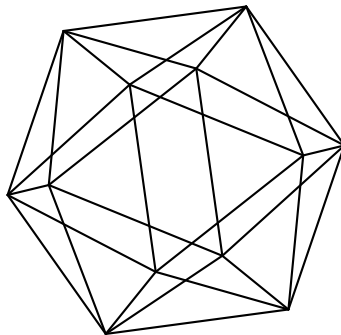


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Fabian Rupp
Christian Scharrer



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Fabian Rupp
Christian Scharrer

Max-Planck-Institut für Mathematik
Vivatsgasse 7
53111 Bonn
Germany

Faculty of Mathematics
University of Vienna
Oskar-Morgenstern-Platz 1
1090 Vienna
Austria

Li–Yau inequalities for the Helfrich functional and applications

Fabian Rupp* Christian Scharrer†

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Abstract: We prove a general Li–Yau inequality for the Helfrich functional where the spontaneous curvature enters with a singular volume type integral. In the physically relevant cases, this term can be converted into an explicit energy threshold that guarantees embeddedness. We then apply our result to the spherical case of the variational Canham–Helfrich model. If the infimum energy is not too large, we show existence of smoothly embedded minimizers. Previously, existence of minimizers was only known in the class of immersed bubble trees.

Keywords: Canham–Helfrich model, Li–Yau inequality, embeddedness, Willmore energy, biological membranes, oriented varifolds, Alexandrov immersions.

MSC(2020): 53C42 (primary), 49Q10, 92C10 (secondary).

1 Introduction

Lipid bilayers make up the cellular membranes of most organisms. These extremely thin structures commonly form vesicles, so mathematically they are naturally modeled as two-dimensional structures, i.e. closed surfaces. The *Canham–Helfrich model* [8, 17] characterizes the equilibrium shapes of lipid bilayers as (constrained) minimizers of a curvature dependent bending energy. For an oriented surface Σ and an immersion $f: \Sigma \rightarrow \mathbb{R}^3$, the *Helfrich energy* is defined by

$$\mathcal{H}_{c_0}(f) := \frac{1}{4} \int_{\Sigma} |H - c_0 n|^2 d\mu.$$

*Faculty of Mathematics, University of Vienna, Oskar-Morgenstern-Platz 1, 1090 Vienna, Austria.
fabian.rupp@univie.ac.at

†Max Planck Institute for Mathematics, Vivatsgasse 7, 53111 Bonn, Germany.
scharrer@mpim-bonn.mpg.de

Here, $H = H_f$ is the mean curvature vector of the immersion, i.e. the trace of the second fundamental form, $n = n_f$ is the unit normal induced by the orientation of Σ (see (2.6) below) and $\mu = \mu_f$ denotes the Riemannian measure associated to the pullback metric $g = g_f = f^*\langle \cdot, \cdot \rangle$ on Σ , where $\langle \cdot, \cdot \rangle$ denotes the Euclidean inner product in \mathbb{R}^3 . The constant $c_0 \in \mathbb{R}$ is called *spontaneous curvature*. Since H is normal to the surface, the Helfrich energy can also be written as

$$\mathcal{H}_{c_0}(f) = \frac{1}{4} \int_{\Sigma} (H_{\text{sc}} - c_0)^2 d\mu,$$

where $H_{\text{sc}} := \langle H, n \rangle$ is the scalar mean curvature with respect to n . Reversing the orientation on Σ corresponds precisely to replacing n by $-n$. Thus, we have

$$\mathcal{H}_{c_0}(f) = \mathcal{H}_{-c_0}(\hat{f}), \quad (1.1)$$

where $\hat{\Sigma}$ is the surface Σ with reversed orientation and $\hat{f}: \hat{\Sigma} \rightarrow \mathbb{R}^3$, $\hat{f}(p) = f(p)$. Clearly, the Helfrich functional is not scale-invariant. However, we observe the following scaling property involving both arguments:

$$\mathcal{H}_{c_0}(f) = \mathcal{H}_{\frac{c_0}{r}}(rf) \quad \text{for } r > 0.$$

In particular, we see that

$$\lim_{r \rightarrow 0^+} \mathcal{H}_{c_0}(rf) = \lim_{r \rightarrow 0^+} \mathcal{H}_{r c_0}(f) = \mathcal{H}_0(f) =: \mathcal{W}(f). \quad (1.2)$$

The right hand side is well known as the *Willmore energy*. In contrast to the Helfrich functional, the Willmore functional \mathcal{W} is scale-invariant and does not depend on the unit normal field n or the orientation of the underlying surface Σ . One may also consider the L^2 -CMC-deficit

$$\bar{\mathcal{H}}(f) := \inf_{c_0 \in \mathbb{R}} \mathcal{H}_{c_0}(f) = \frac{1}{4} \int_{\Sigma} (H_{\text{sc}} - \bar{H}_{\text{sc}})^2 d\mu, \quad (1.3)$$

where $\bar{H}_{\text{sc}} := \int_{\Sigma} H_{\text{sc}} d\mu$ is the average scalar mean curvature. Also the functional $\bar{\mathcal{H}}$ is scale-invariant and does not depend on the orientation of Σ . For more details and corresponding results, see Section 5.2 and Section 6.4.

We are primarily interested in the following minimization problem suggested by Canham [8] and Helfrich [17] in order to study the shape of red blood cells. Our main contribution is stated in Theorem 1.6 below (see also Theorem 6.10).

Problem 1.1. Let $c_0 \in \mathbb{R}$ and $A_0, V_0 > 0$ be given constants. Let the unit sphere $\mathbb{S}^2 \subset \mathbb{R}^3$ be oriented by the inner unit normal. Minimize the functional \mathcal{H}_{c_0} in the class of smooth embeddings $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ subject to the constraints

$$\mathcal{A}(f) := \int_{\mathbb{S}^2} 1 d\mu = A_0, \quad \mathcal{V}(f) := -\frac{1}{3} \int_{\mathbb{S}^2} \langle f, n \rangle d\mu = V_0.$$

We consider the following example of Problem 1.1 where the infimum cannot be attained by a smooth embedding, cf. [29].

Example 1.2. Let $\iota_{\mathbb{S}^2}: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ be the inclusion of the unit sphere. Let $c_0 := 2$, $A_0 := 2\mathcal{A}(\iota_{\mathbb{S}^2})$, and $V_0 := 2\mathcal{V}(\iota_{\mathbb{S}^2})$. There exists a sequence of smooth embeddings $f_k: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ satisfying $\mathcal{A}(f_k) = A_0$ and $\mathcal{V}(f_k) = V_0$ which converges in the varifold topology to the set $T \subset \mathbb{R}^3$ given by two translations of the unit sphere that meet in exactly one point (see Figure (1)) such that

$$\lim_{k \rightarrow \infty} \mathcal{H}_{c_0}(f_k) = 2\mathcal{H}_{c_0}(\iota_{\mathbb{S}^2}) = 0.$$

In other words, the infimum of Problem 1.1 is attained by the set T . Clearly, T cannot be written as the image of a smooth immersion of \mathbb{S}^2 . If on the other hand $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ is a smooth immersion with $\mathcal{H}_{c_0}(f) = 0$ then, by a result of Hopf [18, Chapter VI, Theorem 2.1], the image of f must be a round sphere. In particular, if $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ is a smooth embedding with $\mathcal{H}_{c_0}(f) = 0$, then $\mathcal{A}(f) \neq A_0$ and $\mathcal{V}(f) \neq V_0$.

In the terminology of [28, 29], the set T in Example 1.2 is the *bubble tree* consisting of two unit spheres. Bubbling phenomena have also been observed in nature and are known as *budding transition* [39]. Thus, the space of bubble trees appears to be a natural class in which to minimize the Helfrich functional. Indeed, in [29, Theorem 1.7], the existence of minimizers for the Helfrich functional in the class of immersed bubble trees was proven. Each of the bubbles is given by a map $\mathbb{S}^2 \rightarrow \mathbb{R}^3$ which outside of finitely many so called *branch points* is a smooth immersion. For similar results, see [9, 13]. However, not all minimizers of Problem 1.1 are necessarily bubble trees, consider for instance the case $c_0 = 2$, $A_0 = \mathcal{A}(\iota_{\mathbb{S}^2})$, $V_0 = \mathcal{V}(\iota_{\mathbb{S}^2})$. One may conjecture that bubbling can only occur if the parameters A_0 and V_0 are within a certain range depending on c_0 . Apart from the geometric relevance to obtain such qualitative results for the minimizers of Problem 1.1, it is of great interest to confirm mathematically that the Canham–Helfrich model is suitable for the study of red blood cells which are actually embedded — rather than a bubble tree. As a first step in this direction it was proven in [29] that there exists a constant $\varepsilon = \varepsilon(A_0, V_0) > 0$ such that the minimizers are given by smooth embeddings provided $|c_0| < \varepsilon$. However, apart from the fact that $\varepsilon(A_0, V_0)$ is implicitly small, one would rather want to have a criterion of the following type: *For all $c_0 \in \mathbb{R}$, the Problem 1.1 has a solution provided A_0 and V_0 are in a certain explicit range depending on c_0 .*

The proof of embeddedness of minimizers in [29] is based on the fact that for $|c_0|$ small, the Helfrich functional is close to the Willmore functional, see (1.2), and minimizers for $c_0 = 0$ are given by smooth embeddings, see [38]. A crucial tool to prove smoothness and embeddedness of the minimizers in [38] (i.e. solutions of Problem 1.1 for $c_0 = 0$) is the following inequality of Li and Yau [23, Theorem 6]. If Σ is *closed* (i.e. compact and without boundary), for any $x_0 \in \mathbb{R}^3$ we have

$$\mathcal{H}^0(f^{-1}\{x_0\}) \leq \frac{1}{4\pi} \mathcal{W}(f), \tag{1.4}$$

where \mathcal{H}^0 denotes the counting measure. In particular, if $\mathcal{W}(f) < 8\pi$, then f must be an embedding. This observation also played an essential role in the study of the Willmore energy and related topics, cf. for instance [40, 22, 38, 21, 30, 20, 32].

In view of the fact that branch points have multiplicity at least 2, such a tool could be the key to exclude bubbling in the Canham–Helfrich model. Apart from comparing the Helfrich energy with the Willmore energy for small $|c_0|$ via (1.2), one might also try to make use of the Li–Yau inequality (1.4) by estimating the Willmore energy from above in terms of the Helfrich energy (see (2.5)):

$$\mathcal{W}(f) \leq 2\mathcal{H}_{c_0}(f) + \frac{1}{2}c_0^2\mathcal{A}(f). \quad (1.5)$$

Again, if the right hand side is strictly less than 8π , then the Li–Yau inequality (1.4) implies that f is an embedding. However, as one of our results reveals (see Lemma 6.1), in the case of red blood cells where $c_0 < 0$ (see [11]), there holds $\mathcal{H}_{c_0}(f) > 4\pi$. In particular, the right hand side of (1.5) is already strictly larger than 8π and one cannot apply the Li–Yau inequality (1.4) to deduce embeddedness of f .

Another naive attempt to apply (1.4) would be to show that $\mathcal{W} \leq \mathcal{H}_{c_0}$ for $c_0 < 0$. However, this is impossible by the following simple scaling argument. Let $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ be an immersion such that $\int_{\mathbb{S}^2} H_{sc} d\mu < 0$ (such an f exists by [10, Theorem 1.2]). We find that

$$\mathcal{H}_{c_0}(rf) - \mathcal{W}(rf) = -\frac{rc_0}{2} \int_{\mathbb{S}^2} H_{sc} d\mu + \frac{r^2c_0^2}{4} \mathcal{A}(f) \quad (1.6)$$

which becomes negative if $r > 0$ is sufficiently small as $c_0 < 0$.

1.1 Main results

Instead of applying (1.4) by comparing the Helfrich energy with the Willmore energy, the aim of this article is to prove and apply a Li–Yau type inequality directly for the Helfrich functional. In the smooth setting, our multiplicity inequality reads as follows.

Lemma 1.3. *Let $f: \Sigma \rightarrow \mathbb{R}^3$ be a smooth proper immersion of an oriented surface Σ without boundary. Let $c_0 \in \mathbb{R}$, $x_0 \in \mathbb{R}^3$, and suppose that the concentrated volume of f at x_0 defined by*

$$\mathcal{V}_c(f, x_0) := - \int_{\Sigma} \frac{\langle f - x_0, n \rangle}{|f - x_0|^2} d\mu \quad (1.7)$$

exists. Then

$$\mathcal{H}^0(f^{-1}\{x_0\}) \leq \limsup_{\rho \rightarrow \infty} \frac{\mu(f^{-1}(B_\rho(x_0)))}{\pi\rho^2} + \frac{1}{4\pi}\mathcal{H}_{c_0}(f) + \frac{c_0}{2\pi}\mathcal{V}_c(f, x_0). \quad (1.8)$$

In order to apply Equation (1.8), it is of crucial interest to determine the sign of the concentrated volume. Despite singular, the integrand in (1.7) is subcritical and locally integrable, see Lemma 3.2 and Remark 3.3 below. Moreover, the integrand is nonpositive if $f[\Sigma]$ parametrizes the boundary of an open set in \mathbb{R}^3 which is star-shaped with respect to x_0 and n is the inner unit normal, cf. [14, 9.4.2]. However, such an immersion f must be embedded a priori.

It turns out that the sign of the concentrated volume can be determined if we can find a suitable notion of inner unit normal, resulting in an appropriate divergence theorem.

Definition 1.4. We call a smooth immersion $f: \Sigma \rightarrow \mathbb{R}^3$ of a closed surface Σ an *Alexandrov immersion* if there exist a smooth compact 3-manifold M with boundary $\partial M = \Sigma$, a smooth inner unit normal field ν to Σ and a smooth immersion $F: M \rightarrow \mathbb{R}^3$ such that $F|_{\Sigma} = f$. The surface Σ is then necessarily orientable. Moreover, we choose the orientation on Σ such that the induced normal field along f (see (2.6) below) satisfies $n = dF(\nu)$.

Our orientation on Σ does not coincide with the usual Stokes orientation. The reason for this is that we want to work with the inner unit normal such that the standard embedding of a round sphere has positive scalar mean curvature.

In the setting of Definition 1.4, the Li–Yau inequality (1.8) can be put into the following more convenient form.

Theorem 1.5. *Let Σ be a closed surface and let $f: \Sigma \rightarrow \mathbb{R}^3$ be an Alexandrov immersion with $f = F|_{\Sigma}$, $F: M \rightarrow \mathbb{R}^3$ as in Definition 1.4. Then for all $x_0 \in \mathbb{R}^3$ we have*

$$\mathcal{H}^0(f^{-1}\{x_0\}) \leq \frac{1}{4\pi} \mathcal{H}_{c_0}(f) + \frac{c_0}{2\pi} \int_{F[M]} \frac{\mathcal{H}^0(F^{-1}\{x\})}{|x - x_0|^2} d\mathcal{L}^3(x).$$

In particular, in case $c_0 \leq 0$ we infer

$$\mathcal{H}^0(f^{-1}\{x_0\}) \leq \frac{1}{4\pi} \mathcal{H}_{c_0}(f).$$

Due to round spheres, the above extension of (1.4) can only hold if $c_0 \leq 0$ and n is the inner unit normal. Of course, in view of (1.1) we could simply reverse the orientation on Σ , but this will generically make it impossible to find an Alexandrov immersion where M in Definition 1.4 is compact.

As a key application of our Li–Yau inequalities, we prove the following contribution to Problem 1.1 based on the previous result in [29].

Theorem 1.6. *Let $c_0 \in \mathbb{R}$ and suppose $A_0, V_0 > 0$ satisfy the isoperimetric inequality $36\pi V_0^2 \leq A_0^3$. Set*

$$\eta(c_0, A_0, V_0) := \inf\{\mathcal{H}_{c_0}(f) \mid f \in C^\infty(\mathbb{S}^2; \mathbb{R}^3) \text{ embedding, } \mathcal{A}(f) = A_0, \mathcal{V}(f) = V_0\}. \quad (1.9)$$

There exists $\Gamma(c_0, A_0, V_0) > 0$ such that if

$$\eta(c_0, A_0, V_0) < \begin{cases} 8\pi + \Gamma(c_0, A_0, V_0) & \text{if } c_0 < 0, \\ 8\pi - \Gamma(c_0, A_0, V_0) & \text{if } c_0 > 0, \end{cases}$$

then the infimum in (1.9) is attained. Moreover, $\lim_{c_0 \rightarrow 0} \Gamma(c_0, A_0, V_0) = 0$ for all $A_0, V_0 > 0$ and for any $c_0 \leq 0$ there exist $A_0, V_0 > 0$ with $\eta(c_0, A_0, V_0) < 8\pi$.

Thanks to further geometric applications of our Li–Yau inequality (see Lemma 6.4), we are even able to give explicit estimates of the constant $\Gamma(c_0, A_0, V_0)$, see Remark 6.11. As a consequence of Theorem 1.6, the only missing step to exclude bubbling in Problem 1.1 are estimates from above for $\eta(c_0, A_0, V_0)$. These could be derived numerically.

1.2 A suitable setup for the Li–Yau inequalities

We now discuss the different notions of (generalized) surfaces that we want to prove and apply Li–Yau inequalities for. There are four key points to be considered.

- (i) In order to even define the Helfrich energy, the surface needs to have a unit normal vector field. In the smooth case, this naturally means that the surface is orientable.
- (ii) One of the main applications of the classical Li–Yau inequality for the Willmore functional is to deduce embeddedness of immersions whose energy lies strictly below 8π . Therefore, the Li–Yau inequality should hold for surfaces that are not already embedded a priori, i.e. we want to allow for multiplicity points.
- (iii) In order to actually apply the Li–Yau inequality for the Helfrich energy (see Lemma 1.3), it is necessary to determine the sign of the concentrated volume (1.7). A sufficient tool to do so would be a divergence theorem.
- (iv) Another important application of the classical Li–Yau inequality is to infer regularity and embeddedness of minimizers. It is therefore of interest to prove Li–Yau inequalities for weak surfaces that have good compactness properties.

Oriented varifolds The most general notion of surface that comprises all shapes shown in Figures (1)-(9) and that naturally satisfies Items (i), (ii), and (iv) are *oriented varifolds*. They generalize the idea of immersed submanifolds and allow for a generalized concept of mean curvature. Since they also possess strong compactness properties, they have already been applied in several variational settings for the Canham–Helfrich model, see [12, 13, 6]. Our most general version of the Li–Yau inequality for oriented varifolds, Theorem 4.2, is also applicable if the first variation has a nontrivial singular part β (see Hypothesis 2.1). The reason for this generality is that we would like the Li–Yau inequality to be applicable for surfaces like the one shown in Figure (7), see also Example 4.7. Moreover, the Li–Yau inequality can then also be applied in the context of boundary problems, see [12].

Alexandrov immersions In nature, one expects the principle of noninterpenetration of matter to hold true. As for vesicles that means there is a clearly defined inside. Nevertheless, membranes can be squeezed together as in Figures (2), (4), and (5). In order to satisfy a divergence theorem, a surface should possess a well defined inside. In the smooth case, the so called *Alexandrov immersions* (see Definition 1.4) do satisfy a divergence theorem, see Lemma 5.1 below. They allow for multiplicity points as shown in Figures (2) and (3). Moreover, since the underlying 3-manifold of an Alexandrov immersion does not have to be connected, they also allow for multiplicity points that arise from two touching surfaces as shown in Figure (1) or even two intersecting surfaces as shown in Figure (8). However, the rotationally symmetric surface in Figure (4) is not an Alexandrov immersion. The Li–Yau inequality for Alexandrov immersions is stated in Theorem 1.5.

Sets of finite perimeter A nonsmooth notion of surfaces that satisfy a divergence theorem are the boundaries of *sets of finite perimeter*. As opposed to Alexandrov immersions, they allow for multiplicity points as shown in Figure (4) but they do not allow for the multiplicity points in Figure (3). Sets of finite perimeter do have good compactness properties. Moreover, they comprise nonsmooth objects as shown in Figure (7) and discussed in Example 4.7. In Section 4.3, we introduce a weak notion of Alexandrov immersions, the *varifolds with enclosed volume*. Their underlying 3-dimensional structure is a sequence of decreasing sets of finite perimeter. They allow for multiplicity points as in (1)–(5) and (8) and still satisfy a divergence theorem. The corresponding Li–Yau inequality is stated in Corollary 4.11.

Currents Another important class of surfaces that naturally satisfies Items (i) and (iv) above are *currents* (see [16, Chapter 4]). A downside of this concept is that the current associated with the immersion of Figure (5) corresponds to the surface shown in Figure (6). More precisely, a current induced by an immersion with a given unit normal field loses information about multiplicity points that arise by overlapping where the sum of the unit normal vectors vanishes, cf. also (6.18). As a consequence, the varifold corresponding to the surface in Figure (6) has a nontrivial singular part while the varifold corresponding to the immersion of Figure (5) has no singular part.

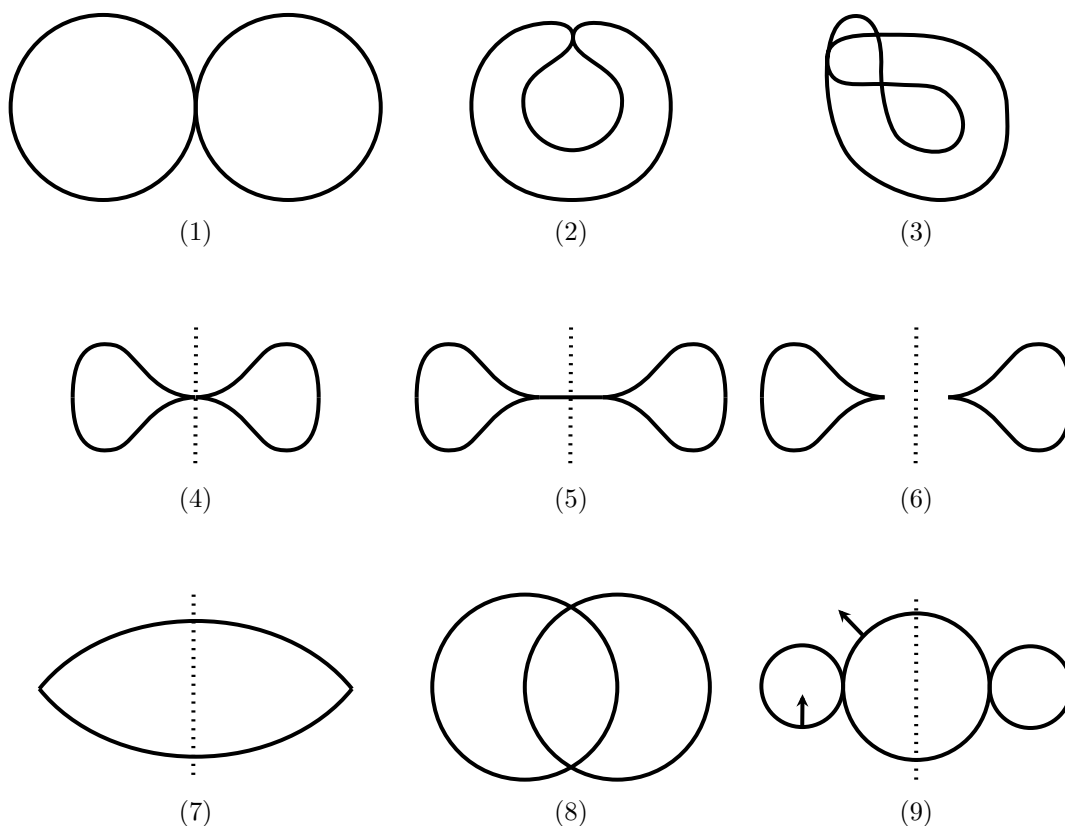
Lipschitz cells An appropriate concept to model cellular membranes is what we term *Lipschitz cells*, a particular class of weak branched immersions (see Section 6.3) which satisfy Items (i)–(iv) above. They describe spherical shapes and comprise the surfaces in Figures (2), (4) and (5), but do not allow for interpenetration as in (3). It turns out that this class is well-suited for the variational discussion of the spherical Canham–Helfrich model which is why we rely on it for the proof of Theorem 1.6.

The only kind of surface where the sign of the concentrated volume cannot be determined in general are those surfaces where the unit normal vector field changes between inner and outer, see Figure (9) and Example 3.8. These are surfaces where interpenetration necessarily happens.

1.3 Structure of this article

After a brief discussion of the geometric and measure theoretic background in Section 2, we examine the concentrated volume and its properties in Section 3. This includes Hölder continuity in $x_0 \in \mathbb{R}^3$ and continuity with respect to varifold convergence. In Section 4, we then derive a monotonicity formula for the Helfrich functional, from which we deduce our most general Li–Yau inequality for varifolds, Theorem 4.2. After that, we review the notion of sets of finite perimeter and introduce the concept of varifolds with enclosed volume. The Li–Yau inequalities in the smooth setting, Lemma 1.3 and Theorem 1.5 are then a direct application, see Section 5. Finally, in Section 6, we derive some geometric estimates and discuss implications of our results. This includes a nonexistence result for

Figure: Profiles of surfaces with different types of multiplicity points. Dotted lines indicate rotationally symmetric surfaces.



the penalized version of Problem 1.1 (Section 6.2), diameter bounds (Section 6.2), the existence and regularity result for the Canham–Helfrich model, Theorem 1.6 (Section 6.3), and a criterion for positive total mean curvature (Section 6.4).

2 Preliminaries

In this section, we will review some of the concepts and tools used throughout this article.

2.1 Notation and definitions

Let μ be a Radon measure over \mathbb{R}^3 and define the closed balls

$$B_\rho(x) := \{y \in \mathbb{R}^3 \mid |x - y| \leq \rho\}$$

for all $x \in \mathbb{R}^3$ and $\rho > 0$. For each nonnegative integer m and $x \in \mathbb{R}^3$, the m -dimensional *lower* and *upper density* of μ at x are defined by

$$\theta_*^m(\mu, x) := \liminf_{\rho \rightarrow 0^+} \frac{\mu(B_\rho(x))}{\omega_m \rho^m}, \quad \theta^{*m}(\mu, x) := \limsup_{\rho \rightarrow 0^+} \frac{\mu(B_\rho(x))}{\omega_m \rho^m},$$

where $\omega_m = \mathcal{L}^m(B_1(0))$ and \mathcal{L}^m is the m -dimensional Lebesgue measure over \mathbb{R}^m . The m -dimensional *density* of μ at x is defined by

$$\theta^m(\mu, x) := \lim_{\rho \rightarrow 0^+} \frac{\mu(B_\rho(x))}{\omega_m \rho^m},$$

provided the limit exists. We define the *support* of μ by

$$\text{spt } \mu := \mathbb{R}^3 \setminus \{x \in \mathbb{R}^3 \mid \exists \rho > 0 \text{ such that } \mu(B_\rho(x)) = 0\}.$$

The m -dimensional Hausdorff measure in Euclidean space is denoted with \mathcal{H}^m . We say that an integral exists if and only if it exists in the Lebesgue sense (i.e. its integrand is *summable* in the terminology of [16, 2.4.2]).

2.2 Oriented 2-varifolds

Let $\mathbb{G}^\circ(3, 2)$ be the set of *oriented 2-dimensional subspaces* of \mathbb{R}^3 . In view of [16, 3.2.28(2)], we identify $\mathbb{G}^\circ(3, 2)$ with

$$\{\xi \in \bigwedge_2 \mathbb{R}^3 \mid \xi \text{ is simple, } |\xi| = 1\}$$

which is a smooth submanifold of the 2-nd exterior power $\bigwedge_2 \mathbb{R}^3$. In particular, $\mathbb{G}^\circ(3, 2)$ is a locally compact Hausdorff space.

Following Hutchinson [19], we say that V is an *oriented 2-varifold* on \mathbb{R}^3 , if and only if V is a Radon measure over $\mathbb{G}_2^\circ(\mathbb{R}^3) := \mathbb{R}^3 \times \mathbb{G}^\circ(3, 2)$. The *weight measure* μ_V on \mathbb{R}^3 is defined by

$$\mu_V(A) := V\{(x, \xi) \in \mathbb{G}_2^\circ(\mathbb{R}^3) \mid x \in A\} \quad \text{whenever } A \subset \mathbb{R}^3.$$

It is the push forward of V under the projection $\mathbb{R}^3 \times \mathbb{G}^\circ(3, 2) \rightarrow \mathbb{R}^3$. The set of oriented 2-varifolds in \mathbb{R}^3 is denoted by $\mathbb{V}_2^\circ(\mathbb{R}^3)$.

For each $\xi \in \mathbb{G}^\circ(3, 2)$ we define the unoriented 2-dimensional subspace $T(\xi)$ of \mathbb{R}^3 by

$$T(\xi) := \{v \in \mathbb{R}^3 \mid v \wedge \xi = 0\}.$$

Since ξ is simple, there exist $v_1, v_2 \in \mathbb{R}^3$ such that $\xi = v_1 \wedge v_2$. Moreover, $|\xi| = 1$ implies that $v_1 \wedge v_2 = e_1 \wedge e_2$ for $e_1 := v_1/|v_1|$ and $e_2 := \tilde{v}_2/|\tilde{v}_2|$ where $\tilde{v}_2 := v_2 - \langle e_1, v_2 \rangle e_1$. In other words, each $\xi \in \mathbb{G}^\circ(3, 2)$ corresponds to an oriented orthonormal basis (e_1, e_2) with $\xi = e_1 \wedge e_2$ and $T(\xi) = \text{span}\{e_1, e_2\}$. In particular, each oriented 2-varifold $V \in \mathbb{V}_2^\circ(\mathbb{R}^3)$ induces a general (unoriented) 2-varifold in the sense of [1, Definition 3.1], given by the push forward of V under the map $q(x, \xi) := (x, T(\xi))$. Notice that the two weight measures of V and $q\#V$ coincide.

For all compactly supported vector fields $X \in C_c^1(\mathbb{R}^3; \mathbb{R}^3)$ and 2-dimensional subspaces T of \mathbb{R}^3 with orthonormal basis $\{e_1, e_2\}$, we define

$$\operatorname{div}_T X(x) := \sum_{j=1}^2 \langle e_j, DX(x)e_j \rangle.$$

The *first variation* of an oriented 2-varifold V in \mathbb{R}^3 is defined by

$$\delta V: C_c^1(\mathbb{R}^3; \mathbb{R}^3) \rightarrow \mathbb{R}, \quad \delta V(X) := \int_{\mathbb{G}_2^o(\mathbb{R}^3)} \operatorname{div}_{T(\xi)} X(x) \, dV(x, \xi). \quad (2.1)$$

Notice that δV coincides with the first variation of the unoriented 2-varifold $q_{\#}V$ as defined in [1, Definition 4.2]. In other words, δV does not depend on the orientation. The singular part of the total variation of δV with respect to μ_V will be denoted by β_V (i.e. $\beta_V = \|\delta V\|_{\text{sing}}$, cf. [1, 4.3]).

For $k = 0, \dots, 3$ we may identify k -vectors in \mathbb{R}^3 with $(3 - k)$ -vectors by means of the Hodge star operator

$$\star: \bigwedge_k \mathbb{R}^3 \rightarrow \bigwedge_{3-k} \mathbb{R}^3.$$

If $v_1, v_2 \in \mathbb{R}^3 \simeq \bigwedge_1 \mathbb{R}^3$, we have $\star(v_1 \wedge v_2) = v_1 \times v_2$, where \times denotes the usual cross product on \mathbb{R}^3 . In particular, for all $\xi \in \mathbb{G}^o(3, 2)$ there holds $|\star \xi| = 1$. Moreover, we have $\star \star v = v$ for all $v \in \mathbb{R}^3 \simeq \bigwedge_1 \mathbb{R}^3$.

In the case where V has locally bounded first variation, δV can be represented by integration as follows (see [1, 4.3(5)]).

Hypothesis 2.1. *Let $V \in \mathbb{V}_2^o(\mathbb{R}^3)$, $\eta \in L^\infty(\beta_V; \mathbb{S}^2)$, and $H \in L_{\text{loc}}^1(\mu_V; \mathbb{R}^3)$. Suppose*

$$\delta V(X) = - \int_{\mathbb{R}^3} \langle X, H \rangle \, d\mu_V + \int_{\mathbb{R}^3} \langle X, \eta \rangle \, d\beta_V \quad (2.2)$$

and

$$H(x) \wedge \star \xi = 0 \quad \text{for } V\text{-almost all } (x, \xi). \quad (2.3)$$

The map H is often referred to as *generalized mean curvature* and $\text{spt } \beta_V$ can be seen as *generalized boundary*. Indeed, one can understand β_V as the boundary measure. However, two boundary parts can fall together as in Figure (7). Typically, one can determine H and β_V using Remark 4.4 and 4.7 in [1]. If V is rectifiable (i.e. $q_{\#}V$ is rectifiable in the sense of [1, 3.5]) then the condition in (2.3) means that the generalized mean curvature is perpendicular. In the absence of the singular part, Hypothesis 2.1 simplifies as follows.

Hypothesis 2.2. *Let $V \in \mathbb{V}_2^o(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^1(\mu_V; \mathbb{R}^3)$. Suppose*

$$\delta V(X) = - \int_{\mathbb{R}^3} \langle X, H \rangle \, d\mu_V \quad (2.4)$$

and

$$H(x) \wedge \star \xi = 0 \quad \text{for } V\text{-almost all } (x, \xi).$$

Let c_0 be a real number, and assume $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^1(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.1 for some $\eta \in L^\infty(\beta_V; \mathbb{S}^2)$. Then we define the *Helfrich energy*

$$\mathcal{H}_{c_0}(V) := \frac{1}{4} \int |H(x) - c_0(\star\xi)|^2 dV(x, \xi) = \frac{1}{4} \int (\langle H(x), \star\xi \rangle - c_0)^2 dV(x, \xi).$$

Notice that the Helfrich energy does not depend on the singular part of the first variation. This is analogous to the definition in [12, Section 2]. For $c_0 = 0$, we obtain the *Willmore functional* $\mathcal{W} := \mathcal{H}_0$.

Remark 2.3. Since μ_V and V are Radon measures, we have $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ if and only if the function $(x, \xi) \mapsto H(x) - c_0(\star\xi)$ is a member of $L_{\text{loc}}^2(V; \mathbb{R}^3)$. Indeed, given any Borel set B in \mathbb{R}^3 , the Cauchy–Schwarz inequality implies

$$\int_{B \times \mathbb{G}^0(3,2)} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) \leq 2 \int_B |H|^2 d\mu_V + 2c_0^2 \mu_V(B).$$

On the other hand,

$$\int_B |H|^2 d\mu_V \leq \int_{B \times \mathbb{G}^0(3,2)} 2|H(x) - c_0(\star\xi)|^2 dV(x, \xi) + 2c_0^2 \mu_V(B). \quad (2.5)$$

In particular, $\mathcal{H}_{c_0}(V) < \infty$ implies $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$.

2.3 Oriented varifolds induced by immersions

A particular class of oriented varifolds will be given by immersions of oriented surfaces. Following [33], we term a surface Σ to be *orientable*, if there exists an atlas $A = \{(U_\alpha, x_\alpha)\}_{\alpha \in I}$ such that the Jacobians $\det D(x_{\alpha_1} \circ x_{\alpha_2}^{-1})$ of all coordinate transformations are positive. The members of A are called *positive charts*. If $f: \Sigma \rightarrow \mathbb{R}^3$ is a smooth immersion, then we define the induced smooth normal field n along f (the *Gauss map*) by

$$n: \Sigma \rightarrow \mathbb{S}^2, \quad n := \frac{\partial_{x^1} f \times \partial_{x^2} f}{|\partial_{x^1} f \times \partial_{x^2} f|}, \quad (2.6)$$

whenever x is a positive chart. Notice that since the Hodge star operator is an isometry, $\star n = \partial_{x^1} f \wedge \partial_{x^2} f / |\partial_{x^1} f \wedge \partial_{x^2} f|$ takes values in $\mathbb{G}^0(3, 2)$. Moreover, in the context of an immersion f , we will always denote by $\mu = \mu_f$ the Riemannian measure induced by the pullback metric $g = g_f := f^*\langle \cdot, \cdot \rangle$, and we define by

$$\mathcal{A}(f) := \int_\Sigma 1 d\mu, \quad \mathcal{V}(f) := -\frac{1}{3} \int_\Sigma \langle f, n \rangle d\mu$$

the area and the (algebraic) volume of f , provided the respective integral exists. If f is an embedding and n is the inner unit normal, $\mathcal{V}(f)$ yields the enclosed volume as a consequence of the divergence theorem, see [31, Appendix A] for a more detailed discussion.

In the sequel, the immersion under consideration will usually be clear from the context, so we will drop the dependence on f of the associated geometric quantities.

Example 2.4 (Oriented varifold associated with immersed surface). Let $f: \Sigma \rightarrow \mathbb{R}^3$ be a smooth proper immersion of an oriented surface Σ without boundary. We define the oriented 2-varifold $V \in \mathbb{V}_2^{\circ}(\mathbb{R}^3)$ associated with (Σ, f) by

$$V(A) := \mu\{p \in \Sigma \mid (f(p), \star n(p)) \in A\} \quad \text{whenever } A \subset \mathbb{G}_2^{\circ}(\mathbb{R}^3),$$

i.e. V is the push forward of μ under the map $\Sigma \rightarrow \mathbb{R}^3 \times \mathbb{G}^{\circ}(3, 2), p \mapsto (f(p), \star n(p))$. Since this map is continuous and proper, V is indeed a Radon measure (see [16, 2.2.17]). Notice that $T(\star n(p)) = df_p[T_p\Sigma]$ for $p \in \Sigma$. In view of [36, Lemma 2.3], there holds

$$\begin{aligned} \mu_V(B) &= (f_{\#}\mu)(B) = \int_B \mathcal{H}^0(f^{-1}\{x\}) d\mathcal{H}^2(x) \quad \text{for all Borel sets } B \text{ in } \mathbb{R}^3, \\ \theta^2(\mu_V, x) &= \mathcal{H}^0(f^{-1}\{x\}) \quad \text{for all } x \in \mathbb{R}^3. \end{aligned}$$

Moreover, by [16, 2.4.18] and [36, Lemma 2.3] we have

$$\begin{aligned} \int_{\mathbb{G}_2^{\circ}(\mathbb{R}^3)} k(x, \xi) dV(x, \xi) &= \int_{\Sigma} k(f(p), \star n(p)) d\mu(p) \\ &= \int_{\mathbb{R}^3} \sum_{p \in f^{-1}\{x\}} k(x, \star n(p)) d\mathcal{H}^2(x) \end{aligned} \quad (2.7)$$

whenever $k: \mathbb{G}_2^{\circ}(\mathbb{R}^3) \rightarrow \mathbb{R}$ is a nonnegative Borel function.

Let $H_f: \Sigma \rightarrow \mathbb{R}^3$ be the classical mean curvature (vector) of f , i.e. the trace of the second fundamental form, and define

$$H(x) := \begin{cases} \frac{1}{\theta^2(\mu_V, x)} \sum_{p \in f^{-1}\{x\}} H_f(p) & \text{if } \theta^2(\mu_V, x) > 0 \\ 0 & \text{if } \theta^2(\mu_V, x) = 0. \end{cases} \quad (2.8)$$

Then, $H \in L_{\text{loc}}^{\infty}(\mu_V; \mathbb{R}^3)$ and in view of [36, Example 2.4], $H(x) \wedge \star \xi = 0$ for V -almost all (x, ξ) , and

$$\delta V(X) = - \int_{\mathbb{R}^3} \langle X, H \rangle d\mu_V.$$

Thus, V, H satisfy Hypothesis 2.2.

In the sequel, we will always use the above notation to distinguish H_f as the classical mean curvature when f is an immersion and H defined by (2.8) as the generalized mean curvature of the associated varifold. By [37, Theorem 4], there holds

$$H_f(p) = H(f(p)) \quad \text{for } \mu\text{-almost all } p \in \Sigma.$$

Thus, by (2.7) we observe

$$\mathcal{H}_{c_0}(V) = \frac{1}{4} \int_{\Sigma} |H_f - c_0 n|^2 d\mu. \quad (2.9)$$

3 On the concentrated volume

In this section, we discuss the concentrated volume (1.7) in the context of varifolds.

Definition 3.1. Suppose $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and $x_0 \in \mathbb{R}^3$. Then we define the *concentrated volume of V at x_0* by

$$\mathcal{V}_c(V, x_0) := - \int_{\mathbb{G}_2^0(\mathbb{R}^3)} \frac{\langle x - x_0, \star \xi \rangle}{|x - x_0|^2} dV(x, \xi)$$

and the *algebraic volume at x_0*

$$\mathcal{V}(V, x_0) := -\frac{1}{3} \int_{\mathbb{G}_2^0(\mathbb{R}^3)} \langle x - x_0, \star \xi \rangle dV(x, \xi)$$

provided the respective integral exists.

If the varifold V is associated with an immersion $f: \Sigma \rightarrow \mathbb{R}^3$, then we also write $\mathcal{V}_c(f, x_0)$ instead of $\mathcal{V}_c(V, x_0)$. By (2.7), this is consistent with (1.7). If Σ is closed, then we have $\mathcal{V}(V, x_0) = \mathcal{V}(f)$ for all $x_0 \in \mathbb{R}^3$ after integration by parts.

In general, the algebraic volume of an oriented varifold depends on the point x_0 . Indeed, one may consider the varifold associated to the 2-dimensional unit sphere in \mathbb{R}^3 where the upper hemisphere is oppositely oriented to the lower hemisphere.

Lemma 3.2. Suppose $m, \rho_0, D > 0$, μ is a Radon measure over \mathbb{R}^3 , $x_0 \in \mathbb{R}^3$, and

$$\mu(B_\sigma(x_0)) \leq D\sigma^m$$

for all $0 < \sigma < \rho_0$. Then, for all $1 \leq p < m$, there exists $C(p, m, D) < \infty$ such that

$$\int_{B_\sigma(x_0)} \frac{1}{|x - x_0|^p} d\mu(x) \leq C(p, m, D)\sigma^{m-p}$$

for all $0 < \sigma < \rho_0$. Moreover, $C(1, 2, D) = 2D$, and if $\mu(\mathbb{R}^3) < \infty$ then

$$\int_{\mathbb{R}^3} \frac{1}{|x - x_0|^p} d\mu(x) \leq C(p, m, D)\rho_0^{m-p} + \frac{\mu(\mathbb{R}^3)}{\rho_0^p} < \infty.$$

Remark 3.3. (i) Assume $V \in \mathbb{V}_2^0(\mathbb{R}^3)$, $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2. By [36, Theorem 3.6], we find that the density $\theta^2(\mu_V, x_0)$ exists and is finite for all $x_0 \in \mathbb{R}^3$. Hence there exist $\rho_0 = \rho_0(V, x_0) > 0$ and $D = D(V, x_0) < \infty$ such that

$$\mu_V(B_\sigma(x_0)) \leq D\sigma^2 \quad \text{for all } 0 < \sigma < \rho_0, x_0 \in \mathbb{R}^3.$$

This immediately implies $\mu_V\{x_0\} = 0$ for all $x_0 \in \mathbb{R}^3$. By Remark 2.3, the condition $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ is in particular satisfied if $\mathcal{H}_{c_0}(V) < \infty$.

(ii) If $V \in \mathbb{V}_2^0(\mathbb{R}^3)$, $\mu_V(\mathbb{R}^3) < \infty$, and $H \in L^2(\mu_V; \mathbb{R}^3)$ satisfies Hypothesis 2.2, then the hypothesis of Lemma 3.2 is satisfied for $m = 2$, $\rho_0 = \infty$, and all $x_0 \in \mathbb{R}^3$ with $D = CW(V)$ for some universal constant $0 < C < \infty$. Indeed, by [22, Appendix (A.16)] there holds

$$\mu_V(B_\sigma(x_0)) \leq CW(V)\sigma^2 \quad \text{for all } \sigma > 0, x_0 \in \mathbb{R}^3.$$

Proof of Lemma 3.2. We apply [24, Theorem 1.15] to compute

$$\begin{aligned} \int_{B_\sigma(x_0)} \frac{1}{|x - x_0|^p} d\mu(x) &= \int_0^\infty \mu(B_\sigma(x_0) \cap B_{t^{-1/p}}(x_0)) dt \\ &= \int_0^{\sigma^{-p}} \mu(B_\sigma(x_0)) dt + \int_{\sigma^{-p}}^\infty \mu(B_{t^{-1/p}}(x_0)) dt \\ &\leq D \left(\int_0^{\sigma^{-p}} \sigma^m dt + \int_{\sigma^{-p}}^\infty t^{-m/p} dt \right) \\ &= D \left(\sigma^{m-p} + \frac{p\sigma^{m-p}}{m-p} \right) = C(p, m, D)\sigma^{m-p}. \end{aligned}$$

The last statement follows by splitting the integral into $\mathbb{R}^3 = B_{\rho_0}(x_0) \cup (\mathbb{R}^3 \setminus B_{\rho_0}(x_0))$. \square

Proposition 3.4. *Suppose $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2 and assume that $\mathcal{V}(V, x_0)$ exists for some $x_0 \in \mathbb{R}^3$. Then also $\mathcal{V}_c(V, x_0)$ exists.*

Proof. Splitting the integral, for $\rho_0 > 0$ we find

$$\begin{aligned} &\int_{\mathbb{G}_2^0(\mathbb{R}^3)} \frac{|\langle x - x_0, \star\xi \rangle|}{|x - x_0|^2} dV(x, \xi) \\ &\leq \int_{B_{\rho_0}(x_0)} \frac{1}{|x - x_0|} d\mu_V(x) + \frac{1}{\rho_0^2} \int_{\mathbb{G}_2^0(\mathbb{R}^3)} |\langle x - x_0, \star\xi \rangle| dV(x, \xi). \end{aligned}$$

By Lemma 3.2 and Remark 3.3(i), on the right hand side the first integral is finite for $\rho_0 > 0$ small, whereas the second integral is finite since $\mathcal{V}(V, x_0)$ exists. \square

We recall the concept of convergence of oriented varifolds.

Definition 3.5. Suppose V_k is a sequence in $\mathbb{V}_2^0(\mathbb{R}^3)$. Then we say that V_k converges to V in $\mathbb{V}_2^0(\mathbb{R}^3)$ and write

$$V_k \rightarrow V \quad \text{in } \mathbb{V}_2^0(\mathbb{R}^3) \text{ as } k \rightarrow \infty$$

if and only if $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and

$$\int_{\mathbb{G}_2^0(\mathbb{R}^3)} \varphi(x, \xi) dV_k(x, \xi) \rightarrow \int_{\mathbb{G}_2^0(\mathbb{R}^3)} \varphi(x, \xi) dV(x, \xi) \quad \text{as } k \rightarrow \infty$$

for all continuous functions $\varphi: \mathbb{G}_2^0(\mathbb{R}^3) \rightarrow \mathbb{R}$ with compact support.

Lemma 3.6. *Suppose V_k is a sequence in $\mathbb{V}_2^o(\mathbb{R}^3)$, $V \in \mathbb{V}_2^o(\mathbb{R}^3)$, $H_k \in L^2(\mu_{V_k}; \mathbb{R}^3)$ and $H \in L^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2,*

$$\mu_V(\mathbb{R}^3) + \mathcal{W}(V) + \sup_{k \in \mathbb{N}} (\mu_{V_k}(\mathbb{R}^3) + \mathcal{W}(V_k)) < \infty \quad (3.1)$$

and

$$V_k \rightarrow V \quad \text{in } \mathbb{V}_2^o(\mathbb{R}^3) \text{ as } k \rightarrow \infty. \quad (3.2)$$

Then for all $x_0 \in \mathbb{R}^3$, the concentrated volume converges: $\lim_{k \rightarrow \infty} \mathcal{V}_c(V_k, x_0) = \mathcal{V}_c(V, x_0)$.

Proof. Let $x_0 \in \mathbb{R}^3$, $0 < \sigma < \rho < \infty$, and pick a continuous function $\chi: \mathbb{R}^3 \rightarrow \mathbb{R}$ with compact support in $\mathbb{R}^3 \setminus \{x_0\}$ such that $0 \leq \chi \leq 1$ and $\chi(x) = 1$ for $\sigma \leq |x - x_0| \leq \rho$. Define the function

$$\varphi: \mathbb{R}^3 \times \mathbb{G}^o(3, 2) \rightarrow \mathbb{R}, \quad \varphi(x, \xi) := \chi(x) \frac{\langle x - x_0, \star \xi \rangle}{|x - x_0|^2}.$$

Then φ has compact support, φ is continuous, and thus, by (3.2),

$$\int_{\mathbb{G}_2^o(\mathbb{R}^3)} \varphi \, dV_k \rightarrow \int_{\mathbb{G}_2^o(\mathbb{R}^3)} \varphi \, dV \quad \text{as } k \rightarrow \infty. \quad (3.3)$$

Let

$$A := \sup_{k \in \mathbb{N}} (\mu_{V_k}(\mathbb{R}^3) + \mu_V(\mathbb{R}^3)), \quad D := \sup_{k \in \mathbb{N}} (\mathcal{W}(V_k) + \mathcal{W}(V)).$$

Then, by (3.1), Lemma 3.2, and Remark 3.3(ii), we have

$$\sup_{k \in \mathbb{N}} \left(\int_{B_\sigma(x_0)} \frac{1}{|x - x_0|} \, d\mu_{V_k}(x) + \int_{B_\sigma(x_0)} \frac{1}{|x - x_0|} \, d\mu_V(x) \right) \leq C(D)\sigma$$

and

$$\sup_{k \in \mathbb{N}} \left(\int_{\mathbb{R}^3 \setminus B_\rho(x_0)} \frac{1}{|x - x_0|} \, d\mu_{V_k}(x) + \int_{\mathbb{R}^3 \setminus B_\rho(x_0)} \frac{1}{|x - x_0|} \, d\mu_V(x) \right) \leq \frac{C(A)}{\rho}.$$

Since $|\varphi(x, \xi)| \leq 1/|x - x_0|$ for all $(x, \xi) \in \mathbb{G}_2^o(\mathbb{R}^3)$, it follows

$$|\mathcal{V}(V, x_0) - \mathcal{V}(V_k, x_0)| \leq \left| \int_{\mathbb{G}_2^o(\mathbb{R}^3)} \varphi \, dV_k - \int_{\mathbb{G}_2^o(\mathbb{R}^3)} \varphi \, dV \right| + C(D)\sigma + \frac{C(A)}{\rho}.$$

Now, the conclusion follows from the convergence in (3.3). \square

Lemma 3.7. *Suppose $V \in \mathbb{V}_2^o(\mathbb{R}^3)$, $\text{spt } \mu_V$ is compact, and $H \in L^2(\mu_V; \mathbb{R}^3)$ satisfies Hypothesis 2.2. Then the concentrated volume $\mathcal{V}_c(V, \cdot)$ is Hölder continuous with exponent α for any $0 < \alpha < 1$ and constant $C = C(\alpha, V)$ depending monotonically nondecreasing on $\mu_V(\mathbb{R}^3)$ and $\mathcal{W}(V)$.*

Proof. Let $0 < \alpha < 1$, $x_0, x_1 \in \mathbb{R}^3$ with $0 < |x_0 - x_1| \leq 1$, and abbreviate $\sigma := |x_0 - x_1|$, $A := \mu_V(\mathbb{R}^3)$ and $D := \mathcal{W}(V)$. By Lemma 3.2 and Remark 3.3(ii) there holds

$$\int_{B_{2\sigma}(x_0)} \frac{1}{|x - x_0|} d\mu_V(x) \leq 4D|x_0 - x_1| \leq C(D)|x_0 - x_1|^\alpha$$

and, since $B_{2\sigma}(x_0) \subset B_{3\sigma}(x_1)$,

$$\int_{B_{2\sigma}(x_0)} \frac{1}{|x - x_1|} d\mu_V(x) \leq \int_{B_{3\sigma}(x_1)} \frac{1}{|x - x_1|} d\mu_V(x) \leq 6D|x_0 - x_1| \leq C(D)|x_0 - x_1|^\alpha.$$

Thus, we have

$$\begin{aligned} & |\mathcal{V}_c(V, x_0) - \mathcal{V}_c(V, x_1)| \\ & \leq C(D)|x_0 - x_1|^\alpha + \int_{\pi^{-1}[\mathbb{R}^3 \setminus B_{2\sigma}(x_0)]} \left| \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} - \frac{\langle x - x_1, \star\xi \rangle}{|x - x_1|^2} \right| dV(x, \xi). \end{aligned} \quad (3.4)$$

where $\pi : \mathbb{R}^3 \times \mathbb{G}^\circ(3, 2) \rightarrow \mathbb{R}^3$ is the projection. For all $x \in \mathbb{R}^3 \setminus B_{2\sigma}(x_0)$ there holds $|x - x_1| \geq |x - x_0|/2$ as well as $|x_0 - x_1| \leq |x - x_0|$ and hence

$$\begin{aligned} & \left| \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} - \frac{\langle x - x_1, \star\xi \rangle}{|x - x_1|^2} \right| \\ & \leq \left| \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} - \frac{\langle x - x_1, \star\xi \rangle}{|x - x_0|^2} \right| + \left| \frac{\langle x - x_1, \star\xi \rangle}{|x - x_0|^2} - \frac{\langle x - x_1, \star\xi \rangle}{|x - x_0||x - x_1|} \right| \\ & \quad + \left| \frac{\langle x - x_1, \star\xi \rangle}{|x - x_0||x - x_1|} - \frac{\langle x - x_1, \star\xi \rangle}{|x - x_1|^2} \right| \\ & \leq 2 \frac{|x_0 - x_1|}{|x - x_0|^2} + \frac{|x_0 - x_1|}{|x - x_0||x - x_1|} \leq 4 \frac{|x_0 - x_1|^\alpha}{|x - x_0|^{1+\alpha}}. \end{aligned}$$

Integrating this inequality and applying Lemma 3.2 for $p = 1 + \alpha$, $m = 2$ and $\rho_0 = 1$, (3.4) becomes

$$\begin{aligned} |\mathcal{V}_c(V, x_0) - \mathcal{V}_c(V, x_1)| & \leq C(D)|x_0 - x_1|^\alpha + 4|x_0 - x_1|^\alpha \int_{\mathbb{R}^3} \frac{1}{|x - x_0|^{1+\alpha}} d\mu_V(x) \\ & \leq [C(D) + 4C(\alpha, D) + 4A]|x_0 - x_1|^\alpha = C(\alpha, A, D)|x_0 - x_1|^\alpha. \end{aligned}$$

For $|x_0 - x_1| \geq 1$ we apply Lemma 3.2 to see

$$|\mathcal{V}_c(V, x_0) - \mathcal{V}_c(V, x_1)| \leq 2(2D + A) \leq C(A, D)|x_0 - x_1|^\alpha$$

which concludes the proof. \square

Example 3.8. Consider $S := \partial B_1(0) \subset \mathbb{R}^3$, the round sphere with radius one centered at the origin. Moreover, for $r > 0$ let $T_r \subset \mathbb{R}^3$ be the torus which is obtained by revolving a circle with radius r and center $(1 + r, 0)$ (in the xz -plane) around the z -axis. Note that if we revolve the corresponding disk instead of the circle, we obtain a full torus T_r^{full}

with $\partial T_r^{\text{full}} = T_r$. We now define a smooth unit normal n on $S \cup T_r$ by taking n to be the outer unit normal on S and the inner unit normal on T_r , cf. Figure (9). It is not difficult to see that $S \cup T_r$ is the image of a $C^{1,1}$ -immersion $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$.

By standard formulas in geometry, the algebraic volume can be computed as

$$\mathcal{V}(f) = -\mathcal{L}^3(B_1(0)) + \mathcal{L}^3(T_r^{\text{full}}) = -\frac{4\pi}{3} + 2\pi^2 r^2(1+r). \quad (3.5)$$

Let $x_0 = 0$ be the origin. Using that $|x - x_0| = 1$ for $x \in S$ and applying the divergence theorem to T_r^{full} , the concentrated volume is given by

$$\mathcal{V}_c(f, x_0) = -\frac{4\pi}{3} + \int_{T_r^{\text{full}}} \frac{1}{|x - x_0|^2} d\mathcal{L}^3(x).$$

Clearly $|x - x_0| > 1$ for \mathcal{L}^3 -almost every $x \in T_r^{\text{full}}$, so that

$$\mathcal{V}_c(f, x_0) < -\frac{4\pi}{3} + \int_{T_r^{\text{full}}} d\mathcal{L}^3(x) = \mathcal{V}(f).$$

By means of (3.5), we thus find $r > 0$ such that $\mathcal{V}(f) = 0$ but $\mathcal{V}_c(f, x_0) < 0$. Slightly increasing the radius, we have $\mathcal{V}_c(f, x_0) < 0 < \mathcal{V}(f)$ by a continuity argument. Lastly, we replace a small disk on S by a thin dent, such that the new surface \tilde{S} satisfies $x_0 \in \tilde{S}$ and such that $\tilde{S} \cup T_r$ is still the image of an immersion $\tilde{f}: \mathbb{S}^2 \rightarrow \mathbb{R}^3$. Making this dent sufficiently thin and smoothing, we can achieve that $\mathcal{V}_c(\tilde{f}, x_0) < 0 < \mathcal{V}(\tilde{f})$ is still satisfied and \tilde{f} is smooth. Therefore, positive algebraic volume does not imply positive concentrated volume, even at points in the support.

4 The Li–Yau inequality in the varifold setting

4.1 A monotonicity formula

Our first essential observation is the following lemma, which can be seen as an extension of the monotonicity formula due to Simon [40, (1.2)]. We follow the varifold approach in [22, Appendix A] relying on the first variation identity and examine the additional terms originating from the spontaneous curvature.

Lemma 4.1. *Suppose $V \in \mathbb{V}_2^{\circ}(\mathbb{R}^3)$, $\eta \in L^\infty(\beta_V; \mathbb{S}^2)$, and $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.1. Let $x_0 \in \mathbb{R}^3$ and abbreviate $B_\rho := B_\rho(x_0)$, $A_\rho := B_\rho \times \mathbb{G}^\circ(3, 2)$ for all $\rho > 0$. Then, for $c_0 \in \mathbb{R}$ and $0 < \sigma < \rho < \infty$, there holds*

$$\begin{aligned} & \frac{\mu_V(B_\sigma)}{\sigma^2} + \int_{A_\rho \setminus A_\sigma} \left(\frac{1}{4} (\langle H(x), \star \xi \rangle - c_0) + \frac{\langle x - x_0, \star \xi \rangle}{|x - x_0|^2} \right)^2 dV(x, \xi) \\ &= \frac{1}{16} \int_{A_\rho \setminus A_\sigma} |H(x) - c_0(\star \xi)|^2 dV(x, \xi) - \frac{c_0}{2} \int_{A_\rho \setminus A_\sigma} \frac{\langle x - x_0, \star \xi \rangle}{|x - x_0|^2} dV(x, \xi) + \frac{\mu_V(B_\rho)}{\rho^2} \\ & \quad - \frac{1}{2\sigma^2} \int_{A_\sigma} \langle x - x_0, H(x) - c_0(\star \xi) \rangle dV(x, \xi) - \frac{c_0}{2\sigma^2} \int_{A_\sigma} \langle x - x_0, \star \xi \rangle dV(x, \xi) \end{aligned}$$

$$\begin{aligned}
& + \frac{1}{2\rho^2} \int_{A_\rho} \langle x - x_0, H(x) - c_0(\star\xi) \rangle dV(x, \xi) + \frac{c_0}{2\rho^2} \int_{A_\rho} \langle x - x_0, \star\xi \rangle dV(x, \xi) \\
& + \frac{1}{2\sigma^2} \int_{B_\sigma} \langle x - x_0, \eta(x) \rangle d\beta_V(x) + \frac{1}{2} \int_{B_\rho \setminus B_\sigma} \frac{\langle x - x_0, \eta(x) \rangle}{|x - x_0|^2} d\beta_V(x) \\
& - \frac{1}{2\rho^2} \int_{B_\rho} \langle x - x_0, \eta(x) \rangle d\beta_V(x). \tag{4.1}
\end{aligned}$$

Proof of Lemma 4.1. Following the computations in [40, p. 284], we consider the smooth vector field $X(x) := x - x_0$ for $x \in \mathbb{R}^3$ and the Lipschitz function

$$\varphi: \mathbb{R} \rightarrow \mathbb{R}, \quad \varphi(t) := (\max\{t, \sigma\}^{-2} - \rho^{-2})_+. \tag{4.2}$$

Choose a sequence φ_k in $C_c^\infty(-\infty, \rho + 1)$ such that $\sup_{k \in \mathbb{N}} \|\varphi_k\|_{C^1(\mathbb{R})} < \infty$,

$$\begin{aligned}
\varphi_k &\rightarrow \varphi && \text{in } C^0(\mathbb{R}) \text{ as } k \rightarrow \infty, \\
\varphi'_k(t) &\rightarrow \varphi'(t) && \text{as } k \rightarrow \infty \text{ for all } t \in \mathbb{R} \setminus \{\sigma, \rho\}
\end{aligned}$$

and such that for all $k \in \mathbb{N}$, there holds $\varphi'_k(\sigma) = 0$ and $\varphi'_k(\rho) = -2\rho^{-3}$. Abbreviating $\Phi_k := \varphi_k \circ |X|$ it follows

$$\lim_{k \rightarrow \infty} \operatorname{div}_{T(\xi)}(\Phi_k X)(x) = \begin{cases} 2(\frac{1}{\sigma^2} - \frac{1}{\rho^2}) & \text{for } (x, \xi) \in A_\sigma \\ \frac{2\langle X(x), \star\xi \rangle^2}{|X(x)|^4} - \frac{2}{\rho^2} & \text{for } (x, \xi) \in A_\rho \setminus A_\sigma \\ 0 & \text{for } (x, \xi) \in \mathbb{G}_2^o(\mathbb{R}^3) \setminus A_\rho. \end{cases}$$

Denoting $|X|_\sigma := \max\{|X|, \sigma\}$, testing the first variation identity (see (2.2), and (2.1)) with the vector fields $\Phi_k X$ and passing to the limit as $k \rightarrow \infty$, we obtain

$$\begin{aligned}
& \frac{2\mu_V(B_\sigma)}{\sigma^2} + \int_{A_\rho \setminus A_\sigma} \frac{2\langle X(x), \star\xi \rangle^2}{|X(x)|^4} dV(x, \xi) \\
& = \frac{2\mu_V(B_\rho)}{\rho^2} - \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, H \rangle d\mu_V + \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, \eta \rangle d\beta_V.
\end{aligned}$$

By (2.3) and since $|\star\xi| = 1$ for $\xi \in \mathbb{G}^o(3, 2)$, we have the pointwise identity

$$\left| \frac{1}{4}(H - c_0(\star\xi)) + \frac{\langle X, \star\xi \rangle \langle \star\xi \rangle}{|X|^2} \right|^2 = \frac{1}{16}|H - c_0(\star\xi)|^2 + \frac{\langle H - c_0(\star\xi), X \rangle}{2|X|^2} + \frac{\langle X, \star\xi \rangle^2}{|X|^4}$$

and consequently

$$\begin{aligned}
& \frac{\mu_V(B_\sigma)}{\sigma^2} + \int_{A_\rho \setminus A_\sigma} \left(\frac{1}{4}(\langle H(x), \star\xi \rangle - c_0) + \frac{\langle X(x), \star\xi \rangle}{|X(x)|^2} \right)^2 dV(x, \xi) \\
& = \frac{1}{16} \int_{A_\rho \setminus A_\sigma} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) + \frac{\mu_V(B_\rho)}{\rho^2} \\
& + \frac{1}{2} \int_{A_\rho \setminus A_\sigma} \frac{\langle H(x) - c_0(\star\xi), X(x) \rangle}{|X(x)|^2} dV(x, \xi) - \frac{1}{2} \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, H \rangle d\mu_V
\end{aligned}$$

$$+ \frac{1}{2} \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, \eta \rangle d\beta_V. \quad (4.3)$$

It follows

$$\begin{aligned} & \frac{1}{2} \int_{A_\rho \setminus A_\sigma} \frac{\langle H(x) - c_0(\star\xi), X(x) \rangle}{|X(x)|^2} dV(x, \xi) - \frac{1}{2} \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, H \rangle d\mu_V \\ &= -\frac{c_0}{2} \int_{A_\rho \setminus A_\sigma} \frac{\langle X(x), \star\xi \rangle}{|X(x)|^2} dV(x, \xi) - \frac{1}{2\sigma^2} \int_{B_\sigma} \langle X, H \rangle d\mu_V + \frac{1}{2\rho^2} \int_{B_\rho} \langle X, H \rangle d\mu_V \\ &= -\frac{c_0}{2} \int_{A_\rho \setminus A_\sigma} \frac{\langle X(x), \star\xi \rangle}{|X(x)|^2} dV(x, \xi) \\ &\quad - \frac{1}{2\sigma^2} \int_{A_\sigma} \langle X(x), H(x) - c_0(\star\xi) \rangle dV(x, \xi) - \frac{c_0}{2\sigma^2} \int_{A_\sigma} \langle X(x), \star\xi \rangle dV(x, \xi) \\ &\quad + \frac{1}{2\rho^2} \int_{A_\rho} \langle X(x), H(x) - c_0(\star\xi) \rangle dV(x, \xi) + \frac{c_0}{2\rho^2} \int_{A_\rho} \langle X(x), \star\xi \rangle dV(x, \xi) \end{aligned} \quad (4.4)$$

as well as

$$\begin{aligned} & \frac{1}{2} \int_{B_\rho} (|X|_\sigma^{-2} - \rho^{-2}) \langle X, \eta \rangle d\beta_V \\ &= \frac{1}{2\sigma^2} \int_{B_\sigma} \langle X, \eta \rangle d\beta_V + \frac{1}{2} \int_{B_\rho \setminus B_\sigma} \frac{\langle X, \eta \rangle}{|X|^2} d\beta_V - \frac{1}{2\rho^2} \int_{B_\rho} \langle X, \eta \rangle d\beta_V. \end{aligned} \quad (4.5)$$

Now, using $X(x) = x - x_0$ and putting (4.4), (4.5) into (4.3), the conclusion follows. \square

4.2 The general varifold case

We now use the monotonicity formula (4.1) to prove our most general Li–Yau inequality.

Theorem 4.2. *Suppose $V \in \mathbb{V}_2^0(\mathbb{R}^3)$, $\eta \in L^\infty(\beta_V; \mathbb{S}^2)$ and $H \in L^1_{\text{loc}}(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.1. Let $c_0 \in \mathbb{R}$ and suppose that*

$$\mathcal{H}_{c_0}(V) < \infty \quad (4.6)$$

and

$$\theta^{*2}(\mu_V, \infty) := \limsup_{\rho \rightarrow \infty} \frac{\mu_V(B_\rho(0))}{\pi\rho^2} < \infty. \quad (4.7)$$

Then, for all $x_0 \in \mathbb{R}^3 \setminus \text{spt } \beta_V$ we have

$$\begin{aligned} \theta^2(\mu_V, x_0) &\leq \theta^{*2}(\mu_V, \infty) + \frac{1}{4\pi} \mathcal{H}_{c_0}(V) \\ &\quad + \limsup_{\rho \rightarrow \infty} \frac{c_0}{2\pi} \left(\int_{B_\rho(x_0) \times \mathbb{G}_2^2(\mathbb{R}^3)} (\rho^{-2} - |x - x_0|^{-2}) \langle x - x_0, \star\xi \rangle dV(x, \xi) \right) \\ &\quad + \limsup_{\rho \rightarrow \infty} \frac{1}{2\pi} \left(\int_{B_\rho(x_0)} (|x - x_0|^{-2} - \rho^{-2}) \langle x - x_0, \eta(x) \rangle d\beta_V(x) \right). \end{aligned} \quad (4.8)$$

Remark 4.3. (i) We do not assume $\mu_V(\mathbb{R}^3) < \infty$ in Theorem 4.2. Indeed, let $r = 1/c_0$ for $c_0 > 0$, let $f: \mathbb{R} \times \mathbb{S}^1 \rightarrow \mathbb{R}^3$, $f(t, \varphi) = (r \cos \varphi, r \sin \varphi, t)$ be the cylinder with radius r , let V be the associated varifold, cf. Example 2.4, and let $x_0 = (r, 0, 0) \in \text{spt } \mu_V$. It is not difficult to see that $\beta = 0$, $\mathcal{H}_{c_0}(V) = 0$ and $\mu_V(B_\rho(x_0)) = O(\rho)$ as $\rho \rightarrow \infty$, so that $\theta^{*2}(\mu_V, \infty) = 0$ whereas $\mu_V(\mathbb{R}^3) = \infty$. Moreover, the third term on the right hand side of (4.8) can be computed to be $4\pi r$.

(ii) We can reverse the orientation of the varifold V by considering \hat{V} , the push forward under the map $(x, \xi) \mapsto (x, -\xi)$, which is continuous and proper so $\hat{V} \in \mathbb{V}_2^o(\mathbb{R}^3)$ by [16, 2.2.17]. In view of (1.1) it is not surprising that

$$\mathcal{H}_{c_0}(V) = \mathcal{H}_{-c_0}(\hat{V}).$$

Similarly, the other term in (4.8) involving c_0 remains unchanged if we replace V by \hat{V} and c_0 by $-c_0$. The singular part does not change under reversing the orientation.

(iii) Equality holds for $c_0 = 0$ if V corresponds to the unit sphere and x_0 is any point on the unit sphere. Equality also holds for $c_0 = 0$ if V corresponds to the unit disk and x_0 is the center, and if V corresponds to a plane and x_0 is any point on the plane.

(iv) If the singular part β_V is regular enough, for instance if $\text{spt } \beta_V$ is given by a smooth embedding $\gamma: \mathbb{S}^1 \rightarrow \mathbb{R}^3$ and $\eta \circ \gamma$ is a unit normal field along γ , then the statement remains valid even for $x_0 \in \text{spt } \beta_V$. Indeed, for x close to x_0 , the vectors $x - x_0$ and $\eta(x)$ are nearly orthogonal. Thus, since $\theta^1(\beta_V, x_0) = 1$, a short argument using Taylor expansion of γ implies

$$x \mapsto |x - x_0|^{-2} \langle x - x_0, \eta(x) \rangle \in L_{\text{loc}}^1(\beta_V).$$

Proof of Theorem 4.2. For $\rho > 0$ let B_ρ and A_ρ be as in Lemma 4.1. By Remark 3.3(i), there exist $D < \infty$ and $\rho_0 > 0$ such that

$$\mu(B_\rho) \leq D\rho^2 \quad \text{for all } 0 < \rho < \rho_0. \quad (4.9)$$

Consequently, Lemma 3.2 yields

$$\int_{B_\rho} \frac{1}{|x - x_0|} d\mu_V(x) \leq C\rho \quad \text{for all } 0 < \rho < \rho_0, \quad (4.10)$$

and thus $x \mapsto |x - x_0|^{-1} \in L_{\text{loc}}^1(\mu_V)$. Moreover we have $\text{dist}(x_0, \text{spt } \beta_V) > 0$, and consequently

$$x \mapsto |x - x_0|^{-2} \langle x - x_0, \eta(x) \rangle \in L_{\text{loc}}^1(\beta_V). \quad (4.11)$$

Using (4.6), (4.10) and (4.11), we find that the function $\gamma: (0, \infty) \rightarrow \mathbb{R}$ with

$$\gamma(\rho) := \frac{\mu_V(B_\rho)}{\rho^2} + \frac{1}{16} \int_{A_\rho} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) - \frac{c_0}{2} \int_{A_\rho} \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} dV(x, \xi)$$

$$\begin{aligned}
& + \frac{1}{2\rho^2} \int_{A_\rho} \langle x - x_0, H(x) - c_0(\star\xi) \rangle dV(x, \xi) + \frac{c_0}{2\rho^2} \int_{A_\rho} \langle x - x_0, \star\xi \rangle dV(x, \xi) \\
& + \frac{1}{2} \int_{B_\rho} (|x - x_0|^{-2} - \rho^{-2}) \langle x - x_0, \eta(x) \rangle d\beta_V(x)
\end{aligned} \tag{4.12}$$

is well defined and, by Lemma 4.1, it is monotonically nondecreasing.

We now examine the limit $\lim_{\sigma \rightarrow 0+} \gamma(\sigma)$. By (4.6), the second term in $\gamma(\sigma)$ goes to zero as $\sigma \rightarrow 0+$ and so does the third term by (4.10). For the fourth term, we use the Cauchy–Schwarz inequality to estimate

$$\begin{aligned}
& \left| \sigma^{-2} \int_{A_\sigma} \langle x - x_0, H(x) - c_0(\star\xi) \rangle dV(x, \xi) \right| \\
& \leq (\sigma^{-2} \mu_V(B_\sigma))^{\frac{1}{2}} \left(\int_{A_\sigma} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) \right)^{\frac{1}{2}},
\end{aligned} \tag{4.13}$$

where the right hand side goes to zero by (4.6), (4.9) and since $\mu_V\{x_0\} = 0$ by Remark 3.3(i). The fifth term in $\gamma(\sigma)$ also goes to zero as $\sigma \rightarrow 0+$, since

$$\sigma^{-2} \left| \int_{A_\sigma} \langle x - x_0, \star\xi \rangle dV(x, \xi) \right| \leq \sigma^{-1} \mu_V(B_\sigma) \leq D\sigma, \tag{4.14}$$

using (4.9). Since $x_0 \notin \text{spt } \beta_V$, we have $\beta_V(B_\sigma) = 0$ for $\sigma > 0$ sufficiently small. Consequently, using $\omega_2 = \pi$, we find $\lim_{\sigma \rightarrow 0+} \gamma(\sigma) = \pi\theta^2(\mu_V, x_0)$.

Now, we discuss the limit $\lim_{\rho \rightarrow \infty} \gamma(\rho)$. It is not too difficult to see that

$$\limsup_{\rho \rightarrow \infty} \frac{\mu_V(B_\rho)}{\pi\rho^2} = \limsup_{\rho \rightarrow \infty} \frac{\mu_V(B_\rho(0))}{\pi\rho^2} = \theta^{*2}(\mu_V, \infty). \tag{4.15}$$

For the fourth term in (4.12), for any $0 < \sigma < \rho$, we estimate by Cauchy–Schwarz

$$\begin{aligned}
& \left| \rho^{-2} \int_{A_\rho} \langle x - x_0, H(x) - c_0(\star\xi) \rangle dV(x, \xi) \right| \\
& \leq (\rho^{-2} \mu_V(B_\rho))^{\frac{1}{2}} \left(\int_{\mathbb{G}_2^o(\mathbb{R}^3) \setminus A_\sigma} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) \right)^{\frac{1}{2}} \\
& \quad + \rho^{-1} \int_{A_\sigma} |H(x) - c_0(\star\xi)| dV(x, \xi).
\end{aligned}$$

Sending first $\rho \rightarrow \infty$ and then $\sigma \rightarrow \infty$, this goes to zero by (4.6), (4.7) and (4.15). The claim then follows from the monotonicity of γ . \square

If the singular part β_V vanishes and $\mathcal{V}_c(V, x_0)$ exists, using Lemma 3.2 we obtain the following

Corollary 4.4. *Suppose $V \in \mathbb{V}_2^o(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2. Let $c_0 \in \mathbb{R}$, $x_0 \in \mathbb{R}^3$ and suppose that $\mathcal{V}_c(V, x_0)$ exists. Then we have*

$$\theta^2(\mu_V, x_0) \leq \theta^{*2}(\mu_V, \infty) + \frac{1}{4\pi} \mathcal{H}_{c_0}(V) + \frac{c_0}{2\pi} \mathcal{V}_c(V, x_0).$$

Proof. Without loss of generality, we may assume $\mathcal{H}_{c_0}(V) < \infty$, $\theta^{*2}(\mu_V, \infty) < \infty$. By Theorem 4.2, we only need to discuss the third term on the right hand side of (4.8). To that end, for $0 < \sigma < \rho$ we estimate

$$\begin{aligned} & \frac{1}{\rho^2} \int_{A_\rho} |\langle x - x_0, \star \xi \rangle| dV(x, \xi) \\ & \leq \int_{\mathbb{G}_2^o(\mathbb{R}^3) \setminus A_\sigma} \frac{|\langle x - x_0, \star \xi \rangle|}{|x - x_0|^2} dV(x, \xi) + \frac{1}{\rho^2} \int_{A_\sigma} |\langle x - x_0, \star \xi \rangle| dV(x, \xi). \end{aligned}$$

Sending first $\rho \rightarrow \infty$ and then $\sigma \rightarrow \infty$ this goes to zero since $\mathcal{V}_c(V, x_0)$ exists by assumption. The result follows. \square

4.3 Varifolds with enclosed volume

In this section we introduce a class of oriented varifolds that satisfy a divergence theorem, see Hypothesis 4.5. These varifolds comprise the surfaces shown in Figures (1)–(5) and (8). We then show that their concentrated volume is positive, see Lemma 4.9. We start with a short review of sets of locally finite perimeter.

Let $E \subset \mathbb{R}^3$. We define the *measure theoretic boundary of E* by

$$\partial_* E = \{x \in \mathbb{R}^3 \mid \theta^{*3}(\mathcal{L}^3 \llcorner E, x) > 0, \theta^{*3}(\mathcal{L}^3 \llcorner (\mathbb{R}^3 \setminus E), x) > 0\}.$$

Moreover, we denote with $n_E: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ the *measure theoretic inner unit normal of E* (see the definition [16, 4.5.5]). In view of *Federer's criterion* [16, 4.5.11, 2.10.6], we say that E has *locally finite perimeter*, if and only if E is an \mathcal{L}^3 -measurable set, and $\mathcal{H}^2(K \cap \partial_* E) < \infty$ for all compact sets $K \subset \mathbb{R}^3$.

Let $E \subset \mathbb{R}^3$ be a set of locally finite perimeter and $B = \{x \in \mathbb{R}^3 \mid n_E(x) \neq 0\}$. We collect the following properties (see [16, 4.5.6]).

- The sets B and $\partial_* E$ are \mathcal{H}^2 -almost equal.
- $\mathcal{H}^2 \llcorner \partial_* E$ is a Radon measure over \mathbb{R}^3 and n_E is $\mathcal{H}^2 \llcorner \partial_* E$ -measurable.
- The divergence theorem reads as

$$- \int_{\partial_* E} \langle X, n_E \rangle d\mathcal{H}^2 = \int_E \operatorname{div} X d\mathcal{L}^3 \quad (4.16)$$

for all Lipschitz maps $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ with compact support.

We define the oriented varifold $V \in \mathbb{V}_2^o(\mathbb{R}^3)$ associated with $\partial_* E$ as the push forward of the Radon measure $\mathcal{H}^2 \llcorner \partial_* E$ under the map

$$\mathbb{R}^3 \rightarrow \mathbb{R}^3 \times \mathbb{G}^o(3, 2), \quad x \mapsto (x, \star n_E(x)). \quad (4.17)$$

In view of [26, Lemma 2.6], V is indeed a Radon measure over $\mathbb{G}_2^o(\mathbb{R}^3)$. There holds $\mu_V = \mathcal{H}^2 \llcorner \partial_* E$ and (4.16) reads

$$\int_{\mathbb{G}_2^o(\mathbb{R}^3)} \langle X(x), \star \xi \rangle dV(x, \xi) = - \int_E \operatorname{div} X d\mathcal{L}^3$$

for all Lipschitz maps $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ with compact support.

This divergence theorem is the main motivation for considering a particular class of varifolds in the sequel.

Hypothesis 4.5. *Suppose $V \in \mathbb{V}_2^o(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^1(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2, $E \subset \mathbb{R}^3$ is an \mathcal{L}^3 -measurable set, $\Theta \in L_{\text{loc}}^1(\mathcal{L}^3 \llcorner E; \mathbb{N})$,*

$$\text{diam spt}(\mathcal{L}^3 \llcorner E) \leq \text{diam spt } \mu_V, \quad (4.18)$$

and

$$- \int_{\mathbb{G}_2^o(\mathbb{R}^3)} \langle X(x), \star \xi \rangle dV(x, \xi) = \int_E (\text{div } X) \Theta d\mathcal{L}^3 \quad (4.19)$$

for all Lipschitz maps $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ with compact support. In this case, we term V a varifold with enclosed volume.

Remark 4.6. (i) Since the divergence theorem (4.16) remains true if we replace E with $\mathbb{R}^3 \setminus E$ and n_E with $-n_E$, condition (4.18) ensures that we pick the correct orientation.

(ii) The function Θ has *locally bounded variation* (see the definition [15, Section 5.1]) and the coarea formula [16, 4.5.9(13)] implies that

$$E_k := \{x \in \mathbb{R}^3 \mid \Theta(x) \geq k\} \quad \text{for } k \in \mathbb{N}$$

defines a sequence of decreasing sets of locally finite perimeter.

(iii) If $\Theta \equiv 1$, then the varifold associated with $\partial_* E$ does not necessarily coincide with V , compare Figures (5) and (6).

(iv) If V is associated with the reduced boundary of a set E of locally finite perimeter, then $q\#V$ is an integral varifold (in the sense of [1, 3.5]). Hence, if additionally V has generalized mean curvature H and vanishing singular part $\beta_V = 0$, then there holds $H(x) \wedge \star \xi = 0$ for V -almost all (x, ξ) by [5, Section 5.8], V, H satisfy Hypothesis 2.2 and thus V, H, E and $\Theta \equiv 1$ satisfy Hypothesis 4.5.

As the following example shows, not all varifolds associated with sets of finite perimeter satisfy Hypothesis 4.5.

Example 4.7. Let C_α be the closed spherical cap of the unit sphere with opening angle $\alpha = \pi/3$ (the hemisphere has opening angle $\pi/2$) whose boundary circle lies in the plane $\{z = 0\}$. Let $S = C_\alpha \cup (-C_\alpha)$, i.e. S is the gluing of the spherical cap C_α with its reflection at the plane $\{z = 0\}$. The surface S looks like a lens, see Figure (7). Its singular part is the circle Γ_a of radius $a = \sqrt{3}/2$ centered at the origin and lying in the plane $\{z = 0\}$. Since $\mathcal{H}^2(S) < 4\pi < \infty$, one can use Federer's criterion to show that S is the boundary of a set E of finite perimeter. However, the varifold V associated with $S = \partial_* E$ (cf. (4.17)) does not satisfy Hypothesis 4.5. In fact, V does not satisfy

Hypothesis 2.2, but the more general Hypothesis 2.1. Indeed, in view of [1, 4.4, 4.7], there holds $\mu_V = \mathcal{H}^2 \llcorner S$, $\beta_V = \sqrt{3}\mathcal{H}^1 \llcorner \Gamma_a$, and

$$\delta V(X) = - \int_S \langle X, H \rangle d\mathcal{H}^2 + \sqrt{3} \int_{\Gamma_a} \frac{\langle X(x), x \rangle}{|x|} d\mathcal{H}^1(x)$$

where H is the mean curvature of the spherical caps $\pm C_\alpha$. Notice that $\theta^1(\beta_V, x) = \sqrt{3}$ for all $x \in \Gamma_a$. In other words, β_V does not have integer multiplicity even though $\theta^2(\mu_V, x) = 1$ for all $x \in S$. Notice also that V satisfies the hypothesis of Theorem 4.2 for all $c_0 \in \mathbb{R}$.

The set E in Hypothesis 4.5 corresponds to an enclosed volume in the following sense, where the algebraic volume does not depend on the point $x_0 \in \mathbb{R}^3$.

Proposition 4.8. *Suppose V, H, E, Θ satisfy Hypothesis 4.5 with $\text{spt } \mu_V$ compact. Then*

$$\mathcal{V}(V, x_0) = \int_E \Theta d\mathcal{L}^3 =: \mathcal{V}(V) \quad \text{for all } x_0 \in \mathbb{R}^3.$$

Proof. Since $\text{spt } \mu_V$ is compact, so is $\text{spt}(\mathcal{L}^3 \llcorner E)$ by Hypothesis 4.5. We may thus apply (4.19) with $X(x) = x - x_0$, suitably cutoff away from $\text{spt } \mu_V$ and $\text{spt}(\mathcal{L}^3 \llcorner E)$. \square

Under suitable assumptions, the concentrated volume can be computed by (4.19), too.

Lemma 4.9. *Suppose V, H, E, Θ satisfy Hypothesis 4.5. Let $x_0 \in \mathbb{R}^3$ and assume*

$$\lim_{\rho \rightarrow \infty} \frac{1}{\rho^2} \int_{E \cap B_\rho(x_0)} \Theta d\mathcal{L}^3 = 0. \quad (4.20)$$

Then we have

$$\mathcal{V}_c(V, x_0) = \int_E \frac{\Theta(x)}{|x - x_0|^2} d\mathcal{L}^3(x), \quad (4.21)$$

provided both sides exist.

Remark 4.10. (i) By Proposition 3.4, if $V \in \mathbb{V}_2^{\text{q}}(\mathbb{R}^3)$ and $H \in L^2_{\text{loc}}(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2 and if $\mathcal{V}(V, x_0)$ exists, then also $\mathcal{V}_c(V, x_0)$ exists.

(ii) Suppose $\int_E \Theta d\mathcal{L}^3 < \infty$. By Lemma 3.2 applied to the measure $\Theta \mathcal{L}^3 \llcorner E$, the right hand side of (4.21) exists if for some $m > 2$ we have

$$\limsup_{\sigma \rightarrow 0^+} \frac{1}{\sigma^m} \int_{E \cap B_\sigma(x_0)} \Theta d\mathcal{L}^3 < \infty. \quad (4.22)$$

As a consequence of the Lebesgue differentiation theorem, this is true for $m = 3$ and \mathcal{L}^3 -almost all $x_0 \in E$ (cf. [16, 2.9.8]). However, not all $\Theta \in L^1_{\text{loc}}(\mathcal{L}^3 \llcorner E; \mathbb{N})$ and $x_0 \in \mathbb{R}^3$ satisfy (4.22). This can be seen by taking $\Theta(x) := \lceil |x - x_0|^{-2} \rceil$. Nevertheless, (4.22) is clearly satisfied if $\Theta \in L^\infty(\mathcal{L}^3 \llcorner E; \mathbb{N})$.

Proof of Lemma 4.9. Since $\mathcal{V}_c(V, x_0)$ exists, we have

$$\int_{\mathbb{G}_2^+(\mathbb{R}^3)} \frac{|\langle x - x_0, \star\xi \rangle|}{|x - x_0|^2} dV(x, \xi) < \infty. \quad (4.23)$$

Now, let $0 < \sigma < \rho$ and let B_ρ, A_ρ be as in Lemma 4.1. Moreover, let φ be as in (4.2), $X(x) := x - x_0$, $\Phi(x) := \varphi(|x - x_0|)$ for $x \in \mathbb{R}^3$. For \mathcal{L}^3 -almost every $x \in \mathbb{R}^3$ we find

$$\operatorname{div}(\Phi X)(x) = \begin{cases} 3(\sigma^{-2} - \rho^{-2}) & \text{for } x \in B_\sigma \\ |X(x)|^{-2} - 3\rho^{-2} & \text{for } x \in B_\rho \setminus B_\sigma \\ 0 & \text{for } x \in \mathbb{R}^3 \setminus B_\rho. \end{cases}$$

Thus (4.19) implies

$$\begin{aligned} & -\frac{1}{\sigma^2} \int_{A_\sigma} \langle X(x), \star\xi \rangle dV(x, \xi) + \frac{1}{\rho^2} \int_{A_\rho} \langle X(x), \star\xi \rangle dV(x, \xi) - \int_{A_\rho \setminus A_\sigma} \frac{\langle X(x), \star\xi \rangle}{|X(x)|^2} dV(x, \xi) \\ & = \frac{3}{\sigma^2} \int_{E \cap B_\sigma} \Theta d\mathcal{L}^3 - \frac{3}{\rho^2} \int_{E \cap B_\rho} \Theta d\mathcal{L}^3 + \int_{E \cap B_\rho \setminus B_\sigma} \frac{\Theta(x)}{|X(x)|^2} d\mathcal{L}^3(x). \end{aligned} \quad (4.24)$$

We analyze each term in (4.24) separately. First, as $\sigma \rightarrow 0+$, the first term on the left vanishes, since (4.23) yields

$$\frac{1}{\sigma^2} \int_{A_\sigma} |\langle X(x), \star\xi \rangle| dV(x, \xi) \leq \int_{A_\sigma} \frac{|\langle X(x), \star\xi \rangle|}{|X(x)|^2} dV(x, \xi) \rightarrow 0.$$

Here we used that $\mu_V\{x_0\} = 0$ by Remark 3.3(i). The first term on the right hand side of (4.24) goes to zero as $\sigma \rightarrow 0$, since the right hand side of (4.21) exists. For the second term on the left, taking $0 < r < \rho$ and splitting the integral we obtain

$$\begin{aligned} & \frac{1}{\rho^2} \int_{A_\rho} |\langle X(x), \star\xi \rangle| dV(x, \xi) \\ & \leq \int_{\mathbb{G}_2^+(\mathbb{R}^3) \setminus A_r} \frac{|\langle X(x), \star\xi \rangle|}{|X(x)|^2} dV(x, \xi) + \frac{1}{\rho^2} \int_{A_r} |\langle X(x), \star\xi \rangle| dV(x, \xi), \end{aligned}$$

which goes to zero by (4.23), if we send first $\rho \rightarrow \infty$ and then $r \rightarrow \infty$. Taking $\rho \rightarrow \infty$ the second term on the right of (4.24) vanishes by (4.20). Thus, if we let first $\sigma \rightarrow 0$ and then $\rho \rightarrow \infty$ in (4.24) and use that both sides of (4.21) exist, the claim follows. \square

By the preceding discussion, the statement of Corollary 4.4 can be simplified if V is a varifold with enclosed volume. For simplicity, we only consider the case where $\operatorname{spt} \mu_V$ is compact.

Corollary 4.11. *Suppose V, H, E, Θ satisfy Hypothesis 4.5 with $\operatorname{spt} \mu_V$ compact. Then*

$$\theta^2(\mu_V, x_0) \leq \mathcal{H}_{c_0}(V) + \frac{c_0}{2\pi} \int_E \frac{\Theta(x)}{|x - x_0|^2} d\mathcal{L}^3(x)$$

for all $x_0 \in \mathbb{R}^3$, provided the second term on the right hand side exist.

Proof. By Hypothesis 4.5 we find that $\operatorname{spt}(\mathcal{L}^3 \llcorner E)$ is compact, so that using Remark 2.3 and Remark 4.10(i) we find that the assumptions of Corollary 4.4 and Lemma 4.9 are satisfied. The result then directly follows using (4.21). \square

5 The smooth setting

In this section, we will transfer the general varifold Li–Yau inequalities to the setting of smoothly immersed surfaces.

5.1 Proofs of the Li–Yau inequalities

Lemma 1.3 is an easy consequence of the varifold result.

Proof of Lemma 1.3. The claim follows directly from Corollary 4.4 if we consider the varifold associated to the immersion f , cf. Example 2.4. \square

We now show that any Alexandrov immersion induces a varifold with enclosed volume.

Lemma 5.1. *Let Σ be a closed surface and let $f: \Sigma \rightarrow \mathbb{R}^3$ be an Alexandrov immersion with $\Sigma = \partial M$, $f = F|_\Sigma$ and $F: M \rightarrow \mathbb{R}^3$ as in Definition 1.4. Let V be the oriented 2-varifold on \mathbb{R}^3 associated to (Σ, f) as in Example 2.4. Then, there holds*

$$-\int_{\mathbb{G}_2^+(\mathbb{R}^3)} \langle X(x), \star \xi \rangle dV(x, \xi) = \int_{F[M]} (\operatorname{div} X)(x) \mathcal{H}^0(F^{-1}\{x\}) d\mathcal{L}^3(x)$$

for all Lipschitz $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ with compact support. In particular, with $E := F[M]$, $\Theta := \mathcal{H}^0(F^{-1}\{\cdot\})$ we see that V, H, E, Θ satisfy Hypothesis 4.5.

Proof. By an approximation argument, it suffices to consider $X \in C_c^1(\mathbb{R}^3; \mathbb{R}^3)$. Denote with Ω the Riemannian measure on M induced by the pullback metric $g_F := F^*\langle \cdot, \cdot \rangle$, let μ be the induced measure on Σ , and let ν be the inner unit normal on Σ . Given any vector field $X \in C^1(\mathbb{R}^3; \mathbb{R}^3)$, we define the vector field X^* on M by $X^*(p) = (dF_p)^{-1}(X(F(p)))$. By (2.7) and since $n = dF(\nu)$, we compute

$$-\int_{\mathbb{G}_2^+(\mathbb{R}^3)} \langle X(x), \star \xi \rangle dV(x, \xi) = -\int_\Sigma \langle X \circ f, n \rangle d\mu = \int_{\partial M} g_F(X^*, -\nu) d\mu.$$

Since (M, g_F) is flat, we have $\operatorname{div}_{g_F} X^* = (\operatorname{div} X) \circ F$. Hence, by the divergence theorem for Riemannian manifolds (see [33, Theorem 5.11(2)]) and [36, Lemma 2.3],

$$\int_{\partial M} g_F(X^*, -\nu) d\mu = \int_M (\operatorname{div} X) \circ F d\Omega = \int_{F[M]} (\operatorname{div} X)(x) \mathcal{H}^0(F^{-1}\{x\}) d\mathcal{L}^3(x)$$

which implies the conclusion. \square

Equipped with this tool we can now prove Theorem 1.5.

Proof of Theorem 1.5. By Lemma 5.1, $V, H, E := F[M], \Theta := \mathcal{H}^0(F^{-1}\{\cdot\})$ satisfy Hypothesis 4.5. Since M is compact and F is a local diffeomorphism, there exists $k \in \mathbb{N}$ such that

$$\Theta(x) = \mathcal{H}^0(F^{-1}\{x\}) \leq k \quad \text{for all } x \in E = F[M],$$

and as a consequence of Lemma 4.9 and Remark 4.10(i) and (ii) we find

$$\mathcal{V}_c(f, x_0) = \int_{F[M]} \frac{\mathcal{H}^0(F^{-1}\{x\})}{|x - x_0|^2} d\mathcal{L}^3(x) \quad \text{for all } x_0 \in \mathbb{R}^3. \quad (5.1)$$

The statement then follows from Corollary 4.11. \square

Remark 5.2. The results of Lemma 1.3 and Theorem 1.5 are sharp in the sense that equality can be achieved asymptotically for every $c_0 \in \mathbb{R}$. Indeed, let $\mathbb{S}^2 \subset \mathbb{R}^3$ be the unit sphere, and let $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3, f(x) = rx$ denote the parametrization of the round sphere $\partial B_r(0) \subset \mathbb{R}^3$ with radius $r > 0$ and the orientation given by the inner unit normal. This is clearly an Alexandrov immersion (with $M = B_1(0), F(x) = rx$) and hence by (5.1), we have

$$\mathcal{V}_c(f, x_0) = \int_{B_r(0)} \frac{1}{|x - x_0|^2} d\mathcal{L}^3(x) = \begin{cases} 2\pi r & \text{if } x_0 \in \partial B_r(0) \\ 4\pi r & \text{if } x_0 = 0, \end{cases} \quad (5.2)$$

where the last equality follows from an explicit calculation in spherical coordinates. If now $x_0 \in \partial B_r(0)$, Inequality (1.8) reads

$$1 = \mathcal{H}^0(f^{-1}\{x_0\}) \leq \frac{1}{4\pi} \mathcal{H}_{c_0}(f) + \frac{c_0}{2\pi} \mathcal{V}_c(f, x_0) = \frac{1}{4} (c_0 r - 2)^2 + c_0 r \quad \text{for all } r > 0,$$

where the right hand side converges to 1 as $r \rightarrow 0+$. In the case $c_0 = 0$, equality is achieved by any round sphere.

5.2 A scale-invariant version

Clearly, for $x_0 = 0$ the left hand sides of the Li–Yau inequalities in Lemma 1.3 and Theorem 1.5 are invariant under rescalings of the immersion, whereas the right hand sides are not. We will now prove a scale-invariant version of the inequality, involving the L^2 -CMC-deficit of an immersion $f: \Sigma \rightarrow \mathbb{R}^3$ of an oriented surface Σ , given by

$$\bar{\mathcal{H}}(f) = \frac{1}{4} \int_{\Sigma} (H_{sc} - \bar{H}_{sc})^2 d\mu = \inf_{c_0 \in \mathbb{R}} \mathcal{H}_{c_0}(f),$$

cf. (1.3). Here $\bar{H}_{sc} := \mathcal{A}(f)^{-1} \int_{\Sigma} H_{sc} d\mu$ denotes the *average scalar mean curvature*, provided the latter integral exists. Note that $\bar{\mathcal{H}}(f) = 0$ if and only if f is an immersion with constant mean curvature, a *CMC-immersion*, justifying the terminology. We obtain the following Li–Yau inequality which is invariant under rescaling and also under reversing the orientation on Σ .

Corollary 5.3. *Let $f: \Sigma \rightarrow \mathbb{R}^3$ be an immersion of a closed oriented surface Σ . Then for all $x_0 \in \mathbb{R}^3$ we have*

$$\mathcal{H}^0(f^{-1}\{x_0\}) \leq \frac{1}{4\pi} \bar{\mathcal{H}}(f) + \frac{1}{2\pi} \bar{H}_{sc} \mathcal{V}_c(f, x_0) - \frac{1}{\pi \mathcal{A}(f)} (\mathcal{V}_c(f, x_0))^2. \quad (5.3)$$

Proof. By Proposition 3.4, we find that $\mathcal{V}_c(f, x_0)$ exists. We may thus use Lemma 1.3 for any $c_0 \in \mathbb{R}$. Expanding the right hand side of (1.8), we obtain a quadratic polynomial in c_0 . By a direct computation, this polynomial is minimal for $c_0 = \frac{\int_{\Sigma} H_{sc} d\mu - 4\mathcal{V}_c(f, x_0)}{\mathcal{A}(f)}$ and the minimal value is precisely the right hand side of (5.3). \square

6 Applications

In this section, we discuss several applications of the Li–Yau inequalities. We first provide a lower bound on the Helfrich energy resulting in nonexistence of minimizers for the penalized Canham–Helfrich model in Section 6.1. In Section 6.2 we prove some important geometric estimates involving the Helfrich energy. We then use these to prove Theorem 1.6. Lastly, we discuss a criterion for positive total mean curvature in Section 6.4.

6.1 Nonexistence of minimizers for the penalized Canham–Helfrich model

Lemma 6.1. *Suppose $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^2(\mu_V)$ satisfy Hypothesis 2.2, $\text{spt } \mu_V$ is compact, $c_0 < 0$, and $x_0 \in \mathbb{R}^3$ such that $\theta^{*2}(\mu_V, x_0) \geq 1$ and $\mathcal{V}_c(V, x_0) > 0$. Then there holds*

$$\mathcal{H}_{c_0}(V) > 4\pi.$$

Proof. This is a consequence of Corollary 4.4 in combination with [36, Theorem 3.6]. \square

Remark 6.2. The proof of the above inequality for the Willmore functional (i.e. $c_0 = 0$) [41, Theorem 7.2.2] also works for the Helfrich functional provided V is given by an Alexandrov immersion $f : \Sigma \rightarrow \mathbb{R}^3$ with inner unit normal field n . Indeed, denoting with K^+ the set of points in Σ where both principal curvatures are nonnegative, we find

$$\begin{aligned} \mathcal{H}_{c_0}(f) &\geq \frac{1}{4} \int_{K^+} |H_f - c_0 n|^2 \, d\mu \\ &\geq \frac{1}{4} \int_{K^+} |H_f|^2 \, d\mu + \frac{c_0^2}{4} \mathcal{A}(f) > \frac{1}{4} \int_{K^+} |H_f|^2 \, d\mu \geq \int_{K^+} K \, d\mu \end{aligned}$$

where K denotes the Gauss curvature. Similarly to [41, Lemma 7.2.1] we see that if f is an Alexandrov immersion, then

$$\int_{K^+} K \, d\mu \geq 4\pi.$$

For all real numbers c_0, λ, p we define the energy functional

$$\mathcal{H}_{c_0}^{\lambda,p}(f) := \mathcal{H}_{c_0}(f) + \lambda \mathcal{A}(f) + p \mathcal{V}(f)$$

for all smooth immersions $f : \Sigma \rightarrow \mathbb{R}^3$ of a closed oriented surface Σ . The constants λ and p are referred to as *tensile stress* and *osmotic pressure*. The energy was considered by Zhong-Can and Helfrich [42, Equation (1)] in the study of spherical vesicles. Each minimizer of the constrained minimization Problem 1.1 is a critical point of the functional $\mathcal{H}_{c_0}^{\lambda,p}$ for some λ, p by the method of Lagrange multipliers. This is one of the reasons why the energy $\mathcal{H}_{c_0}^{\lambda,p}$ is subject of numerous works in mathematical physics, biology and mathematics, see for instance [3] and the references therein.

Denote with \mathcal{S}^∞ the set of smoothly embedded spheres in \mathbb{R}^3 . In view of (1.2), we see that

$$\inf_{f \in \mathcal{S}^\infty} \mathcal{H}_{c_0}^{\lambda,p}(f) \leq 4\pi. \tag{6.1}$$

In [35, Theorem 3] (see also [29, Theorem 1.9]) the existence of spheres minimizing $\mathcal{H}_{c_0}^{\lambda,p}$ was shown, provided $\lambda, c_0 > 0$ and $p \geq 0$. However, in view of [11], $c_0 < 0$ is empirically more relevant in the study of red blood cells. Lemma 6.1 now reveals that the infimum in (6.1) is not attained whenever $c_0 < 0$ and $\lambda, p \geq 0$. This is actually in accordance with the results on the gradient flow in [25, 4]. Notice also the different behaviour of the constrained gradient flow [32]. Again exploiting the scaling properties of $\mathcal{H}_{c_0}^{\lambda,p}$, we see that the energy is unbounded from below if $p < 0$; in particular, the infimum in (6.1) is not attained. Similarly, if $\lambda < 0$ and $c_0^2 + \lambda < 0$, one can use surfaces of degenerating isoperimetric ratio found in [34, Theorem 1.5] to construct a sequence of embeddings f_k in \mathcal{S}^∞ such that $\mathcal{H}_{c_0}^{\lambda,p}(f_k) \rightarrow -\infty$ as $k \rightarrow \infty$.

Despite the nonexistence of minimizers explained above, the energy functional $\mathcal{H}_{c_0}^{\lambda,p}$ remains an important subject of study, since it is the critical points of $\mathcal{H}_{c_0}^{\lambda,p}$ that are of interest.

6.2 Diameter estimates

In this section, we will show that the Helfrich energy can be used to obtain bounds on the diameter.

Lemma 6.3. *Suppose $V \in \mathbb{V}_2^0(\mathbb{R}^3)$ and $H \in L_{\text{loc}}^2(\mu_V; \mathbb{R}^3)$ satisfy Hypothesis 2.2, $\text{spt } \mu_V$ is compact, and $\mathcal{H}_{c_0}(V) > 0$. Then for all $x_0 \in \text{spt } \mu_V$ we have*

$$\frac{|2\mu_V(\mathbb{R}^3) - 3c_0 \mathcal{V}(V, x_0)|}{2\sqrt{\mu_V(\mathbb{R}^3)\mathcal{H}_{c_0}(V)}} \leq \text{diam spt } \mu_V.$$

Proof. Using Hypothesis 2.2 for the vector field $X(x) = x - x_0$ (multiplied with a suitable cut-off function away from $\text{spt } \mu_V$), we have

$$\begin{aligned} \int_{\mathbb{R}^3} 2d\mu_V(x) - 3c_0 \mathcal{V}(V, x_0) &= - \int_{\mathbb{G}_2^0(\mathbb{R}^3)} \langle H(x), x - x_0 \rangle dV(x, \xi) \\ &\quad + c_0 \int_{\mathbb{G}_2^0(\mathbb{R}^3)} \langle \star\xi, x - x_0 \rangle dV(x, \xi). \end{aligned}$$

Thus, by the Cauchy–Schwarz inequality

$$\begin{aligned} |2\mu_V(\mathbb{R}^3) - 3c_0 \mathcal{V}(V, x_0)| &\leq \int_{\mathbb{G}_2^0(\mathbb{R}^3)} |H(x) - c_0(\star\xi)| |x - x_0| dV(x, \xi) \\ &\leq \sqrt{4\mathcal{H}_{c_0}(V)\mu_V(\mathbb{R}^3)} \text{diam spt } \mu_V. \quad \square \end{aligned}$$

In the case $c_0 = 0$, this is just Simon’s lower diameter estimate, cf. [40, Lemma 1.1]. Note that here we did not use the Li–Yau inequality but merely the first variation formula, see (2.1) and (2.4).

Lemma 6.4. *Suppose V, H, E, Θ satisfy Hypothesis 4.5, $\theta^2(\mu_V, x) \geq 1$ for μ_V -almost all x , $\text{spt } \mu_V$ is connected, and $c_0 \leq 0$. If $\mathcal{H}_{c_0}(V) < \infty$, $\mu_V(\mathbb{R}^3) < \infty$, $\Theta \in L^1(\mathcal{L}^3 \llcorner E)$, and*

$$\int_E \frac{\Theta(x)}{|x - x_0|^2} d\mathcal{L}^3(x) < \infty \quad (6.2)$$

for μ_V -almost all x_0 , then $\text{spt } \mu_V$ is compact and

$$\text{diam spt } \mu_V \leq C \sqrt{\mathcal{H}_{c_0}(V) \left(\mu_V(\mathbb{R}^3) + \frac{2}{3} |c_0| \mathcal{V}(V) \right)} \quad (6.3)$$

where $C = \frac{9}{2\pi}$ and $\mathcal{V}(V) = \int_E \Theta \, d\mathcal{L}^3$ is the algebraic volume (see Proposition 4.8).

Remark 6.5. For $c_0 = 0$, we recover the diameter bound in terms of area and Willmore energy by Simon [40, Lemma 1.1]:

$$\text{diam spt } \mu_V \leq C \sqrt{\mathcal{W}(V) \mu_V(\mathbb{R}^3)}. \quad (6.4)$$

This inequality holds true for all 2-varifolds in \mathbb{R}^3 with generalized perpendicular mean curvature, finite Willmore energy, and whose weight measure is finite and has connected support (see [36, Theorem 1.5]). Hence, by (2.5) we obtain

$$\text{diam spt } \mu_V \leq C \sqrt{\mu_V(\mathbb{R}^3) \left(2\mathcal{H}_{c_0}(V) + \frac{1}{2} c_0^2 \mu_V(\mathbb{R}^3) \right)}$$

for all V satisfying Hypothesis 2.2 with $\mathcal{W}(V) < \infty$, $\mu_V(\mathbb{R}^3) < \infty$, and such that $\text{spt } \mu_V$ is connected. Recall that by (1.6) there are smooth embeddings $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$ such that $\mathcal{H}_{c_0}(f) < \mathcal{W}(f)$. Hence, in general, we do not expect that (6.4) holds true for \mathcal{W} replaced by \mathcal{H}_{c_0} . However, for small algebraic volume, the right hand side in (6.3) is close to the right hand side in (6.4) with \mathcal{W} replaced by \mathcal{H}_{c_0} .

Proof of Lemma 6.4. We will follow the proof of [40, Lemma 1.1]. Suppose $\text{spt } \mu_V \neq \emptyset$ (otherwise the statement is trivial), let $x_0 \in \text{spt } \mu_V$ and define the Radon measure

$$\mathcal{H}_{c_0}(V, B) := \frac{1}{4} \int_{B \times \mathbb{G}^\circ(3,2)} |H(x) - c_0(\star\xi)|^2 \, dV(x, \xi) \quad \text{for all Borel sets } B \text{ in } \mathbb{R}^3.$$

Using the Cauchy–Schwarz inequality as in (4.13) and Young’s inequality, we estimate

$$\frac{1}{2\rho^2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} |\langle x - x_0, H(x) - c_0(\star\xi) \rangle| \, dV(x, \xi) \leq \frac{\mu_V(B_\rho(x_0))}{2\rho^2} + \frac{1}{2} \mathcal{H}_{c_0}(V, B_\rho(x_0))$$

for all $\rho > 0$. Hence, since $\theta^2(\mu_V, x_0) \geq 1$ by [36, Theorem 3.6] in combination with Remark 2.3, we can let σ go to zero in Lemma 4.1 and use (4.13), (4.14) to infer

$$\begin{aligned} \pi &\leq \frac{3}{4} \mathcal{H}_{c_0}(V, B_\rho(x_0)) + \frac{3\mu_V(B_\rho(x_0))}{2\rho^2} - \frac{c_0}{2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} \, dV(x, \xi) \\ &\quad + \frac{c_0}{2\rho^2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} \langle x - x_0, \star\xi \rangle \, dV(x, \xi). \end{aligned} \quad (6.5)$$

Exactly as in (4.24), for $0 < \sigma < \rho$ we may use (4.19) to obtain

$$- \frac{c_0}{2} \int_{(B_\rho \setminus B_\sigma)(x_0) \times \mathbb{G}^\circ(3,2)} \frac{\langle x - x_0, \star\xi \rangle}{|x - x_0|^2} \, dV(x, \xi) + \frac{c_0}{2\rho^2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} \langle x - x_0, \star\xi \rangle \, dV(x, \xi)$$

$$\begin{aligned}
&= \frac{3|c_0|}{2\rho^2} \int_{E \cap B_\rho(x_0)} \Theta \, d\mathcal{L}^3 - \frac{|c_0|}{2} \int_{E \cap B_\rho(x_0) \setminus B_\sigma(x_0)} \frac{\Theta(x)}{|x-x_0|^2} \, d\mathcal{L}^3(x) \\
&\quad - \frac{3|c_0|}{2\sigma^2} \int_{E \cap B_\sigma(x_0)} \Theta \, d\mathcal{L}^3 - \frac{|c_0|}{2\sigma^2} \int_{B_\sigma(x_0) \times \mathbb{G}^\circ(3,2)} \langle x-x_0, \star\xi \rangle \, dV(x, \xi).
\end{aligned}$$

Sending $\sigma \rightarrow 0+$ and using (6.2), Lemma 3.2 and Remark 3.3(i), we find that

$$\begin{aligned}
&- \frac{c_0}{2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} \frac{\langle x-x_0, \star\xi \rangle}{|x-x_0|^2} \, dV(x, \xi) + \frac{c_0}{2\rho^2} \int_{B_\rho(x_0) \times \mathbb{G}^\circ(3,2)} \langle x-x_0, \star\xi \rangle \, dV(x, \xi) \\
&= \frac{|c_0|}{2} \int_{E \cap B_\rho(x_0)} \left(\frac{3}{\rho^2} - \frac{1}{|x-x_0|^2} \right) \Theta(x) \, d\mathcal{L}^3(x) \leq \frac{|c_0|}{\rho^2} \int_{E \cap B_\rho(x_0)} \Theta(x) \mathcal{L}^3(x). \quad (6.6)
\end{aligned}$$

Combining (6.5) and (6.6), we thus obtain

$$\pi \leq \frac{3}{4} \mathcal{H}_{c_0}(V, B_\rho(x_0)) + \frac{3}{2\rho^2} \mu_V(B_\rho(x_0)) + \frac{|c_0|}{\rho^2} (\Theta \mathcal{L}^3 \llcorner E)(B_\rho(x_0)). \quad (6.7)$$

The right hand side of this inequality corresponds to the Radon measure

$$\mu_{c_0, V, E} := \frac{3}{4} \mathcal{H}_{c_0}(V, \cdot) + \frac{3}{2\rho^2} \mu_V + \frac{|c_0|}{\rho^2} (\Theta \mathcal{L}^3 \llcorner E).$$

The set of $x_0 \in \mathbb{R}^3$ that satisfy (6.2) is dense in $\text{spt } \mu_V$. Hence, given any $x_0 \in \text{spt } \mu_V$ and any $\varepsilon > 0$ we can always find $x_1 \in \text{spt } \mu_V$ which satisfies (6.7) such that $B_\rho(x_1) \subset B_{\rho+\varepsilon}(x_0)$. Thus, letting $\varepsilon \rightarrow 0+$, we see that (6.7) remains valid for all $x_0 \in \text{spt } \mu_V$. By Remark 3.3(i), we see $\mu_V(N) = 0$ whenever N is finite and consequently $\mu_{c_0, V, E}(N) = 0$ whenever N is finite. Let $d := \text{diam spt } \mu_V$ (possibly $d = \infty$), $0 < \rho < d$, and N be a positive integer such that $2(N-1)\rho < d$. By the connectedness of $\text{spt } \mu_V$, we can choose points $x_0, \dots, x_{N-1} \in \text{spt } \mu_V$ such that $x_i \in \partial B_{2i\rho}(x_0)$ for $i = 1, \dots, N-1$. The balls $B_\rho(x_0), \dots, B_\rho(x_{N-1})$ intersect in at most $N-1$ points. Applying the inequality (6.7) for each x_i and summing over i yields

$$N\pi \leq \mu_{c_0, V, E}(\mathbb{R}^3). \quad (6.8)$$

Since the right hand side is finite, it follows that the diameter d is finite. Hence, we can choose N such that $2(N-1)\rho < d \leq 2N\rho$. Then (6.8) and Proposition 4.8 imply

$$d \leq \frac{3}{2\pi} \left(\rho \mathcal{H}_{c_0}(V) + \frac{2}{\rho} \mu_V(\mathbb{R}^3) + \frac{4|c_0|}{3\rho} \mathcal{V}(V) \right). \quad (6.9)$$

Now, in view of Lemma 6.1, we may take

$$\rho = \sqrt{\frac{2\mu_V(\mathbb{R}^3) + \frac{4}{3}|c_0| \mathcal{V}(V)}{2\mathcal{H}_{c_0}(V)}} = \sqrt{\frac{\mu_V(\mathbb{R}^3) + \frac{2}{3}|c_0| \mathcal{V}(V)}{\mathcal{H}_{c_0}(V)}}.$$

Then, by Lemma 6.3, $\rho < d$ and thus, (6.9) becomes

$$d \leq \frac{9}{2\pi} \sqrt{\mathcal{H}_{c_0}(V) \left(\mu_V(\mathbb{R}^3) + \frac{2}{3}|c_0| \mathcal{V}(V) \right)}$$

which concludes the proof. \square

6.3 Regularity and embeddedness of Canham–Helfrich minimizers

We start with a survey of the variational setting in [30] (see also [21, 20, 29]). This includes the definition of *Lipschitz immersions*. Then we introduce the space \mathcal{S}_Σ of *Lipschitz cells* which consists of those Lipschitz immersions whose associated varifolds are varifolds with enclosed volume, cf. Hypothesis 4.5. We show that each injective Lipschitz immersion (in particular each smooth embedding) is a Lipschitz cell (see Lemma 6.6). Moreover, we prove a weak closure Lemma 6.8 which leads to our main regularity Theorem 6.10.

Let Σ be a closed oriented surface and let g_0 be a reference Riemannian metric on Σ . A map $f: \Sigma \rightarrow \mathbb{R}^3$ is called *weak branched immersion* if and only if

$$f \in W^{1,\infty}(\Sigma; \mathbb{R}^3) \cap W^{2,2}(\Sigma; \mathbb{R}^3), \quad (6.10)$$

there exists a constant $1 < C < \infty$ such that

$$C^{-1}|df|_{g_0} \leq |df \wedge df|_{g_0} \leq C|df|_{g_0} \quad (6.11)$$

where in local coordinates

$$df \wedge df := (dx^1 \wedge dx^2) \partial_{x^1} f \wedge \partial_{x^2} f,$$

there exist finitely many so called *branch points* $b_1, \dots, b_N \in \Sigma$ such that the *conformal factor* satisfies

$$\log |df|_{g_0} \in L_{\text{loc}}^\infty(\Sigma \setminus \{b_1, \dots, b_N\}),$$

and the Gauss map n defined as in (2.6) satisfies

$$n \in W^{1,2}(\Sigma; \mathbb{R}^3). \quad (6.12)$$

If in addition

$$|\partial_{x^1} f| = |\partial_{x^2} f| \quad \text{and} \quad \langle \partial_{x^1} f, \partial_{x^2} f \rangle = 0 \quad (6.13)$$

for all conformal charts x of (Σ, g_0) , then f is called *conformal*. A chart $x = (x^1, x^2)$ that satisfies (6.13) is referred to as *isothermal coordinates*. Notice that (6.13) implies (6.11) and, since Σ is closed, the conditions (6.10)–(6.12) do not depend on the choice of the Riemannian metric g_0 . The space of weak branched immersions is denoted by \mathcal{F}_Σ . The subspace \mathcal{E}_Σ of *Lipschitz immersions* is defined to consist of all $f \in \mathcal{F}_\Sigma$ such that there exists a constant $0 < C < \infty$ with

$$|df \wedge df|_{g_0} \geq C. \quad (6.14)$$

Notice that (6.10) and (6.14) imply $\log |df|_{g_0} \in L^\infty(\Sigma)$.

Let $f \in \mathcal{F}_\Sigma$. Analogously to Example 2.4, we infer a (possibly degenerated) L^∞ -metric $g := f^*\langle \cdot, \cdot \rangle$, the induced Radon measure μ over Σ , the oriented varifold $V := (f, \star n)_\# \mu$, the classical mean curvature H_f of f (in the Sobolev sense), and the induced generalized mean curvature H . If f is conformal, we have by [21, Theorem 3.1] that

$$\mathcal{H}^0(f^{-1}\{x\}) = \theta^2(\mu_V, x) \quad \text{for all } x \in \mathbb{R}^3.$$

In view of [29, Equation (2.11)] there holds

$$\delta V(X) = - \int_{\mathbb{R}^3} \langle X, H \rangle d\mu_V$$

for all $X \in C_c^1(\mathbb{R}^3; \mathbb{R}^3)$. Moreover, by the definition of H and (6.12) we have that

$$\int_{\mathbb{R}^3} |H|^2 d\mu_V \leq \int_{\Sigma} |H_f|^2 d\mu < \infty.$$

In particular, $H \in L^2(\mu_V; \mathbb{R}^3)$ and V, H satisfy Hypothesis 2.2. Now, we can combine [15, Section 6.1, Theorem 4] and [37, Theorem 4.1] to infer

$$H(f(p)) = H_f(p) \quad \text{for } \mu\text{-almost all } p \in \Sigma.$$

As in (2.9), it follows that for all $c_0 \in \mathbb{R}$ we have

$$\mathcal{H}_{c_0}(V) = \frac{1}{4} \int_{\mathbb{R}^3} |H(x) - c_0(\star\xi)|^2 dV(x, \xi) = \int_{\Sigma} |H_f - c_0 n|^2 d\mu = \mathcal{H}_{c_0}(f). \quad (6.15)$$

The space \mathcal{S}_{Σ} of *Lipschitz cells* is defined to consist of all $f \in \mathcal{E}_{\Sigma}$ such that there exists an \mathcal{L}^3 -measurable set E with

$$\text{diam spt}(\mathcal{L}^3 \llcorner E) \leq \text{diam } f[\Sigma] \quad (6.16)$$

and

$$\int_E \text{div } X d\mathcal{L}^3 = - \int_{\Sigma} \langle X \circ f, n \rangle d\mu \quad (6.17)$$

for any Lipschitz map $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$ with compact support, i.e. the triple E, V, H satisfies Hypothesis 4.5 for $\Theta \equiv 1$. The divergence theorem for sets of finite perimeter (4.16), Equation (6.17), and the area formula (see [16, 3.2.22(3)]) imply

$$n_E(x) = \sum_{p \in f^{-1}\{x\}} n(p) \quad \text{for } \mathcal{H}^2\text{-almost all } x \in \mathbb{R}^3. \quad (6.18)$$

Notice that $x \notin \text{spt}(\mathcal{H}^2 \llcorner \partial_* E)$ does not imply $f^{-1}\{x\} = \emptyset$. In particular, in view of Figures (5) and (6), the two oriented varifolds associated with $\partial_* E$ and f do not necessarily coincide. Hence, by Proposition 3.4, Proposition 4.8, Lemma 4.9 and Remark 4.10 there holds

$$\mathcal{V}(V, x_0) = \mathcal{V}(f) = \mathcal{L}^3(E), \quad \mathcal{V}_c(V, x_0) = \mathcal{V}_c(f, x_0) = \int_E \frac{1}{|x - x_0|^2} d\mathcal{L}^3(x) \quad (6.19)$$

for all $x_0 \in \mathbb{R}^3$. If $\mathcal{H}^0(f^{-1}\{x\}) \leq 1$ for \mathcal{H}^2 -almost all $x \in \mathbb{R}^3$ then (6.18) implies $n_E \circ f = n$, $\partial_* E = f[\Sigma]$ up to a set of \mathcal{H}^2 -measure zero, the two oriented varifolds associated with $\partial_* E$ and f coincide, and since $\text{spt}(\mathcal{H}^2 \llcorner \partial_* E) \subset \text{spt}(\mathcal{L}^3 \llcorner E)$, equality holds in (6.16).

Lemma 6.6. *Let Σ be a closed oriented surface and $f \in \mathcal{E}_\Sigma$ be injective. Then, possibly after changing the orientation of Σ , there exists a connected open bounded set $U \subset \mathbb{R}^3$ of finite perimeter such that $\partial_* U = f[\Sigma]$ up to a set of \mathcal{H}^2 -measure zero, and*

$$\int_U \operatorname{div} X \, d\mathcal{L}^3 = - \int_\Sigma \langle X \circ f, n \rangle \, d\mu \quad (6.20)$$

for any Lipschitz map $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$. In particular, $f \in \mathcal{S}_\Sigma$ and \mathcal{S}_Σ contains all smooth embeddings $f: \Sigma \rightarrow \mathbb{R}^3$ (up to orientation). However, not all $f \in \mathcal{S}_\Sigma$ are injective.

Remark 6.7. Notice that changing the orientation on Σ is equivalent to changing the sign of the (nonzero) algebraic volume. Hence, if additionally $\mathcal{V}(f) > 0$, no change of orientation is necessary in Lemma 6.6.

Proof of Lemma 6.6. We may assume that $(\Sigma, g_0) \subset \mathbb{R}^3$ is embedded and g_0 is the metric induced by the inclusion map. Since f is injective, we can apply the Jordan–Brouwer separation theorem [7] to obtain a connected open bounded set $U \subset \mathbb{R}^3$ such that $\partial U = f[\Sigma]$ and $\mathbb{R}^3 \setminus \bar{U}$ is connected. Since $\partial_* U \subset \partial U = f[\Sigma]$ and $\mathcal{H}^2(f[\Sigma]) < \infty$, Federer’s criterion implies that U is a set of finite perimeter. Moreover, for $p \in \Sigma$, one can show that if f is differentiable at p , then $f(p) \in \partial_* U$. Hence, by Rademacher’s theorem, the sets $\partial_* U$ and $f[\Sigma]$ are \mathcal{H}^2 -almost equal. We still need to show that

$$\int_{\partial_* U} \langle X, n_U \rangle \, d\mathcal{H}^2 = \int_\Sigma \langle X \circ f, n \rangle \, d\mu$$

for all Lipschitz maps $X: \mathbb{R}^3 \rightarrow \mathbb{R}^3$, where n_U is the measure theoretic inner unit normal of U (see Section 4.3), and n is the Gauss map of f , cf. (2.6). Let ν be the unit normal induced by the orientation of $\Sigma \subset \mathbb{R}^3$. We define the 2-current T on \mathbb{R}^3 by

$$T(\omega) := - \int_\Sigma \omega_p(\star \nu(p)) \, d\mathcal{H}^2(p)$$

for all differential forms ω of degree 2 on \mathbb{R}^3 . Since Σ is closed, we have

$$\partial T = 0 \quad (6.21)$$

(see for instance [16, 4.1.31(1)]). Given any positive chart x of Σ , there holds

$$\nu = \frac{\partial_{x^1} \times \partial_{x^2}}{|\partial_{x^1} \times \partial_{x^2}|}, \quad (\wedge_2 df)(\star \nu) = \frac{|\partial_{x^1} f \wedge \partial_{x^2} f|}{|\partial_{x^1} \times \partial_{x^2}|}(\star n)$$

where for \mathcal{H}^2 -almost all $p \in \Sigma$, the linear map $\wedge_2 df_p: \wedge_2 T_p \Sigma \rightarrow \wedge_2 df_p[T_p \Sigma]$ is defined as in [16, 1.3.1]. Recalling that in any local chart x , the area elements of the immersion f and the inclusion $\Sigma \subset \mathbb{R}^3$ are given by $|\partial_{x^1} f \wedge \partial_{x^2} f|$ and $|\partial_{x^1} \times \partial_{x^2}|$, respectively, we have by [16, 4.1.30] that

$$R(\omega) := (f_\# T)(\omega) = - \int_\Sigma \omega_{f(p)}(\star n(p)) \, d\mu(p)$$

for all differential forms ω of degree 2 on \mathbb{R}^3 . Thus, by [16, 4.1.14] and (6.21)

$$\partial R = \partial(f\#T) = f\#(\partial T) = 0.$$

Therefore, we can combine [16, 4.5.17] and [16, 4.5.6] to deduce the existence of sets of finite perimeter $E_j \subset E_{j-1} \subset \mathbb{R}^3$, $j \in \mathbb{Z}$ such that

$$R = \sum_{j \in \mathbb{Z}} R_j, \quad \mu_V = \sum_{j \in \mathbb{Z}} (\mathcal{H}^2 \llcorner \partial_* E_j) \quad (6.22)$$

where

$$R_j(\omega) := - \int_{\partial_* E_j} \omega_x(\star n_{E_j}(x)) \, d\mathcal{H}^2(x)$$

are the currents induced by $\partial_* E_j$. Since U is open and connected, we see from [2, Proposition 2] that U is indecomposable. Given any set of finite perimeter $E \subset \mathbb{R}^3$ with $\partial_* E \subset \partial_* U$ up to a set of \mathcal{H}^2 -measure zero, we see that $\partial_*(U \cap E) \subset \partial_* U$ up to a set of \mathcal{H}^2 -measure zero and thus, by [2, Proposition 4], either $\mathcal{L}^3(U \cap E) = 0$ or $\mathcal{L}^3(U \setminus E) = 0$. The same holds true for U replaced by $\mathbb{R}^3 \setminus \bar{U}$. By (6.22) we have for all $j \in \mathbb{Z}$ that $\partial_* E_j \subset \text{spt } \mu_V = \partial_* U$ up to a set of \mathcal{H}^2 -measure zero and therefore either $E_j = \mathbb{R}^3$ or $E_j = U$ or $E_j = \mathbb{R}^3 \setminus \bar{U}$ or $E_j = \emptyset$ up to a set of \mathcal{L}^3 -measure zero. Since f is injective, we have that $\theta^2(\mu_V, \cdot) \leq 1$. We thus deduce the existence of $j_0 \in \mathbb{Z}$ such that (up to a set of \mathcal{L}^3 -measure zero and possibly after changing the orientation on Σ)

$$E_j = \begin{cases} \mathbb{R}^3 & \text{for } j < j_0, \\ U & \text{for } j = j_0, \\ \emptyset & \text{for } j > j_0. \end{cases}$$

In particular, $R = R_{j_0}$ and (6.20) follows. To see that not all $f \in \mathcal{S}_\Sigma$ are injective, one may consider surfaces like in Figures (2), (4), and (5). \square

In the following, we abbreviate $\mathcal{F} := \mathcal{F}_{\mathbb{S}^2}$, $\mathcal{E} := \mathcal{E}_{\mathbb{S}^2}$, and $\mathcal{S} := \mathcal{S}_{\mathbb{S}^2}$.

Lemma 6.8. *Suppose f_k is a sequence in \mathcal{S} , $0 \in f_k[\mathbb{S}^2]$ for all $k \in \mathbb{N}$, $c_0 \in \mathbb{R}$,*

$$A_0 := \sup_{k \in \mathbb{N}} \mathcal{A}(f_k) < \infty, \quad \inf_{k \in \mathbb{N}} \text{diam } f_k[\mathbb{S}^2] > 0, \quad (6.23)$$

and

$$\begin{cases} \liminf_{k \rightarrow \infty} (\mathcal{H}_{c_0}(f_k) + 2c_0 \inf_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x)) < 8\pi & \text{if } c_0 < 0, \\ \liminf_{k \rightarrow \infty} (\mathcal{H}_{c_0}(f_k) + 2c_0 \sup_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x)) < 8\pi & \text{if } c_0 \geq 0. \end{cases} \quad (6.24)$$

Then, after passing to a subsequence, there exists $f \in \mathcal{S}$ injective such that

$$V_k \rightarrow V \quad \text{in } \mathbb{V}_2^0(\mathbb{R}^3) \text{ as } k \rightarrow \infty, \quad (6.25)$$

where V_k, V are the oriented 2-varifolds in \mathbb{R}^3 associated with f_k, f (cf. Example 2.4) and

$$\mathcal{H}_{c_0}(f) \leq \liminf_{k \rightarrow \infty} \mathcal{H}_{c_0}(f_k). \quad (6.26)$$

Proof. Let g_0 be the standard metric on \mathbb{S}^2 . By [30, Theorem 1.4], after reparametrization, we may assume that all f_k are conformal. After passing to a subsequence, we may further assume that for all $k \in \mathbb{N}$,

$$\begin{cases} \mathcal{H}_{c_0}(f_k) + 2c_0 \inf_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x) < 8\pi & \text{if } c_0 < 0, \\ \mathcal{H}_{c_0}(f_k) + 2c_0 \sup_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x) < 8\pi & \text{if } c_0 \geq 0. \end{cases}$$

Let E_k be the sequence of sets of finite perimeter corresponding to f_k according to (6.17). Using (6.19), for all $x_0 \in \mathbb{R}^3$ and $k \in \mathbb{N}$ there holds

$$\mathcal{V}_c(f_k, x_0) \leq \int_{B_1(x_0)} \frac{1}{|x - x_0|^2} d\mathcal{L}^3(x) + \mathcal{L}^3(E_k) = 4\pi + \mathcal{V}(f_k)$$

and thus, by (6.16) we can apply the isoperimetric inequality for sets of finite perimeter (see [16, Corollary 4.5.3]) to deduce from the uniform area bound (6.23) that

$$V_0 := \sup_{k \in \mathbb{N}} \mathcal{V}(f_k) < \infty, \quad \sup_{k \in \mathbb{N}} \sup_{x \in \mathbb{R}^3} \mathcal{V}_c(f_k, x) < C(V_0) < \infty. \quad (6.27)$$

Hence, by [29, Equation (2.8)] and (2.5), there holds

$$\begin{aligned} \int_{\mathbb{S}^2} 1 + |dn_{f_k}|_{g_0}^2 d\mu_{f_k} &\leq A_0 + 4\mathcal{W}(f_k) \leq A_0 + 8\mathcal{H}_{c_0}(f_k) + 2c_0^2 A_0 \\ &\leq A_0 + 8(8\pi + 2|c_0|C(V_0) + c_0^2 A_0) \end{aligned} \quad (6.28)$$

for all $k \in \mathbb{N}$. Therefore, we can apply [29, Theorem 3.3] (see also Theorem 1.5 and Lemma 4.1 in [28]) to infer that after passing to a subsequence, there exist a positive integer N and sequences $\phi_k^1, \dots, \phi_k^N$ of positive conformal C^∞ -diffeomorphisms of \mathbb{S}^2 such that for each $i \in \{1, \dots, N\}$, there exist $f^i \in \mathcal{F}_{\mathbb{S}^2}$ conformal, $N^i \in \mathbb{N}$, and finitely many points $b^{i,1}, \dots, b^{i,N^i} \in \mathbb{S}^2$ with

$$f_k^i := f_k \circ \phi_k^i \rightharpoonup f^i \quad \text{weakly in } W_{\text{loc}}^{2,2}(\mathbb{S}^2 \setminus \{b^{i,1}, \dots, b^{i,N^i}\}; \mathbb{R}^3) \text{ as } k \rightarrow \infty, \quad (6.29)$$

$$\sup_{k \in \mathbb{N}} \|\log |df_k^i|_{g_0}\|_{L_{\text{loc}}^\infty(\mathbb{S}^2 \setminus \{b^{i,1}, \dots, b^{i,N^i}\})} < \infty. \quad (6.30)$$

Moreover, there exist a sequence ψ_k of C^∞ -diffeomorphisms of \mathbb{S}^2 and $f \in W^{1,\infty}(\mathbb{S}^2; \mathbb{R}^3)$ such that

$$f_k \circ \psi_k \rightarrow f \quad \text{in } C^0(\mathbb{S}^2; \mathbb{R}^3) \text{ as } k \rightarrow \infty, \quad f[\mathbb{S}^2] = \bigcup_{i=1}^N f^i[\mathbb{S}^2]. \quad (6.31)$$

Furthermore, there holds

$$\sum_{i=1}^N \mathcal{W}(f^i) \leq \liminf_{k \rightarrow \infty} \mathcal{W}(f_k), \quad \sum_{i=1}^N \mathcal{H}_{c_0}(f^i) \leq \liminf_{k \rightarrow \infty} \mathcal{H}_{c_0}(f_k). \quad (6.32)$$

Denote with V^i the varifolds associated to f^i and set $V := \sum_{i=1}^N V^i$. In order to show (6.25), let $\varphi: \mathbb{R}^3 \times \mathbb{G}^o(3, 2) \rightarrow \mathbb{R}$ be any continuous function with compact support.

Fix an integer $i \in \{1, \dots, N\}$, choose a conformal chart $x: \mathbb{S}^2 \setminus \{b^{i,1}, \dots, b^{i,N^i}\} \rightarrow \mathbb{R}^2$, and let $K \subset \mathbb{S}^2 \setminus \{b^{i,1}, \dots, b^{i,N^i}\}$ be a compact set. Denote by

$$\lambda_k^i := \log |\partial_{x^1} f_k^i|, \quad \lambda^i := \log |\partial_{x^1} f^i|$$

the conformal factors and recall that the area elements of f_k^i and f^i are given by $e^{2\lambda_k^i}$ and $e^{2\lambda^i}$. Let n_k^i , n_k , and n^i be the Gauss maps of f_k^i , f_k , and f^i . Following the proof of [29, Lemma 3.1], we infer that by the weak convergence (6.29), the Rellich–Kondrachov compactness theorem, and the uniform bounds on the conformal factors (6.30), after passing to a subsequence,

$$\begin{aligned} e^{2\lambda_k^i} \circ x^{-1} &\rightarrow e^{2\lambda^i} \circ x^{-1} && \text{in } L^p(x[K]) \text{ as } k \rightarrow \infty \text{ for all } 1 \leq p < \infty, \\ f_k^i \circ x^{-1} &\rightarrow f^i \circ x^{-1} && \text{pointwise almost everywhere on } x[K] \text{ as } k \rightarrow \infty, \\ n_k^i \circ x^{-1} &\rightarrow n^i \circ x^{-1} && \text{pointwise almost everywhere on } x[K] \text{ as } k \rightarrow \infty. \end{aligned} \quad (6.33)$$

Hence, since φ is continuous, and also the Hodge star operator \star is continuous,

$$\varphi(f_k^i, \star n_k^i) \circ x^{-1} \rightarrow \varphi(f^i, \star n^i) \circ x^{-1} \quad \text{pointwise almost everywhere on } x[K]$$

as $k \rightarrow \infty$. Thus, since φ is bounded, the dominated convergence theorem and (6.33) imply

$$\left(\varphi(f_k^i, \star n_k^i) e^{2\lambda_k^i} \right) \circ x^{-1} \rightarrow \left(\varphi(f^i, \star n^i) e^{2\lambda^i} \right) \circ x^{-1} \quad \text{in } L^p(x[K]) \text{ as } k \rightarrow \infty$$

for any $1 \leq p < \infty$. Therefore, inductively passing to a subsequence, we can achieve that for all $k_0 \in \mathbb{N}$ and all $k_0 \leq k \in \mathbb{N}$, there holds

$$\int_x \left[\mathbb{S}^2 \setminus \bigcup_{j=1}^{N^i} B_{\frac{1}{k_0}}(b^{i,j}) \right] \left| \varphi(f_k^i, \star n_k^i) e^{2\lambda_k^i} - \varphi(f^i, \star n^i) e^{2\lambda^i} \right| \circ x^{-1} d\mathcal{L}^2 \leq \frac{1}{k_0}. \quad (6.34)$$

Successively passing to a subsequence, we infer that (6.34) holds true simultaneously for all $i \in \{1, \dots, N\}$. (Notice however that the chart x actually depends on i .) Moreover, since φ is bounded and by the fact that finite sets have μ_{f_i} -measure zero by Remark 3.3(i), there holds

$$\lim_{k_0 \rightarrow \infty} \int_{\bigcup_{j=1}^{N^i} B_{\frac{1}{k_0}}(b^{i,j})} \varphi(f^i, \star n^i) d\mu_{f^i} = 0,$$

for all $i \in \{1, \dots, N\}$. Writing $s_k := 1/k$, by (2.7) it follows

$$\int_{\mathbb{G}_2^0(\mathbb{R}^3)} \varphi dV^i = \int_{\mathbb{S}^2} \varphi(f^i, \star n^i) d\mu_{f^i} = \lim_{k \rightarrow \infty} \int_{\mathbb{S}^2 \setminus \bigcup_{j=1}^{N^i} B_{s_k}(b^{i,j})} \varphi(f_k^i, \star n_k^i) d\mu_{f_k^i}. \quad (6.35)$$

By the proof of [29, Theorem 3.3] (see also the proof of [28, Theorem 1.5]), there exist Borel sets $S_k^{i,j} \subset \mathbb{S}^2$ such that (see Equations (3.19) and (3.20) in [29])

$$\lim_{k \rightarrow \infty} \int_{S_k^{i,j}} 1 d\mu_{f_k^i} = 0 \quad \text{for all } i \in \{1, \dots, N\} \text{ and } j \in \{1, \dots, N^i\} \quad (6.36)$$

and

$$\begin{aligned} \int_{\mathbb{G}_2^{\circ}(\mathbb{R}^3)} \varphi \, dV_k &= \int_{\mathbb{S}^2} \varphi(f_k, \star n_k) \, d\mu_{f_k} = \sum_{i=1}^N \int_{\mathbb{S}^2 \setminus \bigcup_{j=1}^{N^i} B_{s_k}(b^{i,j})} \varphi(f_k^i, \star n_k^i) \, d\mu_{f_k^i} \\ &\quad + \sum_{i=1}^N \sum_{j=1}^{N^i-1} \int_{S_k^{i,j}} \varphi(f_k^i, \star n_k^i) \, d\mu_{f_k^i}. \end{aligned} \quad (6.37)$$

By (6.36) and the boundedness of φ , there holds

$$\left| \int_{S_k^{i,j}} \varphi(f_k^i, \star n_k^i) \, d\mu_{f_k^i} \right| \leq \|\varphi\|_{C^0(\mathbb{G}_2^{\circ}(\mathbb{R}^3))} \int_{S_k^{i,j}} 1 \, d\mu_{f_k^i} \rightarrow 0 \quad \text{as } k \rightarrow \infty.$$

Thus, (6.37) and (6.35) imply

$$\lim_{k \rightarrow \infty} \int_{\mathbb{G}_2^{\circ}(\mathbb{R}^3)} \varphi \, dV_k = \sum_{i=1}^N \int_{\mathbb{G}_2^{\circ}(\mathbb{R}^3)} \varphi \, dV^i = \int_{\mathbb{G}_2^{\circ}(\mathbb{R}^3)} \varphi \, dV$$

which proves (6.25).

By (6.28) there holds

$$D_0 := \sup_{k \in \mathbb{N}} \mathcal{W}(f_k) < \infty.$$

Thus, by Lemma 3.7, there exists a constant $C(A_0, D_0)$ depending only on the energy bound D_0 and the area bound A_0 in (6.23) such that

$$|\mathcal{V}_c(f_k, x) - \mathcal{V}_c(f_k, y)| \leq C(A_0, D_0) |x - y|^{1/2} \quad \text{for all } k \in \mathbb{N} \text{ and all } x, y \in \mathbb{R}^3. \quad (6.38)$$

Hence, by the varifold convergence (6.25), we can apply Lemma 3.6 and (6.31) to deduce first

$$\lim_{k \rightarrow \infty} \mathcal{V}_c(V_k, f_k(p)) = \mathcal{V}_c(V, f(p)) \quad \text{for all } p \in \mathbb{S}^2 \quad (6.39)$$

and secondly, by (6.24), the lower semi-continuity (6.32), and (6.15)

$$\mathcal{H}_{c_0}(V) + 2c_0 \mathcal{V}_c(V, x_0) < 8\pi \quad \text{for all } x_0 \in \text{spt } \mu_V.$$

Therefore, we can apply the Li–Yau inequality for general varifolds Corollary 4.4 to infer $\theta^2(\mu_V, \cdot) < 2$. Now, it follows from (6.31) that $f = f^1 \in \mathcal{F}$ and f is injective. In particular, (6.26) follows from (6.32). Moreover, by [21, Theorem 3.1], f has no branch points. That is $\log |df|_{g_0} \in L^\infty(\mathbb{S}^2)$ and thus $f \in \mathcal{E}$. It remains to show that $f \in \mathcal{S}$. Recalling that $\{x \in \mathbb{R}^3 \mid n_{E_k}(x) \neq 0\} = \partial_* E_k$ up to a set of \mathcal{H}^2 -measure zero, we see from (6.18) that $\partial_* E_k \subset f_k[\mathbb{S}^2]$ up to a set of \mathcal{H}^2 -measure zero, and thus $\mathcal{H}^2(\partial_* E_k) \leq \mathcal{A}(f_k)$. Hence, the uniform area bound (6.23) and the uniform volume bound (6.27) imply that the sequence χ_{E_k} is bounded in $BV(\mathbb{R}^3)$. Therefore, by compactness (see [15, Section 5.2, Theorem 4]), there exists an \mathcal{L}^3 -measurable set E of finite perimeter such that, after passing to a subsequence, $\chi_{E_k} \rightarrow \chi_E$ in $L^1(\mathbb{R}^3)$ and pointwise almost

everywhere as $k \rightarrow \infty$. In particular, the left hand side in (6.17) converges as $k \rightarrow \infty$. Moreover, the right hand side of (6.17) converges by (2.7) as a consequence of the varifold convergence (6.25). Noting that $\mathcal{L}^3 \llcorner E_k \rightarrow \mathcal{L}^3 \llcorner E$ as Radon measures for $k \rightarrow \infty$, we see that by (6.16) and the C^0 -convergence (6.31)

$$\text{diam spt}(\mathcal{L}^3 \llcorner E) \leq \liminf_{k \rightarrow \infty} \text{diam } f_k[\mathbb{S}^2] \leq \text{diam } f[\mathbb{S}^2].$$

Thus, $f \in \mathcal{S}$ and the proof is concluded. \square

Remark 6.9. The minimizer in [29, Theorem 1.7] has positive algebraic volume V_0 . However, in view of Example 3.8 this is in general not enough to deduce that also the concentrated volume is nonnegative. Thus, we could not apply the Li–Yau inequality Corollary 4.4 directly to the minimizer in [29, Theorem 1.7].

Theorem 6.10. *Suppose $c_0 \in \mathbb{R}$, the numbers $A_0, V_0 > 0$ satisfy the isoperimetric inequality $36\pi V_0^2 \leq A_0^3$, and there exists a minimizing sequence f_k of*

$$\bar{\eta}(c_0, A_0, V_0) := \inf\{\mathcal{H}_{c_0}(f) \mid f \in \mathcal{S}, \mathcal{A}(f) = A_0, \mathcal{V}(f) = V_0\} \quad (6.40)$$

such that

$$\begin{cases} \liminf_{k \rightarrow \infty} (\mathcal{H}_{c_0}(f_k) + 2c_0 \inf_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x)) < 8\pi & \text{if } c_0 < 0, \\ \liminf_{k \rightarrow \infty} (\mathcal{H}_{c_0}(f_k) + 2c_0 \sup_{x \in f_k[\mathbb{S}^2]} \mathcal{V}_c(f_k, x)) < 8\pi & \text{if } c_0 \geq 0. \end{cases} \quad (6.41)$$

Then the infimum is attained by a smooth embedding $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$.

Remark 6.11. (i) In view of (6.16), (6.19) and Lemma 6.4, we see that if $c_0 \leq 0$, then

$$\inf_{x_0 \in f[\mathbb{S}^2]} \mathcal{V}_c(f, x_0) \geq \frac{(2\pi)^2 \mathcal{V}(f)}{9^2(\mathcal{A}(f) + \frac{2}{3}|c_0| \mathcal{V}(f))} \frac{1}{\mathcal{H}_{c_0}(f)}$$

for all $f \in \mathcal{S}$. Thus, an elementary computation shows that (6.41) is satisfied provided

$$\bar{\eta}(c_0, A_0, V_0) < 4\pi \left(1 + \sqrt{1 + L(c_0, A_0, V_0)} \right)$$

for

$$L(c_0, A_0, V_0) := \frac{|c_0|V_0}{2 \cdot 9^2(A_0 + \frac{2}{3}|c_0|V_0)} > 0.$$

(ii) Using (6.17), (5.2) and (6.19), for all $r > 0$ and $f \in \mathcal{S}$ we have

$$\sup_{x_0 \in f[\mathbb{S}^2]} \mathcal{V}_c(f, x_0) = \sup_{x_0 \in f[\mathbb{S}^2]} \int_E |x - x_0|^{-2} d\mathcal{L}^3(x) \leq 4\pi r + r^{-2} \mathcal{V}(f).$$

Minimizing over $r > 0$ yields the estimate $\mathcal{V}_c(f, x_0) \leq 3(4\pi^2 \mathcal{V}(f))^{\frac{1}{3}}$. Thus, (6.41) is satisfied for $c_0 > 0$ provided

$$\bar{\eta}(c_0, A_0, V_0) < 8\pi - 6c_0(4\pi^2 V_0)^{\frac{1}{3}}.$$

(iii) For all $c_0 \leq 0$ and $\sigma \geq 36\pi$, there exists $\bar{A}_0, \bar{V}_0 > 0$ such that $\bar{A}_0^3/\bar{V}_0^2 = \sigma$ and $\bar{\eta}(c_0, A_0, V_0) < 8\pi$ for all $0 < A_0 < \bar{A}_0$, $0 < V_0 < \bar{V}_0$ with $A_0^3/V_0^2 = \sigma$. Indeed, in view of (1.2), this is a consequence of [38, Lemma 1].

Proof of Theorem 6.10. By (6.28), we have that

$$\sup_{k \in \mathbb{N}} \mathcal{W}(f_k) \leq C(c_0, A_0, V_0) < \infty.$$

Hence, by Lemma 6.3 applied for $c_0 = 0$, there holds $\inf_{k \in \mathbb{N}} \text{diam } f_k[\mathbb{S}^2] > 0$. Moreover, after translations, we may assume $0 \in f_k[\mathbb{S}^2]$ for all k . Therefore, we can apply Lemma 6.8 to obtain $f \in \mathcal{S}$ injective such that, after passing to a subsequence,

$$V_k \rightarrow V \quad \text{in } \mathbb{V}_2^o(\mathbb{R}^3) \text{ as } k \rightarrow \infty,$$

where V_k, V are the oriented 2-varifolds in \mathbb{R}^3 associated with f_k, f . The varifold convergence implies $\mathcal{A}(f) = A_0$ and $\mathcal{V}(f) = V_0$. Thus, by (6.26), f attains the infimum (6.40). Let $\omega \in C^\infty(\mathbb{S}^2, \mathbb{R}^3)$ and define $f_t := f + t\omega$ for $t \in \mathbb{R}$. By (6.10) and (6.14) we have

$$\begin{aligned} f_t &\rightarrow f && \text{in } W^{1,\infty}(\mathbb{S}^2; \mathbb{R}^3) \cap W^{2,2}(\mathbb{S}^2; \mathbb{R}^3) \text{ as } t \rightarrow 0, \\ df_t \wedge df_t &\rightarrow df \wedge df && \text{in } L^\infty(\mathbb{S}^2; (\wedge_2 T^* \mathbb{S}^2) \otimes \wedge_2 \mathbb{R}^3) \text{ as } t \rightarrow 0, \\ n_t &\rightarrow n && \text{in } L^\infty(\mathbb{S}^2; \mathbb{R}^3) \text{ as } t \rightarrow 0 \end{aligned}$$

and the associated varifolds converge in $\mathbb{V}_2^o(\mathbb{R}^3)$. Moreover, it follows that $f_t \in \mathcal{E}$ for $|t|$ small and $\mathcal{W}(f_t) \rightarrow \mathcal{W}(f)$, $\mathcal{H}_{c_0}(f_t) \rightarrow \mathcal{H}_{c_0}(f)$, and $\mathcal{A}(f_t) \rightarrow \mathcal{A}(f)$ as $t \rightarrow 0$. Hence, similarly as in (6.38) and (6.39), we can combine Lemma 3.6 and Lemma 3.7 to deduce that for some $\varepsilon > 0$ there holds

$$\mathcal{H}_{c_0}(f_t) + 2c_0 \mathcal{V}_c(f_t, x_0) < 8\pi \quad \text{for all } |t| < \varepsilon \text{ and } x_0 \in f_t[\mathbb{S}^2].$$

It follows by Corollary 4.4 that f_t is injective for $|t| < \varepsilon$ and thus, by Lemma 6.6 and Remark 6.7, $f_t \in \mathcal{S}$. Therefore, we can proceed as in [30] and [29] to deduce that f satisfies the Euler–Lagrange equation given in [29, Lemma 4.1]. Now, we can apply [29, Theorem 4.3] to conclude that f is smooth. \square

Theorem 1.6 is now a direct consequence.

Proof of Theorem 1.6. With $L(c_0, A_0, V_0)$ as in Remark 6.11 (i), we define

$$\Gamma(c_0, A_0, V_0) := \begin{cases} 4\pi \left(\sqrt{1 + L(c_0, A_0, V_0)} - 1 \right) & \text{if } c_0 \leq 0 \\ 6c_0 (4\pi^2 V_0)^{\frac{1}{3}} & \text{if } c_0 > 0. \end{cases}$$

Then $\Gamma(c_0, A_0, V_0) > 0$ for $c_0 \neq 0$. Let $f_k \in \mathcal{S}$ be a minimizing sequence for (6.40). By Remark 6.11(i) and (ii), the choice of Γ , and since $\bar{\eta}(c_0, A_0, V_0) \leq \eta(c_0, A_0, V_0)$ as a consequence of Lemma 6.6, we find that (6.41) is satisfied and hence the infimum (6.40) is attained by a smooth embedding $f: \mathbb{S}^2 \rightarrow \mathbb{R}^3$, which implies that f is also a minimizer for (1.9) and thus $\bar{\eta}(c_0, A_0, V_0) = \eta(c_0, A_0, V_0)$. The last part follows from Remark 6.11(iii). \square

6.4 Positive total mean curvature

We recall the following inequality due to Minkowski [27]. If $\Omega \subset \mathbb{R}^3$ is a bounded convex open subset with C^2 -boundary $\partial\Omega$, then

$$\frac{1}{2} \int_{\partial\Omega} H_{\text{sc}} \, d\mathcal{H}^2 \geq \sqrt{4\pi\mathcal{H}^2(\partial\Omega)}, \quad (6.42)$$

with equality if and only if Ω is a ball. The quantity on the left hand side of (6.42) is called *total (scalar) mean curvature*. With the help of Corollary 5.3, we can generalize (6.42) to a class of nonconvex surfaces.

Theorem 6.12. *Let $f: \Sigma \rightarrow \mathbb{R}^3$ be an immersion of an oriented closed surface Σ . If there exists $x_0 \in \mathbb{R}^3$ with $\mathcal{V}_c(f, x_0) > 0$ and $\bar{\mathcal{H}}(f) \leq 4\pi\mathcal{H}^0(f^{-1}\{x_0\})$, then we have*

$$\frac{1}{2} \int H_{\text{sc}} \, d\mu \geq \sqrt{(4\pi\mathcal{H}^0(f^{-1}\{x_0\}) - \bar{\mathcal{H}}(f)) \mathcal{A}(f)}. \quad (6.43)$$

The assumption $\mathcal{V}_c(f, x_0) > 0$ is especially satisfied if f is an Alexandrov immersion and $x_0 \in \mathbb{R}^3$ is arbitrary, see (5.1).

We would like to point out that it is possible to deduce (6.43) with the absolute value on the left hand side from the classical Li–Yau inequality for the Willmore energy. However, the question whether the total mean curvature is positive remains. In general, this has to be answered in the negative; however, under certain convexity or symmetry assumptions on the surface, the total mean curvature can be shown to be positive, cf. [10, Table 1]. In the case of Alexandrov immersions, Theorem 6.12 provides a sufficient criterion for positive total mean curvature if the CMC-deficit is not too large, depending on the concentrated volume and the multiplicity at a point.

Proof of Theorem 6.12. Set $\delta := 4\pi\mathcal{H}^0(f^{-1}\{x_0\}) - \bar{\mathcal{H}}(f) \geq 0$. By Corollary 5.3 we have

$$\delta\mathcal{A}(f) \leq 2 \int_{\Sigma} H_{\text{sc}} \, d\mu \, \mathcal{V}_c(f, x_0) - 4\mathcal{V}_c(f, x_0)^2,$$

and therefore, using Young's inequality and $\mathcal{V}_c(f, x_0) > 0$, we find

$$\frac{\int H_{\text{sc}} \, d\mu}{2} \geq \frac{\delta\mathcal{A}(f)}{4\mathcal{V}_c(f, x_0)} + \mathcal{V}_c(f, x_0) \geq \sqrt{\delta\mathcal{A}(f)}. \quad \square$$

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References

- [1] W. K. Allard. On the first variation of a varifold. *Ann. of Math. (2)*, 95:417–491, 1972.
- [2] L. Ambrosio, V. Caselles, S. Masnou, and J.-M. Morel. Connected components of sets of finite perimeter and applications to image processing. *J. Eur. Math. Soc. (JEMS)*, 3(1):39–92, 2001.
- [3] Y. Bernard, G. Wheeler, and V.-M. Wheeler. Rigidity and stability of spheres in the Helfrich model. *Interfaces Free Bound.*, 19(4):495–523, 2017.
- [4] S. Blatt. A note on singularities in finite time for the L^2 gradient flow of the Helfrich functional. *J. Evol. Equ.*, 19(2):463–477, 2019.
- [5] K. A. Brakke. *The motion of a surface by its mean curvature*, volume 20 of *Mathematical Notes*. Princeton University Press, Princeton, N.J., 1978.
- [6] K. Brazda, L. Lussardi, and U. Stefanelli. Existence of varifold minimizers for the multiphase Canham-Helfrich functional. *Calc. Var. Partial Differential Equations*, 59(3):Paper No. 93, 26, 2020.
- [7] L. E. J. Brouwer. Beweis des Jordanschen Satzes für den n -dimensionalen Raum. *Math. Ann.*, 71(3):314–319, 1911.
- [8] P. Canham. The minimum energy of bending as a possible explanation of the biconcave shape of the human red blood cell. *J. Theor. Biol.*, 26(1):61–81, 1970.
- [9] R. Choksi and M. Veneroni. Global minimizers for the doubly-constrained Helfrich energy: the axisymmetric case. *Calc. Var. Partial Differential Equations*, 48(3-4):337–366, 2013.
- [10] J. Dalphin, A. Henrot, S. Masnou, and T. Takahashi. On the minimization of total mean curvature. *J. Geom. Anal.*, 26(4):2729–2750, 2016.
- [11] H. Deuling and W. Helfrich. Red blood cell shapes as explained on the basis of curvature elasticity. *Biophysical journal*, 16(8):861–868, 1976.
- [12] S. Eichmann. The Helfrich boundary value problem. *Calc. Var. Partial Differential Equations*, 58(1):Paper No. 34, 26, 2019.
- [13] S. Eichmann. Lower semicontinuity for the Helfrich problem. *Ann. Global Anal. Geom.*, 58(2):147–175, 2020.
- [14] L. C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, second edition, 2010.
- [15] L. C. Evans and R. F. Gariepy. *Measure theory and fine properties of functions*. Studies in Advanced Mathematics. CRC Press, Boca Raton, FL, 1992.

- [16] H. Federer. *Geometric measure theory*. Die Grundlehren der mathematischen Wissenschaften, Band 153. Springer-Verlag New York Inc., New York, 1969.
- [17] W. Helfrich. Elastic properties of lipid bilayers: Theory and possible experiments. *Zeitschrift für Naturforschung C*, 28(11):693–703, 1973.
- [18] H. Hopf. *Differential geometry in the large*, volume 1000 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 1983. Notes taken by Peter Lax and John Gray, With a preface by S. S. Chern.
- [19] J. E. Hutchinson. Second fundamental form for varifolds and the existence of surfaces minimising curvature. *Indiana Univ. Math. J.*, 35(1):45–71, 1986.
- [20] L. G. A. Keller, A. Mondino, and T. Rivière. Embedded surfaces of arbitrary genus minimizing the Willmore energy under isoperimetric constraint. *Arch. Ration. Mech. Anal.*, 212(2):645–682, 2014.
- [21] E. Kuwert and Y. Li. $W^{2,2}$ -conformal immersions of a closed Riemann surface into \mathbb{R}^n . *Comm. Anal. Geom.*, 20(2):313–340, 2012.
- [22] E. Kuwert and R. Schätzle. Removability of point singularities of Willmore surfaces. *Ann. of Math. (2)*, 160(1):315–357, 2004.
- [23] P. Li and S. T. Yau. A new conformal invariant and its applications to the Willmore conjecture and the first eigenvalue of compact surfaces. *Invent. Math.*, 69(2):269–291, 1982.
- [24] P. Mattila. *Geometry of sets and measures in Euclidean spaces*, volume 44 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 1995. Fractals and rectifiability.
- [25] J. McCoy and G. Wheeler. Finite time singularities for the locally constrained Willmore flow of surfaces. *Comm. Anal. Geom.*, 24(4):843–886, 2016.
- [26] U. Menne and C. Scharrer. A novel type of Sobolev-Poincaré inequality for submanifolds of Euclidean space, 2017. arXiv:1709.05504.
- [27] H. Minkowski. Volumen und Oberfläche. *Math. Ann.*, 57(4):447–495, 1903.
- [28] A. Mondino and T. Rivière. Immersed spheres of finite total curvature into manifolds. *Adv. Calc. Var.*, 7(4):493–538, 2014.
- [29] A. Mondino and C. Scharrer. Existence and regularity of spheres minimising the Canham-Helfrich energy. *Arch. Ration. Mech. Anal.*, 236(3):1455–1485, 2020.
- [30] T. Rivière. Variational principles for immersed surfaces with L^2 -bounded second fundamental form. *J. Reine Angew. Math.*, 695:41–98, 2014.
- [31] F. Rupp. The volume-preserving Willmore flow, 2020. arXiv:2011.14904.

- [32] F. Rupp. The Willmore flow with prescribed isoperimetric ratio, 2021. arXiv:2106.02579.
- [33] T. Sakai. *Riemannian geometry*, volume 149 of *Translations of Mathematical Monographs*. American Mathematical Society, Providence, RI, 1996. Translated from the 1992 Japanese original by the author.
- [34] C. Scharrer. Embedded Delaunay tori and their Willmore energy, 2021. arXiv:2102.04947.
- [35] C. Scharrer. On the minimisation of bending energies related to the Willmore functional under constraints on area and volume, 2021. PhD thesis, University of Warwick.
- [36] C. Scharrer. Some geometric inequalities for varifolds on Riemannian manifolds based on monotonicity identities. *Ann. Global Anal. Geom.*, 2022.
- [37] R. Schätzle. Lower semicontinuity of the Willmore functional for currents. *J. Differential Geom.*, 81(2):437–456, 2009.
- [38] J. Schygulla. Willmore minimizers with prescribed isoperimetric ratio. *Arch. Ration. Mech. Anal.*, 203(3):901–941, 2012.
- [39] U. Seifert, K. Berndt, and R. Lipowsky. Shape transformations of vesicles: Phase diagram for spontaneous-curvature and bilayer-coupling models. *Physical review A*, 44(2):1182, 1991.
- [40] L. Simon. Existence of surfaces minimizing the Willmore functional. *Comm. Anal. Geom.*, 1(2):281–326, 1993.
- [41] T. J. Willmore. *Riemannian geometry*. Oxford Science Publications. The Clarendon Press, Oxford University Press, New York, 1993.
- [42] O.-Y. Zhong-Can and W. Helfrich. Instability and deformation of a spherical vesicle by pressure. *Physical review letters*, 59(21):2486, 1987.