\frac{1}{2 \pi} \Omega_{t}\right)$ the corresponding Euler form of the bundle E. As

$$
\begin{aligned}
\Omega_{1} e & =\partial_{1}\left(\partial_{1} e\right) \\
& \equiv 0,
\end{aligned}
$$

we can assert that $\chi_{1}=\operatorname{Pfaffian}\left(\frac{1}{2 \pi} \Omega_{1}\right) \equiv 0$.
Indeed, complete the section $e$ to a local orthonormal frame $e_{1}, \ldots, e_{2 n}$ for $E$, with $e_{1}=e$. In this frame, the curvature $\Omega_{1}$ is represented by a skew-symmetric matrix $\left(\Omega_{i j}\right)$ with $\Omega_{1 j}=\Omega_{i 1} \equiv 0$ for all $i, j=1, \ldots, 2 n$. From this, the desired conclusion follows.

Thus,

$$
\begin{aligned}
\chi(E) & =\chi_{0}(E) \\
& =-\int_{0}^{1}\left(\frac{d}{d t} \chi_{t}(E)\right) d t
\end{aligned}
$$

Our next goal is to find a primitive for $\frac{d}{d t} \chi_{t}(E)$. To this end, we invoke the variation formula for the curvature (cf. Proposition 1.2.4 in Fedosov [F]), according to which $\frac{d}{d t} \Omega_{t}=\partial_{t} \dot{\Gamma}$, where

$$
\begin{equation*}
\dot{\Gamma}=e(\cdot, \partial e)_{x}-\partial e(\cdot, e)_{x} \tag{2.4}
\end{equation*}
$$

Hence it follows that

$$
\begin{align*}
\frac{d}{d t} & \text { Pfaffian }\left(\frac{1}{2 \pi} \Omega_{t}\right) e_{1} \wedge \ldots \wedge e_{2 n} \\
& =\frac{1}{(n-1)!}\left(\frac{1}{2} \sum_{i, j} \frac{1}{2 \pi}\left(\frac{d}{d t} \Omega_{t}\right)_{i j} e_{i} \wedge e_{j}\right) \wedge\left(\frac{1}{2} \sum_{i, j} \frac{1}{2 \pi}\left(\Omega_{t}\right)_{i j} e_{i} \wedge e_{j}\right)^{\wedge(n-1)} \\
& =\frac{1}{(n-1)!}\left(\frac{1}{2} \sum_{i, j} \frac{1}{2 \pi}\left(\partial_{t} \dot{\Gamma}\right)_{i j} e_{i} \wedge e_{j}\right) \wedge\left(\frac{1}{2} \sum_{i, j} \frac{1}{2 \pi}\left(\Omega_{t}\right)_{i j} e_{i} \wedge e_{j}\right)^{\wedge(n-1)} \tag{2.5}
\end{align*}
$$

Note that the equality (2.5) remains valid for each oriented orthonormal frame $e_{1}, \ldots, e_{2 n}$. In particular, for each fixed $(t, x) \in[0,1] \times \bar{M}$, we may choose a proper basis. Namely, given any point $(t, x) \in[0,1] \times \bar{M}$, take $e_{1}, \ldots, e_{2 n}$ such that $\partial_{t} e_{j}=0$ at this point. Then, the covariant derivative $\partial_{t}$ coincides with the exterior derivative $d$ at this point. By Bianchi's identity, we obtain

$$
\begin{aligned}
\left.d \Omega_{t}\right|_{(t, x)} & =\left.\partial_{t} \Omega_{t}\right|_{(t, x)} \\
& =0,
\end{aligned}
$$

so equality (2.5) implies

$$
\begin{aligned}
\frac{d}{d t} & \text { Pfaffian }\left(\frac{1}{2 \pi} \Omega_{t}\right) \\
& =\frac{1}{2^{n}(n-1)!(2 \pi)^{n}} \sum_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}(-1)^{\varepsilon_{\left.i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)} d\left((\dot{\Gamma})_{i_{1} j_{1}}\left(\Omega_{t}\right)_{i_{2} j_{2}} \ldots\left(\Omega_{t}\right)_{i_{n} j_{n}}\right)} .
\end{aligned}
$$

Combining this with the above expression for $\chi(E)$, we derive the desired primitive for the Euler form in the form

$$
\begin{align*}
\phi(f) & =-\int_{0}^{1} \frac{1}{2^{n}(n-1)!(2 \pi)^{n}} \\
& \times \sum_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}(-1)^{\varepsilon_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}}(\dot{\Gamma})_{i_{1} j_{1}}\left(\Omega_{t}\right)_{i_{2} j_{2}} \ldots\left(\Omega_{t}\right)_{i_{n} j_{n}} d t \tag{2.6}
\end{align*}
$$

where $\dot{\Gamma}$ is given by formula (2.4) and $\Omega_{t}$ by formula (2.3). This completes the proof.
It follows from Theorem 2.1 that, if $E$ possesses a section $f$ vanishing at no point of $\bar{M}$, then integral (2.1) can be reduced to an integral over the boundary of $\bar{M}$. Namely, we have, by Stokes' formula,

$$
\begin{align*}
\int_{\bar{M}} \chi(E) & =\int_{\bar{M}} d \phi(f) \\
& =\int_{\partial \tilde{M}} \phi(f) . \tag{2.7}
\end{align*}
$$

## 3 Kronecker formula

In the sequel, we restrict our attention to the case where $E$ is the tangent bundle of $\bar{M}$, i.e., $E=T(\bar{M})$. In this case, $f$ is a vector field on $\bar{M}$ vanishing nowhere in $\bar{M}$.

Consider the particular case where $M$ is a bounded domain in $\mathbb{R}^{2 n}$ with a smooth boundary $\partial \bar{M}$ (possibly consisting of several connected components).

Identifying $T(\bar{M})$ with $\bar{M} \times \mathbb{R}^{n}$ and taking $\partial=d$ as the original connection on $E$, we see at once that

$$
\begin{aligned}
\Omega_{t} & =\left(2 t-t^{2}\right) d e \wedge(\cdot, d e)_{x}, \\
\dot{\Gamma} & =e(\cdot, d e)_{x}-d e(\cdot, e)_{x} .
\end{aligned}
$$

Writing $f(x)=\sum_{j=1}^{2 n} f_{j}(x) \partial / \partial x_{j}$, with $f_{j}$ a $C^{\infty}$ function on $\bar{M}$, we get

$$
\begin{aligned}
\left(\Omega_{t}\right)_{i j} & =\left(2 t-t^{2}\right) d e_{i} \wedge d e_{j}, \\
(\tilde{\Gamma})_{i j} & =e_{i} d e_{j}-e_{j} d e_{i}
\end{aligned}
$$

where $e_{j}(x)=\frac{f_{j}(x)}{|f(x)|}, j=1, \ldots, 2 n$. Thus, the primitive $\phi(f)$ becomes

$$
\begin{aligned}
\phi(f)= & -\frac{1}{2^{n-1}(n-1)!(2 \pi)^{n}} \int_{0}^{1}\left(2 t-t^{2}\right)^{n-1} d t \\
& \times \sum_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}(-1)^{\varepsilon_{\left(i_{1}, j_{1}, \ldots, i_{n}, j_{n}\right)}} e_{i_{1}} d e_{j_{1}} \wedge d e_{i_{2}} \wedge d e_{j_{2}} \wedge \ldots \wedge d e_{i_{n}} \wedge d e_{j_{n}} \\
= & \frac{(-1)}{\sigma_{2 n}} \sum_{j=1}^{2 n}(-1)^{j-1} e_{j} d e[j],
\end{aligned}
$$

where $\sigma_{N}$ is the area of the unit sphere in $\mathbb{R}^{N}$ and de $[j]$ is the exterior product of the differentials $d e_{1}, \ldots, d e_{2 n}$ one after another except $d e_{j}$.

Since $\Omega \equiv 0$, it follows that $\chi(T(\bar{M}) \equiv 0$. Hence, formula (2.7) means that

$$
\begin{equation*}
\int_{\partial \bar{M}} \phi(f)=0 \tag{3.1}
\end{equation*}
$$

as a matter of course.
Equality (3.1) can be easily extended to the case where $f$ is a vector field on $\bar{M}$ with isolated critical points (if such is the case, the number of critical points is finite). For this purpose, we cut out small neighborhoods $B(p, \epsilon)$ of the points $p \in f^{-1}(0)$ and apply formula (3.1) to the domain $M \backslash \cup_{p \in f^{-1}(0)} \overline{B(p, \epsilon)}$. This yields

$$
\begin{equation*}
\int_{\partial \bar{M} \backslash \cup_{p \in f^{-1}(0)} B(p, c)} \phi(f)=\sum_{p \in f^{-1}(0)} \int_{\bar{M} \cap \partial B(p, c)} \phi(f) \tag{3.2}
\end{equation*}
$$

for each $\epsilon>0$ small enough.
Corollary 3.1 Suppose $M$ is a bounded domain with smooth boundary in $\mathbb{R}^{2 n}$ and $f$ is a smooth vector field on $\bar{M}$ whose critical points are isolated and do not meet $\partial \bar{M}$. Then,

$$
\begin{equation*}
\int_{\partial \dot{M}} \phi(f)=-\sum_{p \in f^{-1}(0)} \operatorname{deg}(f, p) \tag{3.3}
\end{equation*}
$$

Proof. Indeed, if $p \in M$ is an "interior" critical point of $f$, then we may choose $\epsilon>0$ such that the ball $B(p, \epsilon)$ lies entirely in $M$ and contains no critical points of $f$ different from $p$. Then, by the Poincaré formula,

$$
\begin{aligned}
\int_{\bar{M} \cap \partial B(p, c)} \phi(f) & =-\int_{\partial B(p, c)} \frac{1}{\sigma_{2 n}} \sum_{j=1}^{2 n}(-1)^{j-1} \frac{f_{j}}{|f|} d \frac{f}{|f|}[j] \\
& =-\operatorname{deg}(f, p)
\end{aligned}
$$

(cf. for instance Example 6.1.11 in Tarkhanov [T2]). Thus, if $f$ has no critical points on the boundary of $\vec{M}$, then equality (3.2) reduces to (3.3), as required.

The equality (3.3) is known as the Kronecker formula (ibid., Remark 6.1.8). We emphasize that it remains valid also for domains in an odd-dimensional space $\mathbb{R}^{N}$.

## 4 The generalized Hopf formula

We are now in a position to extend the Hopf formula to the case of compact manifolds with boundary.

Theorem 4.1 Suppose $f$ is a vector field of class $C^{2}$ on an oriented compact manifold $\bar{M}$ (with or without boundary), whose critical points are isolated and do not meet the boundary. Then,

$$
\begin{equation*}
\int_{\bar{M}} \chi(T(\bar{M}))=\int_{\partial \bar{M}} \phi(f)+\sum_{p \in f^{-1}(0)} \operatorname{deg}(f, p) . \tag{4.1}
\end{equation*}
$$

Proof. Choose an $\epsilon_{0}>0$ such that the balls $B\left(p, \epsilon_{0}\right), p \in f^{-1}(0)$, lie in the interior of $\bar{M}$ and do not meet each other. Applying formula ( 2.7 to the manifold with boundary $\bar{M} \backslash \cup_{p \in f^{-1}(0)} B(p, \epsilon)$, we obtain

$$
\int_{\bar{M} \backslash \cup_{p \in f^{-1}(0)} B(p, \epsilon)} \chi(T(\bar{M}))=\int_{\partial \bar{M}} \phi(f)-\sum_{p \in f^{-1}(0)} \int_{\partial B(p, c)} \phi(f),
$$

for all $\epsilon<\epsilon_{0}$.
We are going to pass to the limit in both sides of this equality, when $\epsilon \rightarrow 0$. Since the Euler form $\chi(T(\bar{M}))$ is smooth on the whole manifold $\bar{M}$, the limit of the left-hand side does exist and is equal to $\int_{\bar{M}} \chi(T(\bar{M}))$.

Moreover, each summand $\int_{\partial B(p, c)} \phi(f)$ has a limit, when $\epsilon \rightarrow 0$. Indeed, if $0<\epsilon^{\prime}<\epsilon^{\prime \prime}<\epsilon_{0}$, then we get, by Stokes' formula,

$$
\begin{aligned}
\left|\int_{\partial B\left(p, c^{\prime \prime}\right)} \phi(f)-\int_{\partial B\left(p, c^{\prime}\right)} \phi(f)\right| & =\left|\int_{\partial\left(B\left(p, \epsilon^{\prime \prime}\right) \backslash B\left(p, c^{\prime}\right)\right)} \phi(f)\right| \\
& =\left|\int_{B\left(p, \epsilon^{\prime \prime}\right) \backslash B\left(p, c^{\prime}\right)} \chi(T(\bar{M}))\right| \\
& \leq \sup |\chi(T(\bar{M}))| \operatorname{meas}\left(B\left(p, \epsilon^{\prime \prime}\right)\right) .
\end{aligned}
$$

Hence it follows that $\left(\int_{\partial B(p, c)} \phi(f)\right)_{\epsilon<\varepsilon_{0}}$ is a Cauchy net, as $\epsilon \rightarrow 0$.
The same argument shows that the $\operatorname{limit}^{\lim _{c \rightarrow 0} \int_{\partial B(p, c)} \phi(f) \text { is actually indepen- }}$ dent on the particular choice of the base $(B(p, \epsilon))_{c<\epsilon_{0}}$ for the neighborhood system of the point $p$. In particular, we may require $B(p, \epsilon)$ to be the ball in local coordinates at $p$.

From what has already been proved, it follows that

$$
\begin{equation*}
\int_{\bar{M}} \chi(T(\bar{M}))=\int_{\partial \bar{M}} \phi(f)-\sum_{p \in f^{-1}(0)} \lim _{\epsilon \rightarrow 0} \int_{\partial B(p, c)} \phi(f) \tag{4.2}
\end{equation*}
$$

We are thus left with the task of identifying the individual limits in the right-hand side of (4.2).

To this end, we recall that the definition of the form $\phi(f)$ includes a Euclidean connection $\partial$ on the tangent bundle $T(\bar{M})$ (cf. $\Omega=\partial^{2}$ in (2.3)). In order to evaluate the limit

$$
\begin{equation*}
\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, c)} \phi(f) \tag{4.3}
\end{equation*}
$$

we only need to know the connection $\partial$ in a neighborhood of $p$. We claim that the particular choice of this local connection does not affect the limit (4.3). As each local Euclidean connection can be extended to a global one, what we have to show is that, if $\partial_{1}$ and $\partial_{2}$ are two connections on $T(\bar{M})$, which agree close to the boundary of $\bar{M}$, and $\phi_{1}$ and $\phi_{2}$ are the corresponding differential forms (2.6), then

$$
\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, \epsilon)}\left(\phi_{2}-\phi_{1}\right)=0
$$

Indeed, since $\partial_{1}=\partial_{2}$ close to the boundary of $\bar{M}$, we deduce from (2.6), (2.4) and (2.3) that $\phi_{1}=\phi_{2}$ near $\partial \bar{M}$. Moreover, letting $\chi_{1}$ and $\chi_{2}$ denote the corresponding Euler forms of $T(\bar{M})$, we have $\chi_{2}-\chi_{1}=d \phi$, where $\phi$ is a smooth differential form of degree $n-1$ on $\bar{M}$, vanishing close to the boundary of $\bar{M}$. Let us fix a smooth function $\chi$ with a compact support in the ball $B\left(p, \epsilon_{0}\right)$, such that $\chi=1$ in a neighborhood of the point $p$. Then,

$$
\begin{aligned}
\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, \epsilon)}\left(\phi_{2}-\phi_{1}\right) & =\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, \varepsilon)} \chi\left(\phi_{2}-\phi_{1}\right) \\
& =-\lim _{\epsilon \rightarrow 0} \int_{\bar{M} \backslash B(p, c)} d \chi \wedge\left(\phi_{2}-\phi_{1}\right)+\chi d\left(\phi_{2}-\phi_{1}\right) \\
& =-\int_{\bar{M}} d \chi \wedge\left(\phi_{2}-\phi_{1}\right)+\chi\left(\chi_{2}-\chi_{1}\right)
\end{aligned}
$$

for $d \chi$ vanishes near the singular point $p$.
Since $\operatorname{supp} \chi \subset \subset M$, we get, by Stokes' formula,

$$
\begin{aligned}
\int_{\tilde{M}} \chi\left(\chi_{2}-\chi_{1}\right) & =\int_{\tilde{M}} \chi d \phi \\
& =-\int_{\tilde{M}} d \chi \wedge \phi
\end{aligned}
$$

whence

$$
\begin{aligned}
\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, c)}\left(\phi_{2}-\phi_{1}\right) & =-\int_{\bar{M}} d \chi \wedge\left(\phi_{2}-\phi_{1}-\phi\right) \\
& =-\int_{\bar{M}} d(\chi-1) \wedge\left(\phi_{2}-\phi_{1}-\phi\right) \\
& =-\int_{\partial \bar{M}}(\chi-1)\left(\phi_{2}-\phi_{1}-\phi\right)+\int_{\bar{M}}(\chi-1) \wedge d\left(\phi_{2}-\phi_{1}-\phi\right) \\
& =0
\end{aligned}
$$

because $\left(\phi_{2}-\phi_{1}-\phi\right)=0$ on $\partial \bar{M}$ and $d\left(\phi_{2}-\phi_{1}-\phi\right)=0$ on $\bar{M}$. This is the desired conclusion.

Finally, we can assume, by decreasing $\epsilon_{0}$ if necessary, that the restriction of the bundle $T(\bar{M})$ to $B\left(p, \epsilon_{0}\right)$ is isomorphic to $B\left(p, \epsilon_{0}\right) \times \mathbb{R}^{n}$. Consider the local connection $\partial$ in $T(\bar{M})$ given by the exterior derivative in $B\left(p, \epsilon_{0}\right)$. Then, as in Section 3, the form $\phi(f)$ in $B\left(p, \epsilon_{0}\right)$ becomes $\frac{(-1)}{\sigma_{2 n}} \sum_{j=1}^{2 n}(-1)^{j-1} e_{j} d e[j]$, whence

$$
\begin{aligned}
\lim _{\epsilon \rightarrow 0} \int_{\partial B(p, c)} \phi(f) & =-\lim _{c \rightarrow 0} \int_{\partial B(p, c)} \frac{1}{\sigma_{2 n}} \sum_{j=1}^{2 n}(-1)^{j-1} \frac{f_{j}}{|f|} d \frac{f}{|f|}[j] \\
& =-\operatorname{deg}(f, p)
\end{aligned}
$$

which gives (4.1) when substituted in (4.2).
Formula (4.1) contains the Hopf formula as a very particular case. Indeed, if $M$ is closed, then the integral over the boundary in the right-hand side of (4.1) disappears and we deduce that the integral of the Euler form is equal to the algebraic sum of the local degrees of the vector field $f$ at its critical points.

## 5 The relative Gauss-Bonnet Theorem

Let $\bar{M}$ be an oriented compact manifold with boundary, of dimension $N=2 n$ and let $\bar{M}$ be embedded in a differentiable manifold $M^{\prime}$ without boundary, of the same dimension.

We denote by $M$ the interior of $\bar{M}$ and by $\partial \bar{M}$ the boundary of $\bar{M}$. We assume that $\partial \bar{M}$ is a submanifold of codimension 1 and class $C^{2}$ in $M^{\prime}$ (hypersurface).

There exists a real-valued function $\rho \in C_{l o c}^{2}\left(M^{\prime}\right)$ such that $M=\left\{x \in M^{\prime}\right.$ : $\rho(x)<0\}$ and $d \rho(x) \neq 0$ for $x \in \partial \bar{M}$.

If $\rho_{1}$ and $\rho_{2}$ are two functions with these properties, then there is a positive function $k \in C_{l o c}^{1}\left(M^{\prime}\right)$ such that $\rho_{2}=k \rho_{1}$ on $M^{\prime}$. In this way we obtain what is referred to as the defining function of the oriented hypersurface $\partial \bar{M}$.

By a stationary point of $\rho$ is meant a point $x^{0} \in M^{\prime}$ such that $d \rho\left(x^{0}\right)=0$. Recall that a number $v \in \mathbb{R}$ is said to be a critical value of $\rho$ if the preimage $\rho^{-1}(v)$ contains no stationary points of $\rho$.

Lemma 5.1 Let $[a, b], a<b$, be a closed interval containing no critical values of $\rho$ and let $\rho^{-1}[a, b]$ be compact. Then, the hypersurface $\rho^{-1}(a)$ is $C^{2}$ diffeomorphic to $\rho^{-1}(b)$ and the manifold with boundary $\rho^{-1}(-\infty, a]$ is $C^{2}$ diffeomorphic to $\rho^{-1}(-\infty, b]$.

Proof. The idea of the proof is to push $\rho^{-1}(b)$ down to $\rho^{-1}(a)$ along the orthogonal trajectories of the hypersurfaces $\rho(x)=$ constant. See Milnor [M, p.12].

Choose a Riemannian metric on $M^{\prime}$ (i.e., an inner product in the fibers $T_{x}\left(M^{\prime}\right)$ of the tangent bundle); and let $\left(f^{\prime}, f^{\prime \prime}\right)_{x}$ denote the inner product of two tangent vectors, as determined by this metric. The gradient of $\rho$ is the vector.field $\nabla \rho$ on $M^{\prime}$ which is characterized by the identity $(f, \nabla \rho)_{x}=f(\rho)$ ( $=$ directional derivative of $\rho$ along $f$ ) for any vector field $f$. In classical notation, in terms of local coordinates $x=\left(x_{1}, \ldots, x_{N}\right)$, the gradient has components $\sum_{j=1}^{N} g^{i j} \frac{\partial \rho}{\partial x_{j}}$, where $\left(g^{i j}\right)$ is the inverse to the matrix ( $g_{i j}$ ) of the Riemannian metric. This vector field $\nabla \rho$ vanishes precisely at the stationary points of $\rho$.

When applied to the vector field $f=\nabla \rho$ on $\bar{M}$, Theorem 4.1 yields the following result.

Theorem 5.2 Let $\bar{M} \subset M^{\prime}$ be an oriented compact manifold with $C^{2}$ boundary and let $\rho \in C_{l o c}^{2}\left(M^{\prime}\right)$ be a defining function of $\partial \bar{M}$. Then,

$$
\begin{equation*}
\int_{\bar{M}} \chi(T(\bar{M}))=\int_{\partial \bar{M}} \phi(\nabla \rho)+\chi(\bar{M}) . \tag{5.1}
\end{equation*}
$$

If $\bar{M}$ is a compact closed manifold, then (5.1) becomes $\int_{\bar{M}} \chi(T(\bar{M}))=\chi(\bar{M})$ which is the content of the Gauss-Bonnet-Chern Theorem, as above. For this reason, we call Theorem 5.2 the relative Gauss-Bonnet Theorem.

To prove Theorem 5.2 we recall several results of the Morse theory. The best general reference here is the book of Milnor [M].

By a Morse function on $M^{\prime}$ is meant a smooth function $u$ all of whose stationary points are non-degenerate, i.e., the Hessian $\left(\frac{\partial^{2} u}{\partial x_{i} \partial x_{j}}\right)$ is non-singular at those points. It can be checked directly that non-degeneracy does not depend on the coordinate system.

According to the Morse Lemma (cf. ibid., p. 6), if $p$ is a non-degenerate stationary point of $u$, then there are local coordinates $x=\left(x_{1}, \ldots, x_{N}\right)$ in some neighborhood of $p$, in terms of which the function $u$ is given in that neighborhood by $u(x)=u(p)-x_{1}^{2}-\ldots-x_{i}^{2}+x_{i+1}^{2}+\ldots+x_{N}^{2}$.

The set of Morse functions is known to be dense in the space of all smooth functions on $M^{\prime}$ in the topology of uniform convergence along with a finite number of derivatives on compact subsets of $M^{\prime}$.
,The number $i=$ index $(u, p)$ is called the index of the stationary point; obviously, it is independent of the local coordinates. If by pushing $\rho^{1}(b)$ down to $\rho^{-1}(a)$ we cross a stationary point of $\rho$ of index $i$, then the set $\rho^{-1}(-\infty, b]$ has the homolopy type of $\rho^{-1}(-\infty, a]$ with an $i$-cell attached (ibid., p. 14).

There is a close connection between the number of stationary points of a function $u$ on a compact closed manifold $M^{\prime}$ and certain topological invariants of the manifold. In particular, the number

$$
\sum_{i \geq 0}(-1)^{i} \#\left\{p \in \nabla u^{-1}(0): \text { index }(u, p)=i\right\}
$$

is actually independent of the function $u$, coinciding as it does with the Euler characteristic of $M^{\prime}$ (Weak Morse Inequalities, cf. ibid., p. 29).

We are now in a position to deduce Theorem 5.2 as a consequence of Theorem 4.1 and the Morse theory.

Proof. We first observe that the equality (5.1) depends only on the restriction of the function $\rho$ to an infinitesimal neighborhood of the boundary of $\bar{M}$. By assumption, $\rho$ has no stationary points near $\partial \bar{M}$. Therefore, we can assume without loss of generality that $\rho$ is a Morse function on $M^{\prime}$, for if not, we correct $\rho$ away from a neighborhood of $\partial \bar{M}$ in $M^{\prime}$.

If $\rho$ is a Morse function on $M^{\prime}$, then the critical points of the vector field $\nabla \rho$ are non-degenerate and do not meet $\partial \bar{M}$. Hence it follows, by Theorem 4.1, that

$$
\begin{equation*}
\int_{\bar{M}} \chi(T(\bar{M}))=\int_{\partial \bar{M}} \phi(\nabla \rho)+\sum_{p \in \nabla_{\rho}-1}(0) \cap M, \tag{5.2}
\end{equation*}
$$

If $\nabla \rho(p)=0$, then $p$ is a stationary point of the function $\rho$. Let $i$ be the index of $p$. Then there are local coordinates $x=\left(x_{1}, \ldots, x_{N}\right)$ with center at $p$, such that $\rho(x)=\rho(p)-x_{1}^{2}-\ldots-x_{i}^{2}+x_{i+1}^{2}+\ldots+x_{N}^{2}$.

To evaluate the local degree of $\nabla \rho$ at $p$, we invoke formula (3.3) according to which

$$
\operatorname{deg}(\nabla \rho, p)=\int_{|x|=c} \frac{1}{\sigma_{N}} \sum_{j=1}^{N}(-1)^{j-1} \frac{\frac{\partial \rho}{\partial x_{j}}}{|\nabla \rho|} d \frac{\nabla \rho}{|\nabla \rho|}[j],
$$

where $\epsilon$ is small enough. Thus,

$$
\begin{aligned}
\operatorname{deg}(\nabla \rho, p) & =(-1)^{i} \int_{|x|=c} \frac{1}{\sigma_{N}} \sum_{j=1}^{N}(-1)^{j-1} \frac{x_{j}}{|x|} d \frac{x}{|x|}[j] \\
& =(-1)^{i}
\end{aligned}
$$

Hence

$$
\begin{aligned}
\sum_{p \in \nabla \rho^{-1}(0)} \operatorname{deg}(\nabla \rho, p) & =\sum_{i \geq 0}(-1)^{i} \#\left\{p \in \nabla \rho^{-1}(0) \cap M: \operatorname{index}(\rho, p)=i\right\} \\
& =\sum_{i \geq 0}(-1)^{i} \#\{i \text {-cells in } \bar{M}\} \\
& =\chi(\bar{M})
\end{aligned}
$$

the last equality being a consequence of Theorem 3.2 in Milnor [ $\mathrm{M}, \mathrm{p} .14$ ].
When substituted into (5.2), this gives (5.1), which completes the proof.

## 6 Poincare index

In this section we briefly discuss a particular case of the relative Gauss-Bonnet theorem, when $M$ is a bounded domain in $\mathbb{R}^{N}$ with a boundary of class $C^{2}$. We need not assume that $N$ is even.

Fix a defining function $\rho \in C_{l o c}^{2}\left(\mathbb{R}^{N}\right)$ of the surface $\partial \bar{M}$, so that $M=\left\{x \in \mathbb{R}^{N}\right.$ : $\rho(x)<0\}$ and $\left.d \rho\right|_{\partial \bar{M}} \neq 0$. As described in Section 3, we consider the differential form

$$
\phi(\nabla \rho)=\frac{(-1)}{\sigma_{N}} \sum_{j=1}^{N}(-1)^{j-1} \frac{\frac{\partial \rho}{\partial x_{j}}}{|\nabla \rho|} d \frac{\nabla \rho}{|\nabla \rho|}[j] .
$$

The integral $-\int_{\partial \bar{M}} \phi(\nabla \rho)$ is called the rotation of the vector field $\nabla \rho$ on the boundary of $\bar{M}$. This is known also as the Poincaré index of the cycle $(\nabla \rho)_{\sharp} \partial \bar{M}$ with respect to the point 0 . A simple verification shows that the integral is actually independent of the particular choice of the defining function $\rho$ for $\partial \bar{M}$. Thus, it can be referred to as the rotation of the unit outward normal vector on $\partial \bar{M}$.

Corollary 6.1 For each domain $M \subset \subset \mathbb{R}^{N}$ with $C^{2}$ boundary, the rotation of the unit outward normal vector on $\partial \bar{M}$ is equal to the Euler characteristic of $\bar{M}$.

Proof. The assertion of the corollary amounts to the fact that

$$
-\int_{\partial \bar{M}} \phi(\nabla \rho)=\chi(\bar{M}) .
$$

For even $N$, this follows from (5.1) because $\chi(T(\bar{M}))=0$. For odd $N$, the proof runs just as in Theorem 5.2, with Theorem 4.1 replaced by Corollary 3.1.

## 7 Double differential forms $K_{q}(F)$

Given $N$-dimensional vectors $v_{1}, \ldots, v_{\nu}$ with entries in an algebra, we denote by $D_{N_{1}, \ldots, N_{\nu}}\left(v_{1}, \ldots, v_{\nu}\right)$ the determinant of order $N$ whose first $N_{1}$ columns are $v_{1}$, the second $N_{2}$ columns are $v_{2}$ etc., the last $N_{\nu}$ columns are $v_{\nu}$, where $N_{1}+\ldots+N_{\nu}=N$. The determinant is calculated by columns, i.e., $\operatorname{det}\left(v_{i j}\right)=\sum_{I}(-1)^{\varepsilon_{i}} v_{i_{1} 1} \ldots v_{i_{N} N}$ where $\varepsilon_{I}$ denotes the parity of the rearrangement $I=\left(i_{1}, \ldots, i_{N}\right)$ of the integers $(1, \ldots, N)$.

For a smooth mapping $F=F(x, y, t)$ of an open set $\Omega \subset M^{\prime} \times M^{\prime} \times \mathbb{R}^{1}$ to $\mathbb{R}^{N}$, we consider the double differential forms $K_{q}(F), 1 \leq q \leq N$, of degree $q-1$ in $x$ and of degree $N-q$ in $y$ and $t$, given by

$$
K_{q}(F)=\frac{(-1)^{q}}{\sigma_{N}(q-1)!(N-q)!} D_{1, q-1, N-q}\left(\left(\begin{array}{c}
\frac{F_{1}}{|F|}  \tag{7.1}\\
\dddot{F_{N}} \\
\frac{F F \mid}{|F|}
\end{array}\right),\left(\begin{array}{c}
d_{x} \frac{F_{1}}{\ldots \mid} \\
\cdots \\
d_{x} \frac{F_{N}}{|F|}
\end{array}\right),\left(\begin{array}{c}
\left(d_{y}+d_{t}\right) \frac{F_{1}}{|F|} \\
\cdots \\
\left(d_{y}+d_{t}\right) \frac{F_{N}}{F \mid}
\end{array}\right)\right) .
$$

Moreover, we set $K_{0}=K_{N+1}=0$.
It is worth mentioning that these double forms were first introduced by the third author in [T1] (see also Tarkhanov [T2, 6.1.8]).

Lemma 7.1 The form $K_{q}(F(x, y, t))$ is the component of degree $q-1$ in $x$ of the "pull-back" $F^{\sharp} K_{1}(z)$ times $(-1)^{(N-1)(q-1)}$ where we first have to take full differentials in $x, y$ and $t$, then to place all the differentials in $x$ after the differentials in $y$ and $t$, and finally to declare the form to be a double one.

Proof. The proof is straightforward.
Recall that two continuous mappings $f_{0}, f_{1}: T_{1} \rightarrow T_{2}$ of topological spaces are said to be homotopic if there exists a continuous mapping $f_{t}: T_{1} \times[0,1] \rightarrow T_{2}$ which coincides with $f_{0}$ for $t=0$ and with $f_{1}$ for $t=1$.

If $f_{0}$ and $f_{1}$ are homotopic and differentiable, then one can choose the homotopy $f_{t}$ to be differentiable too, roughly speaking, of the same class as $f_{0}$ and $f_{1}$ (see Theorem 8 of Pontryagin [ $\mathrm{P}, \mathrm{p} .64]$ ).

From the topological point of view, if smooth mappings $f_{0}, f_{1}: \partial \bar{M} \rightarrow \mathbb{R}^{N} \backslash\{0\}$ are homotopic, then they have the same rotations on $\partial \bar{M}$ (see Krasnosel'skii and Zabreiko [KZ, p.16]).

From the analytical point of view, it follows from the Poincare formula that $\int_{\partial \bar{M}} f_{1}^{\sharp} K_{1}(z)-f_{0}^{\sharp} K_{1}(z)=0$. Hence the difference $f_{1}^{\sharp} K_{1}(z)-f_{0}^{\sharp} K_{1}(z)$ is, by the de Rham Theorem, exact on $\partial \bar{M}$.

Our next objective is to extend this obvious observation to the double forms $K_{q}(F(x, y, t))$.

Lemma 7.2 Let $F=F(x, y, t)$ be a mapping of an open set $\Omega \subset M^{\prime} \times M^{\prime} \times \mathbb{R}^{1}$ to $\mathbb{R}^{N}$, of class $C^{2}$. Then the following equality holds on the set $\Omega \backslash F^{-1}(0)$ :

$$
\begin{equation*}
d_{x} K_{q}(F)+(-1)^{q+1}\left(d_{y}+d_{t}\right) K_{q+1}(F)=0 . \tag{7.2}
\end{equation*}
$$

Proof. Using Lemma 7.1 and the equality $d K_{1}(z)=-\delta(z) d z$ in $\mathbb{R}^{N 2}$, we obtain

$$
\begin{aligned}
d_{x} & K_{q}(F)+(-1)^{q+1}\left(d_{y}+d_{t}\right) K_{q+1}(F) \\
& =(-1)^{N-q} d_{x}\left((-1)^{(N-1)(q-1)} F^{\sharp} K_{1}(z)\right)+(-1)^{q+1}\left(d_{y}+d_{t}\right)\left((-1)^{(N-1) q} F^{\sharp} K_{1}(z)\right) \\
& =(-1)^{N q+1} d F^{\sharp} K_{1}(z) \\
& =(-1)^{N q} F^{\sharp} \delta(z) d z \\
& =0,
\end{aligned}
$$

for $(x, y, t) \in \Omega \backslash F^{-1}(0)$, as desired.
For another proof of Lemma 7.2, see the proof of Lemma 1.2 in the book of Aizenberg and Dautov [AD].

Lemma 7.3 If $f_{0}(x, y), f_{1}(x, y)$ are homotopic mappings of the set $M \times \partial \bar{M}$ to $\mathbb{R}^{N} \backslash\{0\}$, of class $C^{2}$, and $f_{t}(x, y), t \in[0,1]$, is a $C^{2}$ homotopy between them, then, for every $0 \leq q \leq N-1$, we have on $M \times \partial \bar{M}$ :

$$
\begin{align*}
& K_{q+1}\left(f_{1}\right)-K_{q+1}\left(f_{0}\right) \\
& \left.\left.\quad=d_{x}\left((-1)^{q} \int_{0}^{1} d t\right\rfloor K_{q}\left(f_{t}\right) d t\right)-(-1)^{q} d_{y}\left((-1)^{q+1} \int_{0}^{1} d t\right\rfloor K_{q+1}\left(f_{t}\right) d t\right) . \tag{7.3}
\end{align*}
$$

Proof. It suffices to integrate equality (7.2) over $t \in[0,1]$ and to take into account that

$$
\left.\left.d_{y} \int_{0}^{1} d t\right\rfloor K_{q+1}\left(f_{t}\right) d t=-\int_{0}^{1} d t\right\rfloor d_{y} K_{q+1}\left(f_{t}\right) d t
$$

## 8 Homotopy formula

Denote by [diagonal $\left.\left(M^{\prime} \times M^{\prime}\right)\right]$ the current of integration over the diagonal submanifold of $M^{\prime} \times M^{\prime}$.

Example 8.1 If $M^{\prime}=\mathbb{R}^{N}$, then

$$
\left[\text { diagonal }\left(M^{\prime} \times M^{\prime}\right)\right]=(-1)^{N} \Delta^{\sharp} \delta(z) d z
$$

where $z=\Delta(x, y)$ is the mapping of $\mathbb{R}^{N} \times \mathbb{R}^{N} \rightarrow \mathbb{R}^{N}$, given by $\Delta(x, y)=y-x$.

[^2]We call a $C^{2}$ mapping $z=\Delta(x, y)$ of $M^{\prime} \times M^{\prime} \rightarrow \mathbb{R}^{N}$ defining for the diagonal of $M^{\prime} \times M^{\prime}$, if $\Delta^{\sharp} \delta(z) d z=(-1)^{N}$ [diagonal $\left.\left(M^{\prime} \times M^{\prime}\right)\right]$.

For a mapping $\Delta: M^{\prime} \times M^{\prime} \rightarrow \mathbb{R}^{N}$, being defining is equivalent to the following conditions:

- diagonal $\left(M^{\prime} \times M^{\prime}\right)=\left\{(x, y) \in M^{\prime} \times M^{\prime}: \Delta(x, y)=0\right\} ;$
- $\operatorname{det} \frac{\partial \Delta}{\partial y}(x, y)>0$ for each $(x, y) \in \Delta,-$
as well as a number of similar conditions on other maximal minors of the Jacobi matrix $\frac{\partial \Delta(x, y)}{\partial(x, y)}$.

Example 8.2 Let $M^{\prime}$ be an open set in $\mathbb{R}^{N}$ and

$$
\Delta(x, y)=\left(\begin{array}{c}
\phi_{1}(x, y)\left(y_{1}-x_{1}\right) \\
\cdots \\
\ldots \\
\phi_{N}(x, y)\left(y_{N}-x_{N}\right)
\end{array}\right)
$$

where $\phi_{j} \in C_{l o c}^{2}\left(M^{\prime} \times M^{\prime}\right), j=1, \ldots, N$. Then,

$$
\operatorname{det} \frac{\partial \Delta}{\partial y}(x, x)=\phi_{1}(x, x) \ldots \phi_{N}(x, x)
$$

so $\Delta$ is a defining mapping for the diagonal of $M^{\prime} \times M^{\prime}$ if and only if

$$
\phi_{1}(x, x) \ldots \phi_{N}(x, x)>0 \quad \text { for all } \quad x \in M^{\prime} .
$$

When working with holomorphic mappings of Stein manifolds, we are not able to ensure the existence of such a mapping in general. Moreover, if $M^{\prime}$ is not parallelizable, one can deduce from the results of Schneider [ $S]$ that it is even impossible to satisfy only the second condition.

To obtain a defining mapping for the diagonal of $M^{\prime} \times M^{\prime}$ in case $M^{\prime}$ is parallelizable, we may use the following simple construction (cf. Khenkin and Leiterer [KhL, 4.2.4]).

Lemma 8.3 There exists a smooth mapping $\Delta: M^{\prime} \times M^{\prime} \rightarrow T\left(M^{\prime}\right)$ such that the following conditions are fulfilled:

1) $\Delta(x, y) \in T_{x}\left(M^{\prime}\right)$ for all $(x, y) \in M^{\prime} \times M^{\prime}$, i.e., $\Delta$ is a section of the pull-back of the bundle $T\left(M^{\prime}\right)$ by the mapping $(x, y) \mapsto x$ of $M^{\prime} \times M^{\prime} \rightarrow M^{\prime}$; and
2) for every fixed $x \in M^{\prime}$, we have $\Delta(x, x)=0$ and the mapping $\Delta(x, y)$ from $M^{\prime}$ to $T_{x}\left(M^{\prime}\right)$ is a diffeomorphism in some neighborhood of $y=x$.

Proof. According to the Whitney Theorem, $M^{\prime}$ can be embedded as a submanifold into a space $\mathbb{R}^{N^{\prime}}$ for some large but finite $N^{\prime}$. In other words, there exists a smooth one-to-one mapping $f: M^{\prime} \rightarrow \mathbb{R}^{N^{\prime}}$ such that, for all $x \in M^{\prime}$, the Jacobi matrix of $f$ (with respect to local coordinates) has maximal rank. Let $T f: T\left(M^{\prime}\right) \rightarrow T\left(\mathbb{R}^{N^{\prime}}\right)$ stand for the tangent mapping of $f$ which is locally defined
by the Jacobi matrix of $f$. Then $T f$ is injective, and so there exists a smooth bundle homomorphism $T f^{(-1)}: f\left(M^{\prime}\right) \times \mathbb{R}^{N^{\prime}} \rightarrow T\left(M^{\prime}\right)$ with the property that $T f^{(-1)} T f$ is the identity mapping of $T\left(M^{\prime}\right)$. Set

$$
\Delta(x, y)=T f^{(-1)}(f(x), f(y)-f(x)), \quad(x, y) \in M^{\prime} \times M^{\prime}
$$

Then it is clear that condition 1) is fulfilled. Moreover, it follows from the Taylor formula that 2) is satisfied.

The double differential forms $K_{q}(\Delta)$ provide us with reproducing kernels for differential forms on $M^{\prime}$.

Lemma 8.4 If $M$ is a relatively compact open subset of $M^{\prime}$ with a piecewise smooth boundary and $u \in C^{1}\left(\Lambda^{q} T^{*}(\bar{M})\right)$, then

$$
\begin{align*}
& -\int_{\partial \bar{M}} u(y) \wedge K_{q+1}(\Delta(x, y)) \\
& \quad+\int_{\bar{M}} d u(y) \wedge K_{q+1}(\Delta(x, y))+d \int_{\bar{M}} u(y) \wedge K_{q}(\Delta(x, y))=\left\{\begin{array}{cc}
u(x), & x \in M \\
0, & x \in M^{\prime} \backslash \bar{M}
\end{array}\right. \tag{8.1}
\end{align*}
$$

Proof. From the invariance of the exterior derivative under a differentiable change of variables it follows that the current $\Delta^{\sharp} K_{1}(z)$ satisfies the fundamental equation

$$
d \Delta^{\sharp} K_{1}(z)=(-1)^{N}\left[\text { diagonal }\left(M^{\prime} \times M^{\prime}\right)\right]
$$

on $M^{\prime} \times M^{\prime}$. Combining this with Lemma 7.1 , we deduce that the family of double forms $K_{q}(\Delta), q=0,1, \ldots, N$, is a fundamental solution of the de Rham complex on $M^{\prime}$. To complete the proof, it suffices to use Corollary 2.5.1 in Tarkhanov [T2].

Having disposed of this preliminary step, we can now invoke Lemma 7.3 to modify formula (8.1) thus arriving at a homotopy formula with an arbitrary "barrier function."

Theorem 8.5 Let $\bar{M}$ be a compact manifold with a boundary of class $C^{2}$ and let $f=f(x, y)$ be a $C^{2}$ mapping of $M \times \partial \bar{M}$ to $\mathbb{R}^{N} \backslash\{0\}$. If $f$ is homotopic to the mapping $\Delta$ on $M \times \partial \bar{M}$ and $f_{t}, t \in[0,1]$, is a $C^{2}$ homotopy between $f$ and $\Delta$, then, for every form $u \in C^{1}\left(\Lambda^{q} T^{*}(\bar{M})\right)$, we have

$$
\begin{equation*}
u(x)=-\int_{\partial \bar{M}} u(y) \wedge K_{q+1}(f(x, y))+h d u+d h u, \quad x \in M \tag{8.2}
\end{equation*}
$$

where

$$
h_{q} u(x)=-\int_{\partial \bar{M} \times[0,1]} u(y) \wedge\left((-1)^{q} K_{q}\left(f_{t}(x, y)\right)\right)+\int_{\bar{M}} u(y) \wedge K_{q}(\Delta(x, y)) .
$$

Proof. Setting $f_{1}=\Delta$ and $f_{0}=f$, we decompose the kernel $K_{q+1}(\Delta(x, y))$ on the set $M \times \partial \bar{M}$ in accordance with equality (7.3). Substituting this decomposition into the integral over the boundary in (8.1) and using Stokes' formula, we derive (8.2), as desired.

Thus, any smooth deformation of the mapping $\Delta(x, y)$ on $M \times \partial \bar{M}$ leads to an integral representation of differential forms in $\bar{M}$.

Example 8.6 Suppose $f=f(x, y)$ is a continuous mapping of $M \times \partial \bar{M}$ to $\mathbb{R}^{N} \backslash\{0\}$ such that $\langle f, \Delta\rangle_{(x, y)}>-|f||\Delta|$ on this set. Then, $f$ is homotopic to $\Delta$. Indeed, consider the continuous mapping $f_{t}=t \Delta+(1-t) f$ of $M \times \partial \bar{M} \times[0,1]$ to $\mathbb{R}^{N}$. In order to prove that $f_{t} \neq 0$, we write

$$
\begin{aligned}
\left|f_{t}\right|^{2} & =t^{2}|\Delta|^{2}+2 t(1-t)\langle\Delta, f\rangle_{(x, y)}+(1-t)^{2}|f|^{2} \\
& >(t|\Delta|-(1-t)|f|)^{2} \\
& \geq 0,
\end{aligned}
$$

whence the desired conclusion follows.

## 9 "Convex" manifolds

In this section, we consider a sample application of Theorem 8.5. To begin, we introduce a class of manifolds to be considered.

Definition 9.1 A manifold $\bar{M}$ is said to be convex if there exists a $C^{1}$ mapping $b: \partial \bar{M} \rightarrow \mathbb{R}^{N}$ such that $\langle b(y), \Delta(x, y)\rangle_{(x, y)}>0$ for all $(x, y) \in M \times \partial \bar{M}$.

The important point to note here is that, in contrast to what was required in Sections 7 and 8 , we are going to apply the above results with $f(x, y)=b(y)$ which is of merely class $C^{1}$. However, the same arguments still go when $f$ is of class $C^{1}$, provided it depends on only one variable $x$ or $y$.

Suppose $\bar{M}$ is convex and $b$ is a mapping guaranteed by Definition 9.1. Using Example 8.6 we conclude that the mapping $b: M \times \partial \bar{M} \rightarrow \mathbb{R}^{N} \backslash\{0\}$ is homotopic to the mapping $\Delta$ on the set $M \times \partial \bar{M}$. Moreover the homotopy constructed above is $f_{t}(x, y)=t \Delta(x, y)+(1-t) b(y)$ which is smooth. Applying now Theorem 8.5 we derive the following series of formulas valid for $u \in C^{1}\left(\Lambda^{g} T^{*}(\bar{M})\right)$ :

$$
\begin{equation*}
u(x)=-\int_{\partial \bar{M}} u(y) \wedge K_{q+1}(b(y))+h d u+d h u, \quad x \in M . \tag{9.1}
\end{equation*}
$$

Since $b$ does not depend on $x$, the kernels $K_{q+1}(b(y))$ vanish provided $q>0$. Therefore the equalities (9.1) all together mean that the operator $h$ is a fundamental solution, at positive degrees, of the de Rham complex on $\bar{M}$.

Example 9.2 Let $M$ be a convex domain in $\mathbb{R}^{N}$ with twice differentiable boundary. Choose a real-valued function $\rho$ of class $C^{2}$ in a neighborhood of $\bar{M}$, such that $M=\{x: \rho(x)<0\}$ and $d \rho(y) \neq 0$ for $y \in \partial \bar{M}$. Since the tangent hyperplane to $\partial \bar{M}$ at a point $y \in \partial \bar{M}$ does not meet $M$, we have $\langle\nabla \rho(y), y-x\rangle_{(x, y)}>0$ on the set $M \times \partial \bar{M}$, where $\nabla \rho$ is the gradient of $\dot{\rho}$. Thus, taking $b(y)=\nabla \rho(y)$, we see that $\bar{M}$ meets Definition 9.1.

Thus Theorem 8.5 shows a way to find new fundamental solutions of the de Rham complex on a manifold with boundary.

## 10 Residue formula

We now consider another method for achieving $K_{q+1}(f(x, y))=0$. If $f(x, y)$ is independent of $y$, then $K_{q+1}(f(x, y))=0$ for $q \leq N-2$ ! This works well for concave boundaries by simply switching the variables $x$ and $y$ for convex manifolds. We discuss in detail the simple case of convex domains in $\mathbb{R}^{N}$ with convex "holes."

Suppose $\bar{M}=\bar{M}_{0} \backslash\left(M_{1} \cup \ldots \cup M_{I}\right)$, where $M_{0}$ and $M_{i} \subset \subset M_{0}$ are convex domains with $C^{2}$ boundaries in $\mathbb{R}^{N}$.

For each $i=0,1, \ldots, I$, we choose a convex functions $\rho_{i} \in C_{l o c}^{2}\left(\mathbb{R}^{N}\right)$ such that $M_{i}=\left\{x \in \mathbb{R}^{N}: \rho_{i}(x)<0\right\}$ and $d \rho_{i}(x) \neq 0$ for $x \in \partial \bar{M}_{i}$. Consider the mapping $f: M \times \partial \bar{M} \rightarrow \mathbb{R}^{N}$ given by

$$
f(x, y)=\left\{\begin{array}{rll}
\nabla \rho_{0}(y) & \text { if } & (x, y) \in M \times \partial \bar{M}_{0} \\
-\nabla \rho_{1}(x) & \text { if } & (x, y) \in M \times \partial \bar{M}_{1} ; \\
\cdots & \cdots & \cdots \\
-\nabla \rho_{I}(x) & \text { if } & (x, y) \in M \times \partial \bar{M}_{I} .
\end{array}\right.
$$

By assumption, the tangent hyperplane to $\partial \bar{M}_{0}$ at a point $y \in \partial \bar{M}_{0}$ does not intersect $M_{0}$. Moreover, for each $i=1, \ldots, I$, the tangent hyperplane to the hypersurface $\rho_{i}^{-1}\left(\rho_{i}(x)\right)$ does not meet $\partial \bar{M}_{i}$ provided $x \notin \bar{M}_{i}$. Hence it follows that $\langle f(x, y), y-x\rangle_{(x, y)} \geq 0$ on the set $M \times \partial \bar{M}$. Using Example 8.6 we deduce that the mapping $f: M \times \partial \bar{M} \rightarrow \mathbb{R}^{N} \backslash\{0\}$ is homotopic to the mapping $\Delta$ on the set $M \times \partial \bar{M}$. Moreover the homotopy constructed above is $f_{t}(x, y)=t(y-x)+(1-t) f(x, y)$ which is smooth.

Applying now Theorem 8.5 we derive the following series of formulas valid for $u \in C^{1}\left(\Lambda^{q} T^{*}(\bar{M})\right):$

$$
\begin{align*}
u(x) & =-\int_{\partial \bar{M}_{0}} u(y) \wedge K_{q+1}\left(\nabla \rho_{0}(y)\right) \\
& -\sum_{i=1}^{I} \int_{\partial \bar{M}_{i}} u(y) \wedge K_{q+1}\left(\nabla \rho_{i}(x)\right)+h d u+d h u, \quad x \in M . \tag{10.1}
\end{align*}
$$

If $q>0$, then $K_{q+1}\left(\nabla \rho_{0}(y)\right)=0$ for $(x, y) \in M \times \partial \bar{M}_{0}$. Thus, (10.1) yields the following "residue formula."

Theorem 10.1 Suppose $1 \leq q \leq N$. Then, for each closed differential form $u \in C^{1}\left(\Lambda^{q} T^{*}(\bar{M})\right)$, we have

$$
\begin{equation*}
u(x)=-\sum_{i=1}^{I} \int_{\partial \bar{M}_{i}} u(y) \wedge K_{q+1}\left(\nabla \rho_{i}(x)\right)+d h u, \quad x \in M . \tag{10.2}
\end{equation*}
$$

We emphasize that, if $q<N-1$, then $K_{q+1}\left(\nabla \rho_{i}(x)\right)=0$ for all $i=1, \ldots, I$. Thus, for such $q$, the sum on the right-hand side of (10.2) vanishes. On the other hand, for $q=N-1$ this sum becomes

$$
-\sum_{i=1}^{I}\left(\int_{c_{i}} u\right) K_{N}\left(\nabla \rho_{i}(x)\right)
$$

where $c_{i}$ is an arbitrary $(N-1)$-dimensional cycle in $M$ surrounding $\bar{M}_{i}$ with multiplicity 1 .

Corollary 10.2 If, for a closed form $u \in C^{1}\left(\Lambda^{N-1} T^{*}(\bar{M})\right)$, all the periods $\int_{c_{i}} u$, $i=1, \ldots, I$, are zero, then $u=d h u$ in $M$.

Proof. This follows from (10.2).

## 11 Rouchet principle

We endow the real space $\mathbb{R}^{2 n}$ with the complex structure $z_{j}=x_{j}+\sqrt{-1} x_{n+j}$, $j=1, \ldots, n$, thus obtaining a complex space $\mathbb{C}^{n}$. Under this structure, one considers the complex derivatives $\partial / \partial z_{j}=\frac{1}{2}\left(\partial / \partial x_{j}-\sqrt{-1} \partial / \partial x_{n+j}\right)$, for $j=1, \ldots, n$.

Let $M$ be a bounded domain in $\mathbb{C}^{n}$ with a boundary of class $C^{2}$. As above, we write $M=\left\{z \in \mathbb{C}^{n}: \rho(z)<0\right\}$ with a suitable real-valued function $\rho \in C_{l o c}^{2}\left(\mathbb{C}^{n}\right)$ satisfying $d \rho(z) \neq 0$ for $z \in \partial \bar{M}$.

For a point $\zeta \in \partial \bar{M}$, we have

$$
\langle\nabla \rho(y), x-y\rangle_{(x, y)}=2 \operatorname{Re}\left\langle\nabla_{\mathbb{C}} \rho(\zeta), z-\zeta\right\rangle_{(z, \zeta)}
$$

where $\zeta_{j}=y_{j}+\sqrt{-1} y_{n+j}, j=1, \ldots, n$, and

$$
\nabla_{\mathbb{C}} \rho(\zeta)=\left(\begin{array}{c}
\partial \rho / \partial \zeta_{1} \\
\ldots \\
\partial \rho / \partial \zeta_{n}
\end{array}\right)
$$

Theorem 11.1 Let $F$ be a holomorphic mapping of $M \rightarrow \mathbb{C}^{n}$, of class $C^{1}(\bar{M})$. If the image of $\partial \vec{M}$ by the mapping $\zeta \mapsto\left\langle\nabla_{\mathbf{C}} \rho, F\right\rangle_{\zeta}$ does not separate 0 from $\infty$ in the complex plane, then the number of zeroes of $F$ in $M$, along with their multiplicities, is equal to $\chi(\bar{M})$.

Proof. It follows from the hypothesis of the theorem that $F$ has no zeroes on the boundary of $\bar{M}$. Hence all zeroes of $F$ in $M$ are isolated and the number of zeroes, along with their multiplicities, is equal to

$$
\begin{equation*}
N=\int_{\partial \bar{M}} \frac{(n-1)!}{(2 \pi \sqrt{-1})^{n}} \sum_{j=1}^{n}(-1)^{j-1} \frac{\bar{F}_{j}}{|F|^{2 n}} d F \wedge d \bar{F}[j] \tag{11.1}
\end{equation*}
$$

(cf. Theorem 2.4 in Aizenberg and Yuzhakov [AYu]).
Since the image of $\partial \bar{M}$ by the mapping $\zeta \mapsto\left\langle\nabla_{\mathbf{c}} \rho, F\right\rangle_{\zeta}$ does not separate 0 from $\infty$ in the complex plane and $\left|\nabla_{\mathbb{C}} \rho(\zeta)\right| \geq c>0$ for $\zeta \in \partial \bar{M}$, it follows that the image of $\partial \bar{M}$ by the mapping

$$
\begin{equation*}
\zeta \mapsto \frac{\left\langle\nabla_{\mathbf{c}} \rho, F\right\rangle_{\zeta}}{\left|\nabla_{\mathbb{C}} \rho(\zeta)\right|^{2}} \tag{11.2}
\end{equation*}
$$

does not separate 0 from $\infty$ in the complex plane. Consequently, there exists a path $p:[0,1) \rightarrow \mathbb{C}$ such that

$$
\begin{aligned}
p(0) & =0 \\
\lim _{t \rightarrow 1}|p(t)| & =\infty
\end{aligned}
$$

and $p(t)$ does not meet the image of $\partial \bar{M}$ by the mapping (11.2). We can certainly assume that $p(t) \neq 1$ for all $t \in[0,1)$. Consider the mapping $F_{t}: \partial \bar{M} \times[0,1] \rightarrow \mathbb{C}^{n}$ given by

$$
F_{t}(\zeta)=\frac{-1}{p(t)-1} F(\zeta)+\frac{p(t)}{p(t)-1} \bar{\nabla}_{\mathbb{C}} \rho(\zeta)
$$

where

$$
\bar{\nabla}_{\mathbb{C}} \rho(\zeta)=\left(\begin{array}{c}
\partial \rho / \partial \bar{\zeta}_{1} \\
\cdots \\
\partial \rho / \partial \bar{\zeta}_{n}
\end{array}\right)
$$

We have $F_{0}(\zeta)=F(\zeta)$ and $F_{1}(\zeta)=\bar{\nabla}_{\mathbf{c}} \rho(\zeta)$. Moreover, an easy verification shows that

$$
\left\langle F_{t}, \nabla_{\mathbf{C}} \rho\right\rangle_{\zeta}=\frac{-1}{p(t)-1}\left\langle F, \nabla_{\mathbf{c}} \rho\right\rangle_{\zeta}+\frac{p(t)}{p(t)-1}\left|\nabla_{\mathbb{C}} \rho(\zeta)\right|^{2}
$$

vanishes only if $p(t)=\frac{\left\langle\nabla_{c} \rho, F\right\rangle_{\ell}}{\left|\nabla_{\mathcal{c}} \rho(\zeta)\right|^{2}}$ for some $\zeta \in \partial \bar{M}$, which is impossible. From what has already been proved we see that $F_{t}: \partial \bar{M} \times[0,1] \rightarrow \mathbb{C}^{n} \backslash\{0\}$ is a smooth homotopy between $F$ and $\bar{\nabla}_{\mathbb{C}} \rho$.

As the integral on the right-hand side of (11.1) is an integer number (see, for instance, Lemma 2.6 in [AYu]), we get

$$
N=\int_{\partial \bar{M}} \frac{(n-1)!}{(2 \pi \sqrt{-1})^{n}} \sum_{j=1}^{n}(-1)^{j-1} \frac{\frac{\partial \rho}{\partial \zeta_{j}}}{\left|\nabla_{\mathbb{C}} \rho\right|^{2 n}} d \bar{\nabla}_{\mathbb{C}} \rho \wedge d \nabla_{\mathbb{C}} \rho[j] .
$$

To complete the proof it remains to use the Poincaré index together with the observation that the differential form under integration is equal to $K_{1}(\nabla \rho)$ up to an
exact form on the boundary of $\bar{M}$. Indeed, for each $C^{1}$ mapping $F$ from an open subset of $\mathbb{C}^{n}$ to $\mathbb{C}^{n}$, we have, away from $F^{-1}(0)$,

$$
\begin{aligned}
& \frac{(n-1)!}{(2 \pi \sqrt{-1})^{n}} \sum_{j=1}^{n}(-1)^{j-1} \frac{\bar{F}_{j}}{|F|^{2 n}} d F \wedge d \bar{F}[j] \\
& \quad=-K_{1}(f)+\sqrt{-1} d\left(\frac{1}{\sigma_{2 n}(2 n-2)} \sum_{j=1}^{n} \frac{1}{|f|^{2 n-2}} d f[j, n+j]\right),
\end{aligned}
$$

where $F_{j}=f_{j}+\sqrt{-1} f_{n+j}, j=1, \ldots, n^{3}$.
Obviously, if $\operatorname{Re}\left\langle\nabla_{\mathbf{c}} \rho, F\right\rangle_{\zeta} \neq 0$ for all $\zeta \in \partial \bar{M}$, then the cycle $\left\{\left\langle\nabla_{\mathbf{c}} \rho, F\right\rangle_{\zeta}: \zeta \in\right.$ $\partial \bar{M}\}$ in the complex plane does not separate 0 from $\infty$. Hence the number of zeroes of $F$ in $M$, along with their multiplicities, is equal to the Euler characteristic of $\bar{M}$.

The condition " $\operatorname{Re}\left\langle\nabla_{\mathbf{C}} \rho, F\right\rangle_{\zeta} \neq 0$ for all $\zeta \in \partial \bar{M}$ " cannot be relaxed to the condition " $\left\langle\nabla_{\mathbb{C}} \rho, F\right\rangle_{\zeta} \neq 0$ for each $\zeta \in \partial \bar{M}$."

Example 11.2 For the mapping $F(z)=z^{2}$ of the closed unit disc in $\mathbb{C}^{1}$, we have

$$
\begin{aligned}
\left\langle\nabla_{\mathbb{C}} \rho, F\right\rangle_{\zeta} & =\left\langle\frac{\partial}{\partial \zeta}(\sqrt{\zeta \underline{\zeta}}-1), \zeta^{2}\right\rangle_{\zeta} \\
& =\frac{1}{2} \zeta \\
& \neq 0
\end{aligned}
$$

for $|\zeta|=1$, while $F$ vanishes at the origin with multiplicity 2 .

We mention a corollary of Theorem 11.1 which concerns the so-called linearly convex domains in complex analysis (cf. [AYu, §24]).

Corollary 11.3 If there is a point $z^{0} \in M$ with the property that the cycle $\left\{\left\langle\nabla_{\mathbf{C}} \rho(\zeta), z^{0}-\zeta\right\rangle_{\zeta}: \zeta \in \partial \bar{M}\right\}$ in the complex plane does not separate 0 from $\infty$, then $\chi(\bar{M})=1$.

Proof. Apply Theorem 11.1 to $F(z)=z-z^{0}$.
If $M$ is a strictly pseudoconvex domain in $\mathbb{C}^{n}$ with $C^{2}$ boundary, then the Levi determinant of the function $\rho$

$$
\mathcal{L}(\rho)=-\operatorname{det}\left(\begin{array}{cccc}
0 & \frac{\partial \rho}{\partial \zeta_{1}} & \cdots & \frac{\partial \rho}{\partial \zeta_{n}} \\
\frac{\partial \rho}{\partial \zeta_{1}} & \frac{\partial^{2} \rho}{\partial \zeta_{1} \partial \zeta_{1}} & \cdots & \frac{\partial^{2} \rho}{\partial \zeta_{n} \partial \zeta_{1}} \\
\cdots & \cdots & \cdots & \cdots \\
\frac{\partial \rho}{\partial \zeta_{n}} & \frac{\partial^{2} \rho}{\partial \zeta_{1} \partial \zeta_{n}} & \cdots & \frac{\partial^{2} \rho}{\partial \zeta_{n} \partial \zeta_{n}}
\end{array}\right)
$$

[^3]is positive at each point $\zeta \in \partial \bar{M}$ (cf. [AYu, $\S 8]$ ). Then, instead of comparing $F$ with the complex gradient of $\rho$, we may compare it with the mapping of $\bar{M} \rightarrow \mathbb{C}^{n}$ whose components are complementary minors of the elements $\partial \rho / \partial \zeta_{j}, j=1, \ldots, 2$, in the Levi determinant $\mathcal{L}(\rho)$. In this way we obtain the following result.

Theorem 11.4 Suppose $M$ is a strictly pseudoconvex domain with twice differentiable boundary in $\mathbb{C}^{n}$ and $F$ is a holomorphic mapping of $M \rightarrow \mathbb{C}^{n}$, of class $C^{1}(\bar{M})$. If the image of $\partial \bar{M}$ by the mapping

$$
\zeta \mapsto \operatorname{det}\left(\begin{array}{cccc}
0 & F_{1} & \ldots & F_{n} \\
\frac{\partial \rho}{\partial \zeta_{1}} & \frac{\partial^{2} \rho}{\partial \zeta_{1} \partial \zeta_{1}} & \ldots & \frac{\partial^{2} \rho}{\partial \zeta_{n} \partial \zeta_{1}} \\
\ldots & \ldots & \ldots & \ldots \\
\frac{\partial \rho}{\partial \zeta_{n}} & \frac{\partial^{2} \rho}{\partial \zeta_{1} \partial \zeta_{n}} & \ldots & \frac{\partial^{2} \rho}{\partial \zeta_{n} \partial \zeta_{n}}
\end{array}\right)
$$

does not separate 0 from $\infty$ in the complex plane, then the number of zeroes of $F$ in $M$, along with their multiplicities, is equal to $\chi(\bar{M})$.

## A A proof of the classical Hopf formula

Let $f: M \rightarrow T(M)$ be a vector field on a oriented compact closed manifold $M$ of dimension $N$. Suppose all critical points of $f$ are non-degenerate.

Recall that each vector field $f$ on $M$ generates a flow on $M$, that is a family of diffeomorphisms $F_{t}(x)$ of $M$, parametrized by $t \in[0, T)$. For a point $x \in M$, it is obtained by solving the system of ordinary differential equations $\frac{d}{d t} F_{t}(x)=$ $f(x)$ with the initial value $F_{0}(x)=x$. (In local coordinates near $x$, write $f(x)=$ $\sum_{j=1}^{N} f_{j}(x) \partial / \partial x_{j}$, then the system becomes $\frac{d}{d t}\left(F_{t}\right)_{j}(x)=f_{j}(x), j=1, \ldots, N$.) It is well-known from the theory of ordinary differential equations that the $F_{t}(x)$ exists for each $x$, provided $T$ is small enough, and depends smoothly on $x$ varying over a compact set (see Milnor [M, p.10] for more details).

We now invoke the general Lefschetz Fixed Point Theorem for $C W$-complexes (cf. Dold [D]). Namely, each smooth mapping $F: M \rightarrow M$ is known to induce a homomorphism ("pull-back") $F^{\sharp}$ of the cohomology $H(M)$ of $M^{4}$. Thus, for each $i=0,1, \ldots, N$, the $\left.\operatorname{trace} \operatorname{tr} F^{\sharp}\right|_{H^{i}(M)}$ is well-defined. On the other hand, if $\sigma$ is a closed subset of $M$ and $U$ a neighborhood of $\sigma$, then $F$ induces a smooth mapping of the pair $\left(F^{-1}(U), F^{-1}(U) \backslash F^{-1}(\sigma)\right) \rightarrow(U, U \backslash \sigma)$ and consequently a homomorphism of the relative cohomology

$$
F^{\sharp}: H\left(F^{-1}(U), F^{-1}(U) \backslash F^{-1}(\sigma)\right) \rightarrow H(U, U \backslash \sigma)
$$

In particular, if $\sigma$ is invariant under $F$, i.e., $F(\sigma) \subset \sigma$, this gives rise to a homomorphism $F^{\sharp}$ of $H \cdot(U, U \backslash \sigma)$. Then, the Lefschetz Fixed Point Theorem states that if

[^4]the set of fixed points of $F$ is discrete, then
\[

$$
\begin{equation*}
\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F^{\sharp}\right|_{H^{i}(M)}=\sum_{p=F(p)}\left(\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F^{\sharp}\right|_{H^{i}(U, U \backslash p)}\right) . \tag{A.1}
\end{equation*}
$$

\]

In case $p$ is an interior point of $M$ the contribution $\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F^{\natural}\right|_{H^{i}(U, U \backslash p)}$ of $p$ reduces to the local degree of $1-\frac{\partial F}{\partial x}$ at $p$.

We are going to apply (A.1) to the diffeomorphism $F_{t}$. Granting that the set of fixed points of $F_{t}$ is discrete, we obtain

$$
\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F_{t}^{*}\right|_{H^{i}(M)}=\sum_{p=F_{t}(p)} \operatorname{deg}\left(1-\frac{\partial F_{t}}{\partial x}, p\right)
$$

The left-hand side of this equality is independent of $t$ small enough, for homotopic mappings of $M \rightarrow M$ induce the same homomorphism of the cohomology. Therefore, letting $t \rightarrow 0$, we deduce that

$$
\begin{aligned}
\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F_{t}^{*}\right|_{H^{i}(M)} & =\left.\sum_{i=0}^{N}(-1)^{i} \operatorname{tr} F_{0}^{*}\right|_{H^{i}(M)} \\
& =\chi(M)
\end{aligned}
$$

for all $t \in[0, T)$.
On the other hand, as the solution of the local Cauchy problem for ordinary differential equations is unique, we can assert that a point $p \in M$ is a fixed point of $F_{t}$ if and only if $p$ is a critical point of $f$, i.e., $f(p)=0$. Moreover, we claim that

$$
\operatorname{deg}\left(1-\frac{\partial F_{t}}{\partial x}, p\right)=(-1)^{N} \operatorname{deg}(f, p)
$$

Indeed, choose a local chart on $M$ with center at $p$, then

$$
f_{j}(x)=\sum_{k=1}^{N} c_{j k} x_{k}+O\left(|x|^{2}\right), \quad j=1, \ldots, N
$$

for $x$ close to $p$. Hence it follows that

$$
\left(F_{t}\right)_{j}(x)=x_{j}+\sum_{k=1}^{N} c_{j k} x_{k} t+O\left(t^{2}+|x|^{2}\right), \quad j=1, \ldots, N,
$$

provided $t^{2}+|x|^{2} \ll 1$. Thus,

$$
\begin{aligned}
\operatorname{deg}\left(1-\frac{\partial F_{t}}{\partial x}, p\right) & =\operatorname{sign} \operatorname{det}\left(1-\frac{\partial F_{t}}{\partial x}\right)(p) \\
& =\operatorname{sign} \operatorname{det}\left(-c_{j k}\right) \\
& =(-1)^{N} \operatorname{deg}(f, p),
\end{aligned}
$$

as desired.
Summarizing, we obtain $\sum_{p \in f^{-1}(0)} \operatorname{deg}(f, p)=(-1)^{N} \chi(M)$. We now apply this argument again, with $f$ replaced by $-f$, to obtain $\sum_{p \in f^{-1}(0)} \operatorname{deg}(f, p)=\chi(M)$. These formulas are certainly equivalent for, by the Poincaré duality, the Euler characteristic of an odd-dimensional manifold is equal to zero (this follows also from our argument). The proof is complete.

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[^0]:    *Supported by the Deutsche Forschungsgemeinschaft.

[^1]:    ${ }^{1}$ A connection with this property is called Euclidean.

[^2]:    ${ }^{2}$ Here, $\delta(z)$ stands for the Dirac functional in $\mathbb{R}^{N}$.

[^3]:    ${ }^{3}$ If $n=1$, one must replace $\frac{1}{2 n-2}\left[\frac{1}{\left[\mid{ }^{2 n-2}\right.}\right.$ by $\log |f|$.

[^4]:    ${ }^{4}$ Here, $M$ need not be closed.

