

# **KAUFMANN BRACKET OF PLANE CURVES**

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# KAUFFMAN BRACKET OF PLANE CURVES

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There exists a straightforward way to get an invariant of an immersed cooriented hypersurface  $C$  in a smooth manifold  $N$ . We lift  $C$  to the manifold  $M$  of cooriented contact elements of  $N$ . This gives us an embedded submanifold  $L_C$ . Now we take the value of a known invariant of embeddings on  $L_C \hookrightarrow M$  as the invariant of our initial immersion  $C \looparrowright N$ .

The manifold  $M$  of cooriented contact elements is the spherisation of the cotangent bundle of  $N$ :  $M = ST^*N$ . It has a natural contact structure. Our lifting  $L_C$  is a Legendrian submanifold with respect to this structure. The hypersurface  $C$  is called *the front* of  $L_C$ . The above procedure defines an invariant not only on immersed  $C \looparrowright N$  but also on submanifolds with some “admissible” singularities which may appear as singularities of fronts of smooth Legendrian submanifolds generically embedded into  $M$ .

The simplest situation is  $N = \mathbf{R}^2$ . The “admissible” singularities in this case are cusps. Thus we can induce an invariant on collections of closed oriented and cooriented plane curves which may have only double points and cusps as singularities. The manifold  $M$  of contact elements of the plane is the solid torus  $M = \mathbf{R}^2 \times S^1$ . So the lifted submanifolds are Legendrian links in it. This general approach was used in [11] to define an invariant of an immersed plane curve. There a Kontsevich type integral [10] was taken as a known invariant of knots in a solid torus. In the similar way, a polynomial invariant of knots in a solid torus defined in [1] was lowered to plane curves in [2].

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In this paper we take the Kauffman bracket for links in a solid torus (see [15]) as a known invariant to be induced on plane fronts. It turns out that it can be entirely defined in terms of front  $C$  itself without using the Legendrian lifting. The Kauffman bracket is a polynomial in two variables  $A$  and  $h$ , Laurent in  $A$  and ordinary in  $h$ . We show that, after the substitution  $A = e^t$  and Taylor expansion in a power series in  $t$ , the coefficient at  $t^n$  is an invariant of Arnold's  $J^+$ -theory [3, 4] of order at most  $n$  in Vassiliev sense. These coefficients are polynomials in  $h$ . We calculate the corresponding symbols (weight systems). As a corollary we obtain that the first coefficient is essentially the quantum deformation of the Bennequin invariant introduced recently by M.Polyak [18].

For an application of the same general idea to induce order 1 invariants in a higher-dimensional situation see [12].

## 1 Definitions and known results

In this section we recall some basic facts about our curves, corresponding Legendrian links and their invariants. See [3, 4] for more details.

### 1.1 Legendrian links and their fronts

A *contact element* at a point of a plane is a line in the tangent plane. Its *coorientation* is a choice of one of two half-planes into which it divides the tangent plane. The manifold  $M = ST^*\mathbf{R}^2$  of all cooriented contact elements of the plane is diffeomorphic to the solid torus  $\mathbf{R}^2 \times S^1$ , since the coorienting normal vector is defined by its angle  $\phi$ . Manifold  $M$  has a natural contact structure defined as zeros of the form  $(\cos \phi)dx + (\sin \phi)dy$ , where  $(x, y)$  are coordinates on  $\mathbf{R}^2$ . A *Legendrian link*  $L$  in  $M$  is an embedding of a number of oriented circles into  $M$  tangent to the contact planes at each of its points. A Legendrian link has a natural framing by transversals to the contact planes. The canonical projection of  $L$  to  $\mathbf{R}^2$  gives a collection of plane curves. We call it *the front* of  $L$ . It has an orientation (coming from  $L$ ) and a coorientation (the coordinate  $\phi$  forgotten by the projection defines not only the line tangent to the front, but the side of this line as well). A generic front may have only transverse double points and cusps as singularities. We call such a front *a normal front*. Since a front is cooriented the number

of cusps on each component of a normal front is even.

Any cooriented plane curve  $C$  lifts to a Legendrian curve  $L_C \in M$  by taking the cooriented tangent direction as a contact element at each point of  $C$ . The lifting of a collection of curves with normal front singularities is a link (Fig.1).

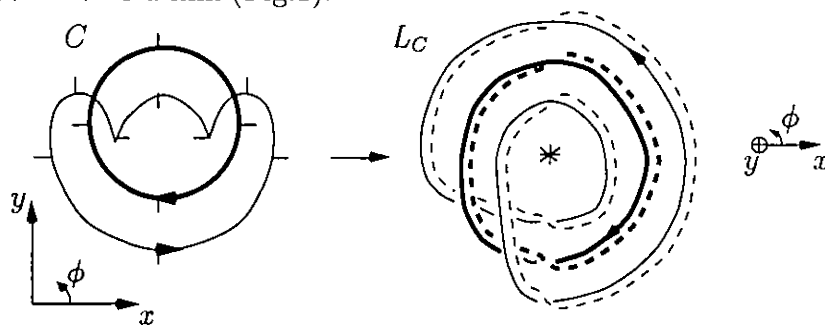


Figure 1. Legendrian lifting.

Left picture: A normal front  $C$  with two components;  $(x, y)$  are coordinates on  $\mathbb{R}^2$ ;  $C$  lies in the halfplane  $x > 0$ .

Right picture: Framed Legendrian link  $L_C$  is drawn as a diagram of the projection to the punctured plane with polar coordinates  $(x, \phi)$ ; the  $y$ -axis is perpendicular to the plane and directed from the reader.

## 1.2 Index, Maslov index and perestroikas

To each component  $C_j$  of a normal front  $C = \cup C_j$  we assign two integers, *index*<sup>1</sup>  $\text{ind}(C_j)$  and *Maslov index*  $\mu(C_j)$ .  $\text{ind}(C_j)$  is the number of full rotations made by the coorienting vector as it moves along  $C_j$ .  $\mu(C_j)$  is the difference between the numbers of positive and negative cusps of  $C_j$ . A cusp is called *positive* if the 1-form which coorients the curve at the cusp point is positive on the neighbouring orienting vectors and *negative* otherwise (Fig.2).



Figure 2. Negative and positive cusps.

The number  $\mu(C_j)$  is always even. Reversing of the orientation of  $C_j$  changes the signs of both  $\text{ind}(C_j)$  and  $\mu(C_j)$ . Reversing of the coorientation of  $C_j$  changes only the sign of  $\mu(C_j)$ .

<sup>1</sup>Other names are *winding number*, *rotation number*, *Whitney index*.

There are four types of generic degenerations of a normal front. We show them in Fig.3 in perestroikas in generic 1-parameter families.

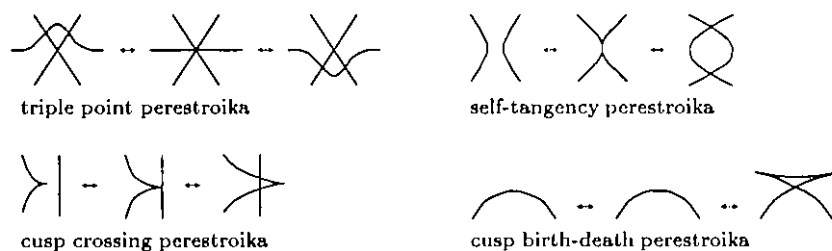


Figure 3. Perestroikas.

**Theorem** [13] (see also [3, 4]). *The collection of pairs  $(\text{ind}(C_j), \mu(C_j))$  is a complete invariant of a normal front  $C = \cup C_j$  under plane isotopies and the four types of perestroikas of Fig.3.*

For each pair of integers  $r \geq 1$  and  $s \geq 0$  let  $K_{r,s}$  be the curve of Fig.4.  $\text{ind}(K_{r,s}) = r - 1$ ;  $\mu(K_{r,s}) = 2s$ .

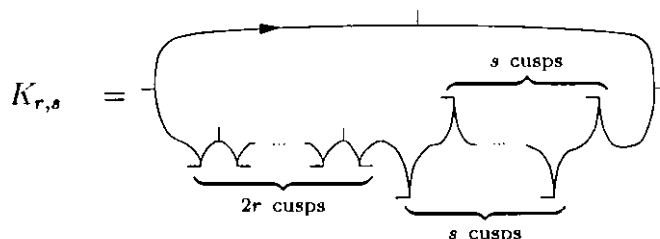


Figure 4. Canonical curves. <sup>2</sup>

The theorem says that each component  $C_j$  of a normal front can be transformed to one of the curves  $K_{r,s}$  (possibly with changed orientation or coorientation or both) by a sequence of the perestroikas and isotopies of the plane. Fig.5 provides an example.

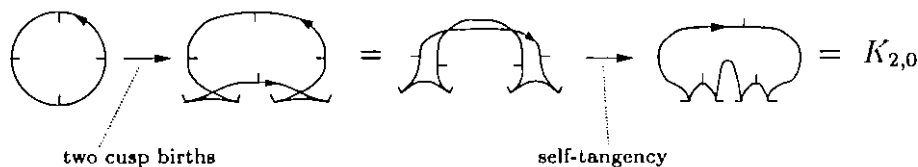


Figure 5. Transformation of the circle to  $K_{2,0}$ .

<sup>2</sup>Our choice of canonical curves slightly differs from Arnold's one [3].

### 1.3 $J^+$ -type invariants

It is convenient to subdivide self-tangency perestroikas into the following four types according to the orientations and coorientations. A self-tangency is called *dangerous* if both the tangent branches are cooriented by the same half-plane and *safe* otherwise. A self-tangency is called *direct* if both tangent branches are oriented by the same tangent vector and *inverse* otherwise.

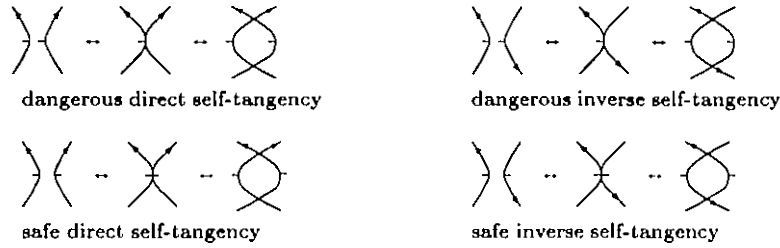


Figure 6. Four types of self-tangencies.

Note that, if two tangent branches belong to the same component of a front, the property of the tangency point to be direct or inverse (resp. dangerous or safe) does not depend on the orientation (resp. coorientation) of the component.

It is easy to see that the topological type of a Legendrian link  $L_C$  in the solid torus  $M$  does not change under all the perestroikas except dangerous self-tangencies. A dangerous self-tangency perestroika corresponds to an interchanging of overcrossing and undercrossing in a link diagram like that in Fig.1. But not all interchangings can be done in the class of Legendrian links, and so in the class of corresponding fronts.

We will say that two fronts are  $J^+$ -equivalent if one can be transformed to another without dangerous self-tangencies. Fig.5 shows a  $J^+$ -equivalence of the circle to  $K_{2,0}$ . Similarly one can show that the circle with the opposite coorientation is also  $J^+$ -equivalent to  $K_{2,0}$ . Another example is  $J^+$ -equivalence of figure-eight curves with different choices of orientation and coorientation (Fig.7).

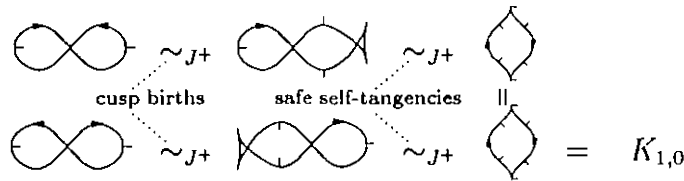


Figure 7.  $J^+$ -equivalence of figure-eight curves.

By a  $J^+$ -type invariant we mean an invariant of normal fronts which does not change under all the perestroikas except dangerous self-tangencies. The first example of such an invariant was an invariant introduced by V.I.Arnold in [3, 4] and named  $J^+$ . This is an invariant of a one component front defined by its values on the canonical curves:

$$J^+(K_{r,s}) = -s \quad (\text{for any choice of the orientation and coorientation})$$

and by its behavior under the dangerous self-tangency perestroikas:

$$J^+(\text{crossing}) - J^+(\text{self-tangency}) = 2; \quad J^+(\text{self-tangency}) - J^+(\text{crossing}) = 2.$$

According to Theorem 1.2 this data is sufficient for calculating  $J^+$  on any normal front. Here is an example.

$$\begin{aligned} J^+(\text{two cusp crossings}) &= J^+(\text{dangerous self-tangency}) = J^+(\text{self-tangency}) - 2 \\ &= J^+(\text{two cusp births}) - 2 = J^+(\text{two safe self-tangencies}) - 2 \\ &= J^+(K_{3,0}) - 2 = -2. \end{aligned}$$

There are several combinatorial formulas for calculating the values of  $J^+$  on curves without cusps (see a review in [8]) and Polyak's formula [17] for curves with cusps.

In Vassiliev sense  $J^+$  is an invariant of order 1.

**Remark.** Reversing orientations of both the local branches in the two dangerous self-tangency perestroikas of Fig.6, one obtains two more dangerous perestroikas which look different from those above. But their behaviour in all our constructions is absolutely identical to the behaviour of the corresponding "twins". So we spell all the formulas involving dangerous self-tangencies only for the two perestroikas of Fig.6.



## 1.4 The Bennequin invariant and its quantization

For a Legendrian knot  $K$  in a contact  $\mathbf{R}^3$  Bennequin [6] defined a self-linking number  $\beta$  as the linking number of  $K$  with a small shift of  $K$  in a direction everywhere transversal to the contact planes. This definition was generalized to a non simply-connected case of Legendrian knots in the solid torus  $ST^*\mathbf{R}^2$ , with its standard contact structure, by S.Tabachnikov [19].

As the usual linking number [16] the Bennequin–Tabachnikov invariant can be read from a diagram of a knot and its framing like that in Fig.1. For example, for the bold component of the link of Fig.1  $\beta = 1$  (we have two positive crossings of the projections of the knot and its framing), for the thin component  $\beta = 3$  (we have six positive crossings there).

Arnold proved [3] that  $\beta = 1 - J^+$ . So any combinatorial formula for  $\beta$  gives a formula for  $J^+$  and vice versa. Several such formulas are in [19] (see also [9]).

M.Polyak [18] invented the following state sum formula for the Bennequin–Tabachnikov invariant  $\beta$ . To each crossing  $p$  of a one component normal front  $C$  we attach the sign  $\sigma(p) = +1$ , if the pairs (orienting vector, coorienting vector) for the two intersecting branches give the same orientation of the plane, and  $\sigma(p) = -1$  otherwise. According to this sign we split  $C$  at  $p$  respecting the orientation and coorientation (Fig.8).

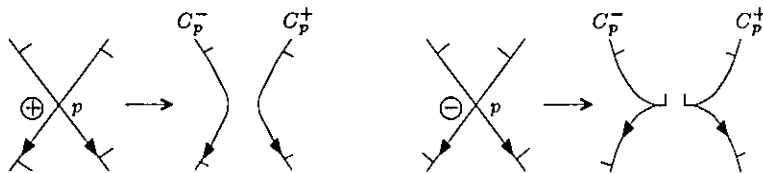


Figure 8. Splittings of a front  $C$  at a crossing saving the orientation and coorientation.

In fact this is a unique natural splitting which gives two component curves with two branches near  $d$  belonging to different components. Denote by  $C_p^-$  (resp.  $C_p^+$ ) the component that contains the left (resp. right) branch assuming both branches oriented downwards (see Fig.8).

**Theorem** [18]. *Let  $C$  be a one component normal front. Denote by  $S$  the state sum*

$$S = \sum_p (\text{ind}(C_p^+) - \text{ind}(C_p^-) - \sigma(p))$$

*over the set of all double points of  $C$ . Then*

$$\beta(L_C) = S - (\text{ind}(C) - 1)n^+ + (\text{ind}(C) + 1)n^- + \text{ind}^2(C),$$

*where  $n^+$  (resp.  $n^-$ ) is half the number of cusps of  $C$  whose neighbourhoods give a positive (resp. negative) contribution to the index of  $C$ .*

This formula admits a quantum deformation [18]. Let  $q$  be a formal quantum parameter and

$$[n]_q = \frac{q^n - q^{-n}}{q - q^{-1}} \in \mathbf{Z}[q, q^{-1}]$$

the corresponding quantum integer.

**Theorem** [18]. *Let  $S_q = \sum[(\text{ind}(C_p^+) - \text{ind}(C_p^-) - \sigma(p))]_q$  be a quantum state sum. Then*

$$\beta_q(L_C) = S_q - [ \text{ind}(C) - 1 ]_q n^+ + [ \text{ind}(C) + 1 ]_q n^- + [ \text{ind}(C) ]_q \text{ind}(C)$$

*is a  $J^+$ -type invariant of a one component normal front  $C$  such that  $\beta_1(L_C) = \beta(L_C)$ .*

**Remark.** The definition of  $\beta_q$  is easily seen to be independent from orientation and coorientation of a normal front.

Taking our canonical curves  $K_{r,s}$  with the orientations as in Fig.4 we get  $n^+ = (2r + s)/2$ ,  $n^- = s/2$  and  $\text{ind}(C) = r - 1$ . So

$$\beta_q(L_{K_{r,s}}) = -(r + \frac{s}{2})[r - 2]_q + \frac{s}{2}[r]_q + (r - 1)[r - 1]_q .$$

Let us describe the behavior of the quantum Bennequin invariant under dangerous self-tangencies. First of all we define *an index  $i_d$  of a self-tangency point  $d$*  which appears during such a perestroika of a normal front. To do this we split the self-tangency point respecting the orientation and coorientation as shown in Fig.9.



Figure 9. Splittings of dangerous self-tangencies saving the orientation and coorientation.

We obtain two curves. Let  $i'$  and  $i''$  be their indices. We set  $i_d = |i' - i''|$ . The jumps of  $\beta_q$  under dangerous self-tangencies of the index  $i_d$  are:

$$\beta_q(\text{self-tangency with arrows out}) - \beta_q(\text{crossing}) = q^{i_d} + q^{-i_d}; \quad \beta_q(\text{crossing}) - \beta_q(\text{self-tangency with arrows in}) = q^{i_d} + q^{-i_d}.$$

These formulas show that  $\beta_q$  is an invariant of order 1 in Vassiliev sense.

## 2 Kauffman bracket

In this section we define the Kauffman bracket and prove its uniqueness. The bracket does not depend on the orientations of curves of a collection.

For a framed link in a solid torus the Kauffman bracket was defined in [15]. Its values belong to  $\mathbf{Z}[A^{\pm 1}, h]$ . Using the Legendrian lifting we can define  $\langle C \rangle = \langle L_C \rangle$ . This is a  $J^+$ -type invariant of a normal front  $C$ . We call it *the Kauffman bracket* of  $C$ .

### 2.1 Main result

**Theorem 1.** *There exists a unique  $J^+$ -type invariant  $\langle C \rangle \in \mathbf{Z}[A^{\pm 1}, h]$  of a normal front  $C$  satisfying the following properties:*

- 1)  $\langle \text{self-tangency with arrows out} \rangle = A^{-1} \langle \text{crossing} \rangle - A^{-2} \langle \text{self-tangency with arrows in} \rangle;$
- 2)  $\langle \infty \rangle = -A^3;$
- 3)  $\langle \bigcirc \rangle = -A^3 h;$
- 4)  $\langle K_1 \cdot K_2 \rangle = -(A^2 + A^{-2}) \langle K_1 \rangle \cdot \langle K_2 \rangle,$   
for  $K_1 \neq \emptyset, K_2 \neq \emptyset.$

Here  $K_1 \cdot K_2$  is a collection of two fronts  $K_1$  and  $K_2$  which lie in different half-planes with respect to a certain line in  $\mathbf{R}^2$ .

**Remarks. 1.** After the Legendrian lifting (see Fig.1) the fragments of links corresponding to the fronts of property 1) have the following diagrams in the  $(x, \phi)$ -plane:

$$\langle \rangle \langle \rangle = A^{-1} \langle \times \rangle - A^{-2} \langle \bowtie \rangle .$$

So property 1) is just the usual skein relation for the Kauffman bracket. All other properties also correspond to the usual properties of the Kauffman bracket in the solid torus (see [15]). So the existence of such a bracket of normal fronts follows directly from [15].

**2.** For calculation of the Kauffman bracket we will use the fact that the curve  $K_{1,0}$  can be moved through other curves of a front. For example

two cusp crossings

So, if one of the components of our front is the  $K_{1,0}$  with nothing inside, we can transfer it far away from everything else and apply property 4) of Theorem 1:

$$\langle \bigcirc \rangle = -A^6 h (A^2 + A^{-2}) = -(A^4 + A^8) h .$$

**3.** One more useful fact is that two circles with opposite coorientations are  $J^+$ -equivalent (see sec.1.3 and Fig.5). So their brackets are equal.

**Proposition 1.** For any  $J^+$ -type invariant satisfying properties 1)-4) of Theorem 1 the following equalities hold:

$$\langle \times \rangle = A \langle \bigcirc \rangle - A^2 \langle \bowtie \rangle ;$$

$$\langle \rangle \langle \rangle - \langle \bigcirc \rangle = (A - A^{-1}) \left( \langle \times \rangle - \langle \bowtie \rangle \right) .$$

**Proof of Proposition 1.**

$$\begin{aligned}
\langle \text{Diagram 1} \rangle &= A^{-1} \langle \text{Diagram 2} \rangle - A^{-2} \langle \text{Diagram 3} \rangle \\
&= A^{-2} \langle \text{Diagram 4} \rangle - A^{-3} \langle \text{Diagram 5} \rangle - A^{-3} \langle \text{Diagram 6} \rangle + A^{-4} \langle \text{Diagram 7} \rangle \\
&= A^{-1} \left( A^{-1} \langle \text{Diagram 8} \rangle - A^{-2} \langle \text{Diagram 9} \rangle \right) \\
&\quad + \left( -A^{-3} \langle \text{Diagram 10} \rangle + A^{-4} (-A^3) (-A^2 - A^{-2}) \langle \text{Diagram 11} \rangle \right) \\
&= A^{-1} \langle \text{Diagram 12} \rangle + (-A^{-3} + A + A^{-3}) \langle \text{Diagram 13} \rangle.
\end{aligned}$$

This implies the first equality of the proposition. The second one follows from it and property 1) of Theorem 1.

## 2.2 Useful lemmas

In lemmas below we prove some relations for the Kauffman bracket which follow from properties 1)-4). We will use these relations in sec.2.3 to prove the uniqueness of the Kauffman bracket.

The relations hold for both possible coorientations of the fragment involved. Therefore we do not indicate its coorientation. The coorientation of the extra circular component also does not matter due to Remark 3 above.

**Lemma 1.**  $\langle \text{Diagram 14} \rangle = A \langle \text{Diagram 15} \rangle - A^2 \langle \text{Diagram 16} \rangle.$

**Proof of Lemma 1.**  $\langle \text{Diagram 16} \rangle = \langle \text{Diagram 17} \rangle = A^{-1} \langle \text{Diagram 15} \rangle - A^{-2} \langle \text{Diagram 14} \rangle.$   
safe self-tangency and cusp death

**Lemma 2.**  $\langle \text{Diagram 18} \rangle = \langle \text{Diagram 19} \rangle = (-A^3) \langle \text{Diagram 20} \rangle.$

**Proof of Lemma 2.**

$$\begin{aligned}
 \langle \text{cusp crossing} \rangle &= \langle \text{cusp crossing} \rangle = A^{-1} \langle \text{cusp crossing} \rangle - A^{-2} \langle \text{cusp crossing} \rangle \\
 &= (A^{-1} - A^{-2}(A^5 + A)) \langle \text{---} \rangle = (-A^3) \langle \text{---} \rangle. \\
 &\text{cusp death and properties 4),2)}
 \end{aligned}$$

The proof of the second equality is similar.

**Lemma 3.**  $\langle \text{---} \rangle = A^{-1} \langle \text{---} \rangle - A^4 \langle \text{---} \rangle.$

**Proof of Lemma 3.**

$$\langle \text{---} \rangle = \langle \text{---} \rangle = A^{-1} \langle \text{---} \rangle - A^{-2} \langle \text{---} \rangle.$$

safe self-tangency and cusp death

The last term is equal to  $A^{-2}(-A^3)^2 \langle \text{---} \rangle = A^4 \langle \text{---} \rangle$  by Lemma 2.

**Lemma 4.**  $\langle \text{---} \rangle = (A^{-1} - A^3) \langle \text{---} \rangle - A^2 \langle \text{---} \rangle.$

**Proof of Lemma 4.**

$$\begin{aligned}
 \langle \text{---} \rangle &= A \langle \text{---} \rangle + A^{-1} \langle \text{---} \rangle \\
 &= A \langle \text{---} \rangle + A^{-1} \langle \text{---} \rangle \\
 &\text{cusp birth and safe self-tangency} \\
 &= A^2 \langle \text{---} \rangle + (A^{-1} - A^3) \langle \text{---} \rangle. \\
 &\text{Lemma 1}
 \end{aligned}$$

### 2.3 Proof of Theorem 1

To prove the uniqueness of the bracket it is enough to show that properties 1)-4) are sufficient for calculation of the bracket on any normal front. We prove this giving an algorithm for such calculation.

First of all we eliminate all double points of the front using the skein relation 1). We obtain a linear combination of brackets of fronts without double points. Each of these fronts is just a union of "ovals" which can have cusps and be nested. Using Lemma 2 we cancel pairs of neighbouring cusps with opposite directions. Then using Lemma 1 we invert the directions of pairs of cusps from inside to outside of their "oval".

After that we reduce the number of cusps on each "oval" to zero or two (Lemma 3). Now consider the deepest "ovals" of the nests. We transfer all those which are  $K_{1,0}$ -curves far away (see Remark 2) reducing our computation to the computation of the bracket of the remaining part. We have left only circles on the deepest level. We decrease their depth by Lemma 4. This brings us to the beginning of this paragraph with the depth of the nests reduced by 1. Theorem 1 is proved.

**Corollary.** *The Kauffman bracket does not distinguish between two fronts which differ by the simultaneous change of coorientations of all the components.*

Corollary follows directly from the proof of Theorem 1.

**Example.** The Kauffman bracket of canonical curves.

$$\begin{aligned} \langle K_{r,s} \rangle &= (-A^3)^s \langle K_{r,0} \rangle \text{ by Lemma 2;} \\ \langle K_{1,0} \rangle &= -A^3 \text{ by property 2);} \\ \langle K_{2,0} \rangle &= -A^3 h \text{ by property 3).} \end{aligned}$$

For  $r > 2$ ,  $\langle K_{r,0} \rangle$  can be computed recurrently:

$$\begin{aligned} \langle K_{r,0} \rangle &= \left\langle \underbrace{\text{---} \underbrace{\text{---} \text{---} \text{---} \text{---}}_{2r \text{ cusps}} \text{---}} \right\rangle \\ &\stackrel{\text{Lemma 3}}{=} A^{-1} \left\langle \underbrace{\text{---} \underbrace{\text{---} \text{---} \text{---} \text{---}}_{2(r-1) \text{ cusps}} \text{---}} \right\rangle - A^{-4} \langle K_{r-2,0} \rangle \\ &\stackrel{\text{properties 4) and 3)}}{=} (A^4 + 1)h \langle K_{r-1,0} \rangle - A^4 \langle K_{r-2,0} \rangle . \end{aligned}$$

Setting  $A = 1$  we get  $\langle K_{r,0} \rangle |_{A=1} = -T_{r-1}(h)$ , where the  $T_n(h)$  are the classical Tchebyshev polynomials in  $h$ :  $T_n(\cos x) = \cos(nx)$ .

So the negative of  $\langle K_{r,0} \rangle$  can be considered as a deformation of the Tchebyshev polynomial with the parameter  $A$ . The number of the polynomial is the absolute value of the index of the canonical curve.

### 3 Taylor coefficients

In this section we prove an analog of Birman-Lin theorem [7] for the Kauffman bracket of a normal front and calculate the symbols of Taylor coefficients of the bracket as functions on marked chord diagrams.

#### 3.1 Finite order $J^+$ -type invariants

The extension of a knot invariant to degenerate knots with double points is basic for the Vassiliev theory. In a similar way any  $J^+$ -type invariant  $f$  recursively extends to fronts with a finite number of dangerous self-tangencies:

$$f(\text{X}) = f(\text{O}) - f(\text{>})\langle \text{<} \rangle; \quad f(\text{>})\langle \text{<} \rangle = f(\text{>})\langle \text{<} \rangle - f(\text{O}).$$

These rules are due to the natural coorientation of the strata of dangerous self-tangencies from [3]. When lifted to  $ST^*\mathbf{R}^2$  both rules are in fact the definition of an extended invariant of the original Vassiliev theory.

Following the above rules we get the extension of the Kauffman bracket of plane fronts. One should note that, though the Kauffman bracket of a normal front does not depend on orientations of its components, the extended bracket does depend on these orientations if a degenerate front has more than one component.

**Definitions.** A  $J^+$ -type invariant  $f$  has *order  $n$  in Vassiliev sense* if  $n$  is the maximal number of dangerous self-tangencies of a front on which the extension of  $f$  does not vanish. *The symbol* of such  $f$  is the restriction of  $f$  to the set of fronts with precisely  $n$  dangerous self-tangencies.

Gromov's theorem (sec.1.2) means that indices and Maslov indices of the components are the only invariants of order zero.

The difference of two invariants of order  $n$  with the same symbol is an invariant of order less than  $n$ .



**Theorem 2.** Set  $A = e^t$  in the Kauffman bracket of a plane front  $C$  and expand the result in a power series in  $t$ . Then the coefficient at  $t^n$  in the series  $\langle C \rangle |_{A=e^t}$  is a  $J^+$ -type invariant of order at most  $n$  in Vassiliev sense.

**Proof of Theorem 2.** Let  $C$  be a front with  $n+1$  dangerous self-tangency points  $d_1, \dots, d_{n+1}$ . We consider two splittings of  $C$  (Fig.10) at a point  $d_i$  and attach to each of the splittings a sign  $\epsilon(d_i)$  which indicates either agreement or disagreement of the surgery with the orientations.

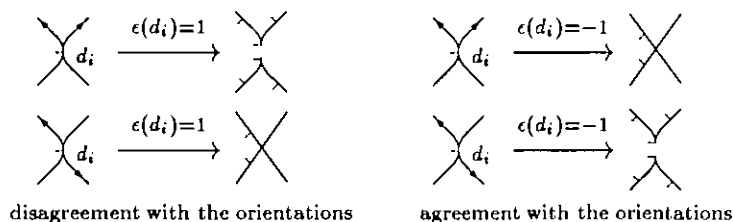


Figure 10. Signs of splittings of dangerous self-tangencies.

Let  $C_{\epsilon_1, \dots, \epsilon_{n+1}}$  be the splitting of  $C$  at all the points  $d_1, \dots, d_{n+1}$  with the signs  $\epsilon(d_i) = \epsilon_i$ . The second equality of Proposition 1 of sec.2.1 implies

$$\langle C \rangle = (A - A^{-1})^{n+1} \sum_{\epsilon_1, \dots, \epsilon_{n+1}} \epsilon_1 \cdot \dots \cdot \epsilon_{n+1} \cdot \langle C_{\epsilon_1, \dots, \epsilon_{n+1}} \rangle,$$

where the sum is taken over all  $2^{n+1}$  possible splittings of the self-tangency points. The substitution  $A = e^t$  and Taylor expansion provide

$$(A - A^{-1})^{n+1} = (2t)^{n+1} + \text{terms of higher degree}.$$

Therefore the coefficient at  $t^n$  in  $\langle C \rangle |_{A=e^t}$  is equal to zero. Theorem 2 is proved.

**Remark.** The proof demonstrates a bit more than the theorem claims. Namely, evaluations at  $A = 1$  and  $A = -1$  of the  $n$ th derivative of the Kauffman bracket with respect to  $A$  turn out to be invariants of order at most  $n$  (cf. [7]). The exponential substitution is a sort of tradition introduced in [7].

### 3.2 Symbols of the coefficients

There are several ways to define a chord diagram of a degenerate front. Say, one can follow the approach of [11] marking chords. But the way which looks most convenient for the study of the Kauffman bracket is as follows.

Consider an oriented  $l$ -component front  $C$  with  $n$  dangerous self-tangencies. Up to an isotopy of the ambient plane we can assume that the coorienting vector at each of the self-tangency points is horizontal and directed to the left. Take  $l$  disjoint circles  $S_1^1 \cup \dots \cup S_l^1$  oriented counter-clockwise in a plane. Consider the front  $C$  as the image of a mapping  $S_1^1 \cup \dots \cup S_l^1 \rightarrow \mathbf{R}^2$ . Connect the two preimages of a direct (resp. inverse) dangerous self-tangency by a solid (resp. dashed) chord. Orient this chord from the inverse image of the right-hand branch of the self-tangency to that of the left-hand one. Mark an arc of a circle between two neighbouring endpoints of chords by a pair of integers  $(i, \mu)$ , where  $i$  is the contribution of this arc to the index of  $C$  and  $\mu$  is its contribution to the Maslov index of  $C$ .

The obtained chord diagram considered up to orientation-preserving diffeomorphisms of the circles  $S_1^1, \dots, S_l^1$  is called *the marked chord diagram* of the front  $C$  and denoted by  $D_C$ . Any abstract marked chord diagram is easily seen to be the marked chord diagram of an appropriate front.

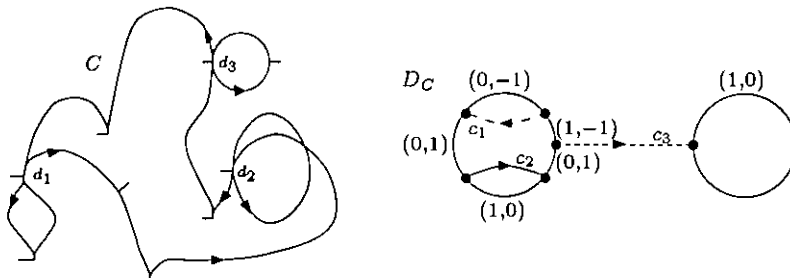


Figure 11. A front with three dangerous self-tangencies and its marked chord diagram.

Gromov's theorem of sec.1.2 implies that two fronts  $C_0$  and  $C_1$  with the same marked  $n$ -chord diagram are related by a homotopy  $\{C_t\}_{t \in [0,1]}$  in which any front  $C_t$  has  $n$  dangerous self-tangencies except a finite number of instants  $t$  when  $C_t$  gets  $n + 1$  dangerous self-tangencies (cf. [11]). Thus the symbol of an invariant of order  $n$  defines a function on marked chord diagrams with  $n$  chords. The main result of this section

(Proposition 2) is a description of this function, denoted by  $\langle D_C \rangle_n$ , for the symbol of the coefficient  $\langle C \rangle_n$  at  $t^n$  of  $\langle C \rangle|_{A=e^t}$  (cf. sec.6.3 of [5]).

To formulate the statement we redraw Fig.10 in terms of diagrams. Order chords  $c_1, \dots, c_n$  of an abstract marked  $n$ -chord diagram  $D$  in an arbitrary way. Define two signed splittings of a chord  $c_i$  as shown in Fig.12. In each of the cases the chord is substituted by two oriented marked arcs. As it will become obvious from what follows, these splittings of  $c_i$  correspond exactly to the similarly signed splittings of the self-tangency point  $d_i$  in Fig.10 if the chord represents the point in the marked diagram of a front.



Figure 12. Splittings of chords, their signs and markings.

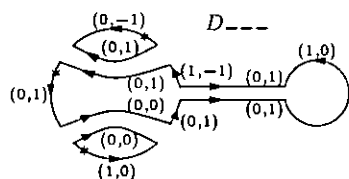
Let  $D_{\epsilon_1, \dots, \epsilon_n}$  be a splitting of the diagram  $D$  with signs  $\epsilon(c_i) = \epsilon_i$ . We denote by  $l(\epsilon_1, \dots, \epsilon_n)$  the number of components of  $D_{\epsilon_1, \dots, \epsilon_n}$ :  $D_{\epsilon_1, \dots, \epsilon_n} = \cup_j (D_{\epsilon_1, \dots, \epsilon_n})_j$ . Each component consists of oriented marked arcs which are either arcs of circles of  $D$  or the results of splittings of chords of  $D$ . For a component  $(D_{\epsilon_1, \dots, \epsilon_n})_j$  we define two integers, *index*  $i_j$  and *Maslov index*  $\mu_j$ , as follows (see Fig.13). Let us walk along the component  $(D_{\epsilon_1, \dots, \epsilon_n})_j$  and sum markings  $(i, \mu)$  of the arcs we visit with appropriate signs. Walking along an arc oriented in (resp. opposite to) the direction of our journey we take its index  $i$  and Maslov index  $\mu$  with the sign plus (resp. minus). Of course,  $i_j$  and  $\mu_j$  change their signs for the trip in the opposite direction. But the statement below does not depend on these signs.

**Proposition 2.** *The value  $\langle D \rangle_n$  of the  $n$ th coefficient of the Kauffman bracket on a marked chord diagram  $D$  is given by the formula*

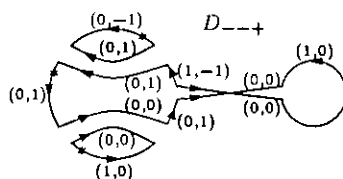
$$\langle D \rangle_n = -2^n \sum_{\epsilon_1, \dots, \epsilon_n} 2^{l(\epsilon_1, \dots, \epsilon_n) - 1} \cdot \epsilon_1 \cdot \dots \cdot \epsilon_n \cdot \prod_{j=1}^{l(\epsilon_1, \dots, \epsilon_n)} (-1)^{\mu_j / 2} T_{|i_j|}(h),$$

where the sum is taken over all  $2^n$  possible splittings of  $D$ , the product is taken over all components of a splitting,  $T_n(\cos x) = \cos(nx)$  are the classical Tchebyshev polynomials.

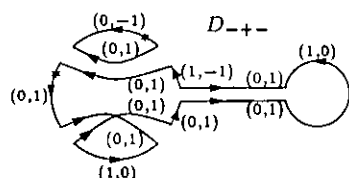
**Example.** For the marked chord diagram  $D$  of Fig.11 we have the eight splittings shown in Fig.13.



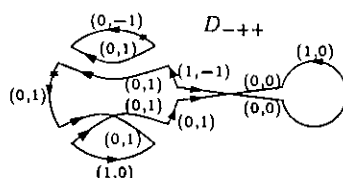
$$\begin{aligned} i_1 &= 0-0=0 & i_3 &= 1-0=1 \\ \mu_1 &= -1-1=-2 & \mu_3 &= 0-0=0 \\ i_2 &= 0+0+0+0+1-0+1+0=2 \\ \mu_2 &= 1+0+1+1+0-1-1+1=2 \end{aligned}$$



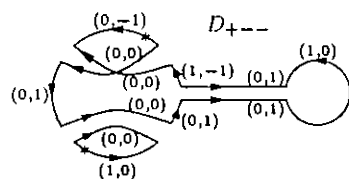
$$\begin{aligned} i_1 &= 0-0=0 & i_3 &= 1-0=1 \\ \mu_1 &= -1-1=-2 & \mu_3 &= 0-0=0 \\ i_2 &= 0+0+0+0-1-0+1+0=0 \\ \mu_2 &= 1+0+1+0-0-0-1+1=2 \end{aligned}$$



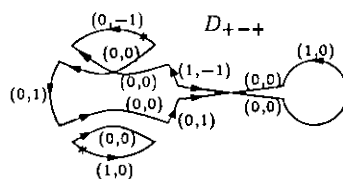
$$\begin{aligned} i_1 &= 0-0=0 & \mu_1 &= -1-1=-2 \\ i_2 &= 0+0-1+0+0+0+1-0+1+0=1 \\ \mu_2 &= 1+1-0+1+1+1+0-1-1+1=4 \end{aligned}$$



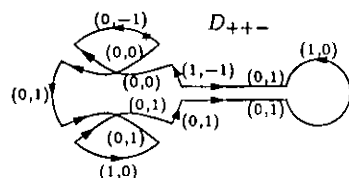
$$\begin{aligned} i_1 &= 0-0=0 & \mu_1 &= -1-1=-2 \\ i_2 &= 0+0-1+0+0+0-1-0+1+0=-1 \\ \mu_2 &= 1+1-0+1+1+0-0-0-1+1=4 \end{aligned}$$



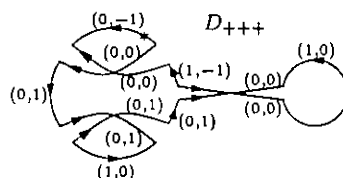
$$\begin{aligned} i_1 &= 0-0-1+0-1-0-0-0-0-0=-2 \\ \mu_1 &= -1-0+1+1-0-1-1-0-1-0=-2 \\ i_2 &= 1-0=1 & \mu_2 &= 0-0=0 \end{aligned}$$



$$\begin{aligned} i_1 &= 0-0-1+0+1-0-0-0-0-0=0 \\ \mu_1 &= -1-0+1+0+0-0-1-0-1-0=-2 \\ i_2 &= 1-0=1 & \mu_2 &= 0-0=0 \end{aligned}$$



$$\begin{aligned} i_1 &= 0-0-1+0-1-0-0-0+1-0-0-0=-1 \\ \mu_1 &= -1-0+1+1-0-1-1-1+0-1-1-0=-4 \end{aligned}$$



$$\begin{aligned} i_1 &= 0-0-1+0+1-0-0-0+1-0-0-0=1 \\ \mu_1 &= -1-0+1+0+0-0-1-1+0-1-1-0=-4 \end{aligned}$$

Figure 13. Eight splittings of the marked chord diagram  $D$  of Fig.11. We assume the chords enumerated as in Fig.11. Calculating the indices of a component we are walking along the component starting from the point  $*$  in the direction of the arc containing  $*$ .

Therefore

$$\begin{aligned}
\langle D \rangle_3 &= -8 \left( \begin{aligned} &-4(-T_0(h))(-T_2(h))T_1(h) \\ &+4(-T_0(h))(-T_0(h))T_1(h) \\ &+2(-T_0(h))T_1(h) - 2(-T_0(h))T_1(h) \\ &+2(-T_2(h))T_1(h) - 2(-T_0(h))T_1(h) \\ &-T_1(h) + T_1(h) \end{aligned} \right) \\
&= -8T_1(h) \left( -4T_0(h)(T_2(h) - T_0(h)) - 2(T_2(h) - T_0(h)) \right) \\
&= 16T_1(h)(T_2(h) - T_0(h))(2T_0(h) - 1) \\
&= 16h(2h^2 - 1 - 1)(2 - 1) = 32h(h^2 - 1),
\end{aligned}$$

since  $T_0(h) = 1$ ,  $T_1(h) = h$ ,  $T_2(h) = 2h^2 - 1$ .

**Proof of Proposition 2.** Let  $C$  be a front with  $n$  dangerous self-tangencies. The proof of Theorem 2 (see sec.3.1) provides an explicit formula for the value  $\langle C \rangle_n$  of the  $n$ th coefficient of the Kauffman bracket on this front:

$$\langle C \rangle_n = 2^n \sum_{\epsilon_1, \dots, \epsilon_n} \epsilon_1 \cdot \dots \cdot \epsilon_n \cdot \langle C_{\epsilon_1, \dots, \epsilon_n} \rangle_0.$$

So for calculation of  $\langle C \rangle_n = \langle D_C \rangle_n$  it is enough to know the zero order coefficients  $\langle C_{\epsilon_1, \dots, \epsilon_n} \rangle_0$ . The lemma below gives an explicit formula for  $\langle C_{\epsilon_1, \dots, \epsilon_n} \rangle_0$  in terms of absolute values of indices and Maslov indices of components of the front  $C_{\epsilon_1, \dots, \epsilon_n}$ . When the ordering of the chords in the diagram  $D_C$  is induced by an ordering of dangerous self-tangencies of the front  $C$ , these absolute values are easily seen to be given by the above algorithm of counting the index information about the splitting  $(D_C)_{\epsilon_1, \dots, \epsilon_n}$ . Thus Proposition 2 follows from

**Lemma 5.** *Let  $C = \cup_{j=1}^l C_j$  be a normal front with  $l$  components. Put  $i_j = \text{ind}(C_j)$  and  $\mu = \sum_{j=1}^l \mu(C_j)$ . Then*

$$\langle C \rangle_0 = -2^{l-1} (-1)^{\mu/2} \prod_{j=1}^l T_{|i_j|}(h),$$

where  $T_n(\cos x) = \cos(nx)$  are the classical Tchebyshev polynomials.

**Proof of Lemma 5.** The second equality of Proposition 1 of sec.2.1 implies that the zero order coefficient of the Kauffman bracket is invariant under dangerous self-tangencies as well. Due to Gromov's theorem (sec.1.2)  $\langle C \rangle_0$  depends only on indices and Maslov indices of components of  $C$ . Therefore it is enough to calculate  $\langle C \rangle_0$  on a collection of canonical curves. Property 4) of Theorem 1 implies

$$\langle K_1 \cdot K_2 \rangle_0 = -2 \langle K_1 \rangle_0 \cdot \langle K_2 \rangle_0 .$$

So Lemma 5 follows from the calculation of the Kauffman bracket on the canonical curves from sec.2.4.

**Remark.** Proposition 2 shows that the orientations of chords in a marked chord diagram do not matter for the value of the symbol of the coefficient. Indeed reorientation of a chord in a diagram  $D$  can effect only the Maslov indices  $\mu_j$  in the formula of the proposition. But for any splitting of the diagram the sum of the  $\mu_j$  modulo 4 is not affected.

In fact the independence from orientations of chords is a general property of the symbol of any  $J^+$ -type invariant  $f$ :

$$\begin{aligned} f(\text{diagram 1}) &= f(\text{diagram 2}) = f(\text{diagram 3}) = f(\text{diagram 4}) - f(\text{diagram 5}) \\ &= f(\text{diagram 6}) - f(\text{diagram 7}) = f(\text{diagram 8}). \end{aligned}$$

The second equality here is due to the fact that we are considering a symbol. The 3rd and 5th ones are the definition. A similar chain of equalities is valid for an inverse dangerous self-tangency.

Thus the orientation of chords in our definition of the marked chord diagram of a front with dangerous self-tangencies should be omitted.

The obtained relation is not the only relation on the values of symbols on our marked chord diagrams. There are a lot of others, some of which are quite obvious. A complete diagrammatic description of symbols of finite order  $J^+$ -type invariants of one component plane fronts has been obtained by J.W.Hill [14]. It turns out that one needs to add one more marking, by the Maslov index of the whole front, to the marked chord diagrams used in [11] in the case of regular plane curve.

### 3.3 The first coefficient

The proposition below means that the first coefficient  $\langle C \rangle_1$  of the Kauffman bracket of a one component normal front  $C$  carries the same information as the quantum Bennequin invariant  $\beta_q(L_C)$  from sec.1.4. Setting  $h = (q + q^{-1})/2$  brings  $\langle C \rangle_1$  to the form whose essential part is  $\beta_q(L_C)$ . A reason for this substitution is that it makes the Tchebyshev polynomials  $T_n(h)$  very simple:  $T_n(\frac{q+q^{-1}}{2}) = \frac{q^n + q^{-n}}{2} = ([n+1]_q - [n]_q)/2$ . Unfortunately the explicit formula relating  $\langle C \rangle_1$  and  $\beta_q(L_C)$  does not look very elegant.

**Proposition 3.** *Let  $C$  be a one component normal front of index  $i$  and Maslov index  $\mu$ . Then*

$$\langle C \rangle_1 |_{h=(q+q^{-1})/2} = (-1)^{\mu/2} \left( \frac{1}{2}(q^i + q^{-i})J^+(C) + 2\beta_q(L_C) \right) + R(i, \mu),$$

where the quantum constant  $R(i, \mu)$  depends only on the index  $i$  and Maslov index  $\mu$  of  $C$ :

$$\begin{aligned} R(i, \mu) = & (-1)^{\mu/2} \left( [|i| + 2]_q - (|i| + |\mu| + \frac{3}{2})[|i| + 1]_q \right. \\ & \left. - (|i| - \frac{|\mu|}{2} - \frac{3}{2})[|i|]_q + (2|i| + \frac{|\mu|}{2} + 2)[|i| - 1]_q + [|i| - 2]_q \right). \end{aligned}$$

**Proof of Proposition 3.** We have to check two points. Firstly, the values of both sides of the identity on the canonical curves should coincide. Secondly, the jumps of both the sides should be the same under a dangerous self-tangency perestroika.

The fact that the canonical curves satisfy the identity follows from the direct computations. We actually introduced the complicated term  $R(i, \mu)$  as the difference between the values of the left-hand side and the remaining part of the right-hand side on the curve  $K_{|i|+1, |\mu|/2}$  with any orientation and coorientation (both sides of the identity do not change when we either reorient or reorient a front). The evaluation of the “main” part of the right-hand side on the canonical curves is provided by the settings and computations of secs. 1.3 and 1.4. The left-hand side of the identity is  $\frac{\partial}{\partial A} \langle C \rangle |_{A=1}$ . Its evaluation on the canonical curves is based on rather elementary calculations (we omit them here) of similar derivatives of the deformations of the Tchebyshev polynomials of sec.2.3.

Now  $R(i, \mu)$  does not change under any perestroika. So the jump of the right-hand side of the identity under a dangerous self-tangency perestroika at a point  $d$  of index  $i_d$  is equal (see secs. 1.3 and 1.4) to

$$(-1)^{\mu/2}(q^i + q^{-i} + 2(q^{i_d} + q^{-i_d})).$$

Let us calculate the jump of the left-hand side. According to sec.3.2 we associate one of the marked chord diagrams of Fig.14 to a dangerous self-tangency point  $d$  of index  $i_d$ .

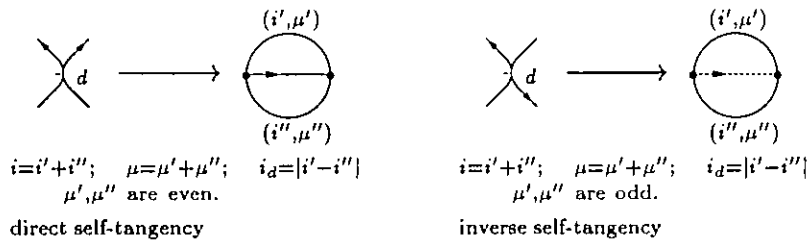


Figure 14. Marked chord diagrams of a self-tangency point  $d$  of index  $i_d$ .

The jump is the value of the symbol of  $\langle C \rangle_1$  on the corresponding diagram. By Proposition 2 (sec.3.2) for calculation of the values we have to consider two splittings of each of the diagrams as it is shown in Fig.15.

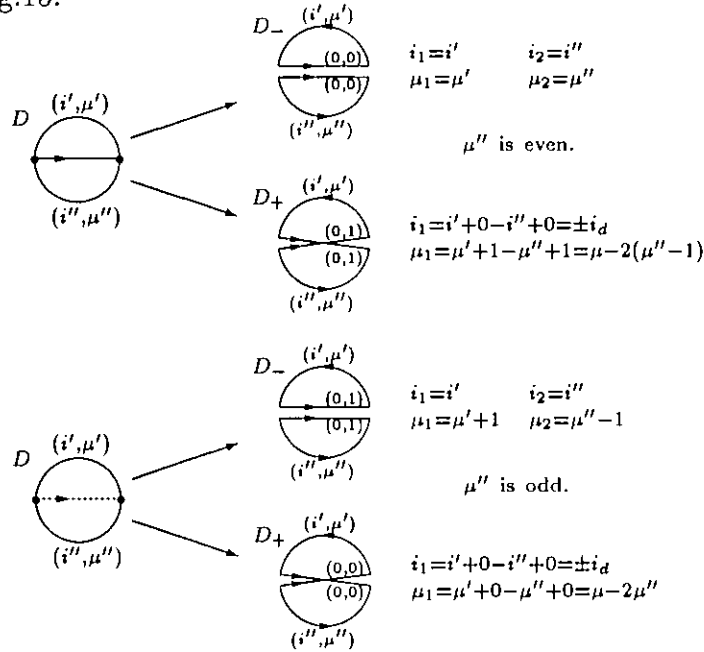


Figure 15. Splittings of marked chord diagrams with one chord.



Therefore the jumps are as follows.

For a direct self-tangency ( $\mu''$  is even):

$$\begin{aligned}
\langle \langle \text{Diagram 1} \rangle \rangle_1 - \langle \langle \text{Diagram 2} \rangle \rangle_1 &= -2 \left( -2(-1)^{\mu'/2} T_{|i'|}(h) (-1)^{\mu''/2} T_{|i''|}(h) \right. \\
&\quad \left. + (-1)^{\mu/2 - (\mu''-1)} T_{i_d}(h) \right) \\
&= 2(-1)^{\mu/2} \left( 2T_{|i'|}(h) T_{|i''|}(h) + T_{i_d}(h) \right) \\
&= 2(-1)^{\mu/2} \left( 2 \cdot \frac{q^i + q^{-i'}}{2} \cdot \frac{q^{i''} + q^{-i''}}{2} + \frac{q^{i_d} + q^{-i_d}}{2} \right) \\
&= (-1)^{\mu/2} (q^i + q^{-i} + 2(q^{i_d} + q^{-i_d})).
\end{aligned}$$

For an inverse self-tangency ( $\mu''$  is odd):

$$\begin{aligned}
\langle \langle \text{Diagram 3} \rangle \rangle_1 - \langle \langle \text{Diagram 4} \rangle \rangle_1 &= -2 \left( -2(-1)^{(\mu'+1)/2} T_{|i'|}(h) (-1)^{(\mu''-1)/2} T_{|i''|}(h) \right. \\
&\quad \left. + (-1)^{\mu/2 - \mu''} T_{i_d}(h) \right) \\
&= 2(-1)^{\mu/2} \left( 2T_{|i'|}(h) T_{|i''|}(h) + T_{i_d}(h) \right) \\
&= (-1)^{\mu/2} (q^i + q^{-i} + 2(q^{i_d} + q^{-i_d})).
\end{aligned}$$

Proposition 3 is proved.

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