UNIQUENESS OF EINSTEIN KÄHLER METRICS

MODULO CONNECTED GROUP ACTIONS

Dedicated to Professor Kunihiko Kodaira on his 70th birthday

by

Shigetoshi Bando*) and Toshiki Mabuchi**)

Max-Planck-Institut für Mathematik Gottfried-Claren-Str. 26 5300 Bonn 3 Federal Republic of Germany

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§ 0 . Introduction

Throughout this paper, we fix an arbitrary n-dimensional compact complex manifold X with positive first Chern class $c_1(X)_{\mathbf{P}} > 0$. We then put

 $K := \text{set of all Kähler forms on } X \text{ representing } 2\pi c_1(X)_{\mathbb{R}},$ $K^+ := \{\omega \in K \mid \omega \text{ has positive definite Ricci tensor}\},$ $E := \{\omega \in K \mid \omega \text{ is Einstein}\},$ $C^{\infty}(X)_{\mathbb{IR}} := \text{space of real valued } C^{\infty} \text{ functions on } X,$ Aut(X) := group of holomorphic automorphisms of X, $G := Aut^0(X) = \text{identity component of } Aut(X).$

Furthermore, Aut(X) is always assumed to act from the <u>right</u> on K by $(\omega, g) \in K \times Aut(X) \rightarrow g^* \omega \in K$.

The main purpose of this paper is to prove the uniqueness of Einstein Kähler metrics, if any, on X up to G-action. Such uniqueness was known only for i) Kähler C-spaces (cf. Matsushima [12]) and ii) some non-homogeneous Einstein manifolds recently discovered by Sakane [13]. Now, the correct statement we obtained has the following stronger form as announced earlier in [9]:

Theorem A. Fix an element ω_1 of K. Let $\mu^+: K^+ \rightarrow \mathbb{R}$ be the restriction to K^+ of the K-energy map $\omega \in K \longmapsto M(\omega_1, \omega) \in \mathbb{R}$ of the Kähler manifold (X, ω_1) (see § 1, also [9]). Assume that $E \neq \phi$. Then

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(i) μ^{\dagger} is bounded from below and takes its absolute minimum exactly on E.

(ii) E consists of a single G-orbit.

We now suppose that $E \neq \phi$, and let K be a maximal compact subgroup of G. By the well-known theorem of Matsushima [11], there exists an element θ of E such that the isotropy subgroup of G at θ coincides with K. Hence E is G-equivariantly diffeomorphic to G/K. Note that G/K has a structure of a Riemannian symmetric space, though the choice of its metric is not unique (even up to constant multiple) if the symmetric space G/K is reducible. We now endow E with the natural Riemannian metric defined in [10] (see also § 9 of the present paper). Then Theorem A allows us to sharpen a result in [10] and one can determine the structure of E as follows:

Theorem B. If $E \neq \phi$, then E is G-equivariantly isometric to the Riemannian symmetric space G/K endowed with a suitable metric, and furthermore, Aut(X) acts isometrically on E.

As a straightforward consequence of Theorem B, we obtain

Theorem C. Let H be an arbitrary (possibly non-connected) compact subgroup of Aut(X). If in addition $E \neq \phi$, then there always exists an H-invariant Einstein Kähler metric on X.

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We now briefly explain how the proof of Theorem A is carried out. Let ω_0 be an arbitrary element of K and $R(\omega_0)$ be the corresponding Ricci form (cf. § 1). (Later in this introduction, we set $\omega_0 := R(\widetilde{\omega})$ for some element $\widetilde{\omega}$ of K^+ and vary ω_0 together with $\widetilde{\omega}$.) Since $R(\omega_0)$ is cohomologous to ω_0 , there uniquely exists a function $f \in C^{\infty}(X)_{\mathbb{R}}$ such that

$$R(\omega_0) = \omega_0 + \sqrt{-1} \partial \overline{\partial} f$$
 and $\int_X \exp(f) \omega_0^n = \int_X \omega_0^n$.

We then consider the following 1-parameter families of equations:

(0.1)
$$\log ((\omega_0 + \sqrt{-1} \ \partial \overline{\partial} \psi_t)^n / \omega_0^n) = -t\psi_t + f, \quad 0 \le t \le 1$$
,
(0.2) $\log ((\omega_0 + \sqrt{-1} \ \partial \overline{\partial} \psi_t)^n / \omega_0^n) = -t\varphi_t - L(0, \varphi_t) + f, \quad 0 \le t \le 1$,

(see § 1 for the definition of L), where in both cases, solutions ψ_+ and ϕ_+ are required to belong to

 $H := \{ \varphi \in C^{\infty}(X)_{\mathbb{R}} \mid \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \varphi \quad \text{is positive definite on } X \}.$

Note that (0.1) above is introduced by Aubin [2] in his study of Einstein Kähler metrics on compact Kähler manifolds with $c_1 > 0$. One can easily pass from the solutions of one of (0.1)and (0.2) to those of the other because for each t, the difference between Ψ_t and φ_t is just a constant (which may depend on t) on X. Now a crucial step of the proof of Theorem A is to show the following fact: (0.3) <u>Given an orbit</u> **O** in E, we can connect **O** with every sufficiently general point $\tilde{\omega}$ of K^+ by a <u>smooth</u> 1-parameter family of solutions $\{\psi_t \mid 0 \leq t \leq 1\}$ (<u>resp.</u> $\{\phi_t \mid 0 \leq t \leq 1\}$) of (0.1) (<u>resp.</u> (0.2)) such that

$$\begin{cases} \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \psi_1 = \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \phi_1 \in \mathbf{O} \ , \ \underline{\text{and}} \\ \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \psi_0 = \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \phi_0 = \widetilde{\omega} \ , \ (\text{i.e., } R(\widetilde{\omega}) = \omega_0) \ . \end{cases}$$

Once one shows (0.3), the proof of Theorem A proceeds as follows:

(i) Consider the K-energy map $\mu: K \to \mathbb{R}$ of the Kähler manifold (X, ω_0) . Recall that μ takes a constant value C on **O** (cf. [9]). Since $\mu(\omega_0 + \sqrt{-1} \ \partial \overline{\partial} \varphi_t)$ is a monotone decreasing function of t (cf. (5.1)), the fact (0.3) above implies $\mu \ge C$ on a dense subset of K^+ . By the continuity of μ , we obtain $\mu \ge C$ on K^+ (cf. (8.1)).

(ii) Note that (0.2) has a unique solution $\varphi_0 \in H$ at t = 0 (cf. (4.3.2)). Hence one can easily show that, over $\{0 \le t \le 1\}$, only one smooth family of solutions of (0.2) is possible (cf. (5.3), (5.4)). We now fix arbitrary G-orbits $\mathbf{O}_1, \mathbf{O}_2$ in E. In view of (0.3), a sufficiently general $\widetilde{\omega} \in K^+$ can be connected with both \mathbf{O}_1 and \mathbf{O}_2 by smooth families $\{\varphi_t^{[i]} \mid 0 \le t \le 1\}$ (i = 1,2) of solutions of (0.2) such that

 $\omega_0 + \sqrt{-1} \partial \overline{\partial} \varphi_1^{[i]} \in \mathbf{O}_i , \ (i = 1, 2), \ \text{where} \ \mathbb{R}(\widetilde{\omega}) = \omega_0 \ .$

Since these two families must coincide, we have $\varphi_1^{[1]} = \varphi_1^{[2]}$ and therefore $\mathbf{0}_1 = \mathbf{0}_2$.

We now give an outline of the proof of (0.3). It roughly consists of the following three steps. (For technical reasons, the actual proof is not divided into such steps.)

Step 1: Given a point $\theta \in \mathbf{O}$, we can always find a solution ψ_1 of (0.1) at t = 1 such that $\theta = \omega_0 + \sqrt{-1} \ \partial \overline{\partial} \psi_1$. This ψ_1 does not necessarily extend to a smooth family $\{\psi_t \mid 1 - \varepsilon \leq t \leq 1\}$ of solutions of (0.1). Because if such a family exists, differentiating (0.1) by t at t = 1, we have $(\Box_{\theta} + 1)(\psi_{t|t=1}) = -\psi_1$. Hence ψ_1 must satisfy

(0.4) $\int_X \psi_1 \varphi \theta^n = 0$ for all $\varphi \in H_{\theta}$,

where H_{θ} denotes Ker $(\sigma_{\theta} + 1)$ in $C^{\infty}(X)_{\mathbb{R}}$. We therefore seek $\theta = \omega_0 + \sqrt{-1} \partial \overline{\partial} \psi_1 \in \mathbf{O}$ which satisfies (0.4), and a method to find such a θ will be given in § 6. However, the condition (0.4) is not enough (cf. (7.2), (7.3)) and a detailed analysis of (0.1) using a bifurcation technique will be effectively employed (cf. § 7). Finally, since the point $\widetilde{\omega} \in K_+$ is sufficiently general, a suitable chosen ψ_1 (resp. ϕ_1) continues to a smooth family { ψ_t | $1 - \varepsilon \leq t \leq 1$ } (resp. { ϕ_t | $1 - \varepsilon \leq t \leq 1$ }) of solutions of (0.1) (resp. (0.2)).

Step 2: By the monotonicity of $\mu(\omega_0 + \sqrt{-1}\partial\partial\phi_t)$ (where we

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always consider such $\phi_{\texttt{t}}\,\texttt{'s}\,$ as depend smoothly on $\,\texttt{t}\,\texttt{)}\,,$ one has

$$(0.5) \qquad \mu(\omega_0 + \sqrt{-1}\partial\overline{\partial}\varphi_t) \ge \mu(\omega_0 + \sqrt{-1}\partial\overline{\partial}\varphi_1) = \mu(\theta)$$

along the solutions of (0.2), and the family $\{\varphi_t \mid 1 - \varepsilon \le t \le 1\}$ in Step 1 uniquely extends to a smooth family $\{\varphi_t \mid 0 < t \le 1\}$ of solutions of (0.2), because for each $0 < t \le 1$, the existence of the lower bound of $\mu(\omega_0 + \sqrt{-1}\partial\overline{\partial}\varphi_t)$ gives us a rough appriori estimate of $\|\varphi_t\|_{C^0}$.

<u>Step</u> 3: Now, another difficulty comes up at t = 0, since the straightforward appriori bound of $\|\varphi_t\|_{C^0}$ obtained from (0.5) tends to infinity as $t \neq 0$. In § 3, we derive a general lower bound of the Green's function of the Laplacian from the isoperimetric inequality of Gallot [6]. This bound allows us to overcome the difficulty and thus we complete the whole extension to $\{\varphi_t \mid 0 \leq t \leq 1\}$.

In conclusion, we wish to thank Professors S. Kobayashi and H. Ozeki and Doctors I. Enoki and R. Kobayshi for their valueable suggestions and constant encouragements. § 1. Notation, convention and preliminaries

(1.1) Throughout this paper (with the only exception of § 8) we fix, once for all, an element ω_0 of K. In addition to the notation defined in Introduction, we put

$$C^{\infty}(X)_{\mathbb{C}}$$
 := space of complex valued C^{∞} functions on X,
 C := space of real d-closed (1,1)-forms on X in
 $2\pi c_1(X)_{\mathbb{R}}$,
 V := set of all volume forms on X,

where on X, everywhere positive real 2n-form is called a volume form. We write an arbitrary element ω of K as

$$\omega = \sqrt{-1} \sum_{\alpha \overline{\beta}} dz^{\alpha} \wedge dz^{\overline{\beta}}$$

in terms of holomorphic local coordinates $z = (z^1, z^2, ..., z^n)$ on X. The corresponding Ricci tensor is denoted by $\sum R(\omega)_{\alpha\overline{\beta}} dz^{\alpha} \bullet dz^{\overline{\beta}}$ and we put $R(\omega) := \sqrt{-1} \sum R(\omega)_{\alpha\overline{\beta}} dz^{\alpha} \wedge dz^{\overline{\beta}}$. Then $R(\omega) = \sqrt{-1}\partial \partial \log det (g_{\alpha\overline{\beta}}) \in C$. We furthermore denote by $\sigma(\omega)$ (resp. \Box_{ω}) the corresponding scalar curvature (resp. Laplacian on functions):

$$\sigma(\omega) := \sum g^{\overline{\beta}\alpha} R(\omega)_{\alpha\overline{\beta}} ,$$
$$\Box_{\omega} := \sum g^{\overline{\beta}\alpha} \partial^2 / \partial z^{\alpha} \partial z^{\overline{\beta}}$$

where $(g^{\overline{\beta}\alpha})$ is the inverse matrix of $(g_{\alpha\overline{\beta}})$. For each $\varphi \in C^{\infty}(X)_{\mathbb{R}}$, we put

(1.1.1) $\omega_0(\varphi) := \omega_0 + \sqrt{-1} \partial \overline{\partial} \varphi$,

(1.1.2)
$$\Omega_{0}(\phi) := \exp(-\phi) \tilde{\omega}^{n}$$
,

where $\tilde{\omega}$ is the unique element of K_{+} such that $R(\tilde{\omega}) = \omega_{0}$. Recall that the following is a straightforward consequence of Yau's affirmative answer [14] to Calabi's conjecture:

(1.1.3) <u>The mapping</u> $\omega \in K \longrightarrow \mathbb{R}(\omega) \in \mathbb{C}$ <u>defines a homeomorphism</u> $\mathbb{R}: (K, \| \|_{C} \mathbb{R} + 2, \alpha) \cong (\mathbb{C}, \| \|_{C} \mathbb{R}, \alpha)$ <u>for each</u> $(\mathbb{K}, \alpha) \in \mathbb{Z} \times \mathbb{R}$ <u>with</u> $\mathbb{R} \ge 0$ <u>and</u> $0 < \alpha < 1$.

Now, (1.1.1) and (1.1.2) above defines the mappings $\omega_0 : \varphi \in C^{\infty}(X)_{\mathbf{IR}} \longrightarrow \omega_0(\varphi) \in C$ and $\Omega_0 : \varphi \in C^{\infty}(X)_{\mathbf{IR}} \longrightarrow \Omega_0(\varphi) \in V$. Let $H := \{\varphi \in C^{\infty}(X)_{\mathbf{IR}} | \omega_0(\varphi) \in K \}$ as in § 0. Then the natural map

$$\begin{array}{c} H \longrightarrow K \\ \varphi \longmapsto \omega_0 \left(\varphi \right) \end{array}$$

is surjective. To each $\omega_0(\varphi) \in K$, the corresponding $\Box_{\omega_0}(\varphi)$, $\sigma(\omega_0(\varphi))$, $R(\omega_0(\varphi))$ will be denoted respectively by \Box_{φ} , $\sigma(\varphi)$, $R(\varphi)$ for simplicity. Finally, we define the mappings $\nu: K \rightarrow V$ and $Ric: V \rightarrow C$ by

 $v(\omega) := \omega^{n} \qquad (\omega \in K) ,$ Ric(Ω) := $\sqrt{-150} \log v \qquad (\Omega \in V) ,$

where we write $\Omega = v(z) \prod_{\alpha=1}^{n} (\sqrt{-1} dz^{\alpha} \wedge dz^{\overline{\alpha}})$ in terms of holomorphic local coordinates $z = (z^1, z^2, \dots, z^n)$ on X. Then the following diagram commutes:



(1.2) Let I be a (not necessarily open or closed) interval in \mathbb{R} , and S be either I or a product $I \times I \times ... \times I$ of I. A family { φ_{g} | $s \in S$ } of functions in $C^{\infty}(X)_{\mathbb{R}}$ is said to be <u>smooth</u> if the mapping

> $S \times X \longrightarrow \mathbb{R}$ $(s,x) \longmapsto \phi_{s}(x)$

is C^{∞} . Any 1-parameter family $\{\varphi_{s} | s \in I\}$ of fuctions in $C^{\infty}(X)_{\mathbb{R}}$ is called a path, and for every smooth path $\{\varphi_{t} | t \in I\}$, the function $\partial \varphi_{t} / \partial t \in C^{\infty}(X)_{\mathbb{R}}$ is denoted by φ_{t} .

(1.3) Let S be a non-empty set. Then a mapping $H: S \times S \longrightarrow \mathbb{R}$ is said to satisfy the 1-cocycle condition if

(i) $H(\sigma_1, \sigma_2) + H(\sigma_2, \sigma_1) = 0$, and

(ii) $H(\sigma_1, \sigma_2) + H(\sigma_2, \sigma_3) + H(\sigma_3, \sigma_1) = 0$

for all $\sigma_1, \sigma_2, \sigma_3 \in S$.

(1.4) (cf. [9]). Let V_0 be the volume $\int_X \omega_0^n / n!$ of the Kähler manifold (X, ω_0) . We put $V := n!V_0$. To each pair $(\phi', \phi'') \in C^{\infty}(X)_{\mathbb{R}} \times C^{\infty}(X)_{\mathbb{R}}$ (resp. $(\phi', \phi'') \in H \times H$), we associate

a real number $L(\phi',\phi'')$ (resp. $M(\phi',\phi'')$) by

(1.4.1)
$$L(\phi',\phi'') := \int_{a}^{b} (\int_{X} \phi_{t} \omega_{0}(\phi_{t})^{n}/V) dt,$$

(1.4.2) (resp. $M(\varphi', \varphi'') := -\int_{a}^{b} \{\int_{X} \dot{\varphi_{t}} (\sigma(\varphi_{t}) - n) \omega_{0} (\varphi_{t})^{n} / V \} dt \}$, where $\{\varphi_{t} \mid a \leq t \leq b\}$ is an arbitrary piecewise smooth path in $C^{\infty}(X)_{\mathbf{R}}$ (resp. H) such that $\varphi_{a} = \varphi'$ and $\varphi_{b} = \varphi''$. Then $L(\varphi', \varphi'')$ (resp. $M(\varphi', \varphi'')$) is independent of the choice of the path $\{\varphi_{t} \mid a \leq t \leq b\}$ and therefore well-defined. Recall that L (resp. M) satisfies the 1-cocycle condition. Furthermore,

(1.4.3)
$$L(\phi_1,\phi_2 + C) = L(\phi_1 - C,\phi_2) = L(\phi_1,\phi_2) + C$$
,

(1.4.4) (resp.
$$M(\phi_1 + C_1, \phi_2 + C_2) = M(\phi_1, \phi_2)$$
),

for all $\varphi_1, \varphi_2 \in C^{\infty}(X)_{\mathbb{R}}$ (resp. $\varphi_1, \varphi_2 \in H$) and all $C \in \mathbb{R}$ (resp. $C_1, C_2 \in \mathbb{R}$). In view of (1.4.4) above, $M: H \times H \longrightarrow \mathbb{R}$ factors through $K \times K$. Hence we can define the mapping $M: K \times K \longrightarrow \mathbb{R}$ (denoted by the same M) by

$$M(\omega', \omega'') := M(\phi', \phi'') \qquad (\omega', \omega'' \in K),$$

where φ', φ'' are elements of H such that $\omega_0(\varphi') = \omega'$ and $\omega_0(\varphi'') = \omega''$. Then the mapping

$$\mu: K \longrightarrow \mathbb{R}$$
$$\omega \longmapsto \mu(\omega) := M(\omega_0, \omega)$$

is called the K-<u>energy map</u> of the Kähler manifold (X, ω_0) . We now put $H_0 := \{ \varphi \in H \mid L(0, \varphi) = 0 \}$. The mapping $\varphi \in H_0 \mapsto \omega_0(\varphi) \in K$ enables us to identify H_0 with K, and we have the following commutative diagram:



(1.5) We regard L as a function on $V \times V$ via the identification $\Omega_0: C^{\infty}(X)_{\mathbb{R}} \cong V$. Let $N: K \times K \longrightarrow \mathbb{R}$ be the pullback $-v \star L$ of -L by $v: K \longrightarrow V$. Then this N is characterized by the following commutative diagram:

$$\begin{array}{cccc} K \times K & \longrightarrow & \mathbf{R} \\ \downarrow & \vee \times \vee & & \uparrow^{-\mathbf{L}} \\ V \times V & \stackrel{\sim}{=} & C^{\infty}(\mathbf{X})_{\mathbf{R}} \times C^{\infty}(\mathbf{X})_{\mathbf{R}}. \end{array}$$

Since L satisfies the 1-cocycle condition, so does N. A straighforward computation shows that, for each pair $(\omega', \omega'') \in K \times K$, the number $N(\omega', \omega'')$ is given by

(1.5.1)
$$N(\omega', \omega'') = \int_{a}^{b} \{ \int_{X} (\Box_{\varphi_{t}} \dot{\varphi_{t}}) R(\varphi_{t})^{n} / V \} dt.$$

where $\{\varphi_t \mid a \leq t \leq b\}$ is an arbitrary piecewise smooth path in *H* such that $\omega_0(\varphi_a) = \omega'$ and $\omega_0(\varphi_b) = \omega''$.

<u>Remark</u> (1.5.2). Several generalizations of L,M,N (which is announced to appear in this paper, cf. [8],[9]) will be given separately in [3] as a self-contained article.

(1.6) (cf. Aubin [2]). For each pair $(\phi', \phi'') \in H \times H$, we put (1.6.1) $I(\phi', \phi'') := \int_X (\phi'' - \phi') (\omega_0 (\phi')^n - \omega_0 (\phi'')^n) / V$, (1.6.2) $J(\phi', \phi'') := -L(\phi', \phi'') + \int_X (\phi'' - \phi') \omega_0 (\phi')^n / V$.

Then both $I: H \times H \longrightarrow \mathbb{R}$ and $J: H \times H \longrightarrow \mathbb{R}$ factors through $K \times K$, i.e., I and J are regarded as functions on $K \times K$ by

$$I(\omega',\omega'') := I(\phi',\phi'')$$
 and $J(\omega',\omega'') := J(\phi',\phi'')$

for all $\omega', \omega'' \in K$, where $\phi', \phi'' \in H$ are such that $\omega_0(\phi') = \omega'$ and $\omega_0(\phi'') = \omega''$. We later need the following properties of I and J:

(1.6.3) $J(\phi',\phi'') \ge 0$ and the equality holds if and only if $\phi' = \phi'' + \text{constant}$.

(1.6.4) $0 \le I(\phi', \phi'') \le (n + 1) (I(\phi', \phi'') - J(\phi', \phi'')) \le n I(\phi', \phi'')$ for all ϕ', ϕ'' .

These follow from Aubin's result [2; p. 146] and the identity $J(\phi',\phi'') + J(\phi'',\phi'') = I(\phi'',\phi'') = I(\phi'',\phi'')$. We now take an arbitrary smooth path $\{\phi_t \mid a \leq t \leq b\}$ in H. Then a simple calculation shows that

$$(1.6.5) \quad \frac{d}{dt} (I(0,\varphi_{t}) - J(0,\varphi_{t})) = -\int_{X} \varphi_{t} (\Box_{\varphi_{t}} \dot{\varphi}_{t}) \omega_{0} (\varphi_{t})^{n} / V.$$

(1.7) Throughout this paper, we always denote by f the function in $C^{\infty}(X)_{TR}$ defined by

(1.7.1)
$$R(\omega_0) = \omega_0 + \sqrt{-1}\partial\overline{\partial}f$$
 and $\int_X \exp(f)\omega_0^n = V$ (cf. § 0).
To each $\varphi \in H_0$, we can similarly associate a function
 $f_{\varphi} \in C^{\infty}(X)_{\mathbb{R}}$ with the following properties (cf. [9]):
(1.7.2) $R(\varphi) = \omega_0(\varphi) + \sqrt{-1}\partial\overline{\partial}f_{\varphi}$,
(1.7.3) $\mu(\omega_0(\varphi)) = -\int_X f_{\varphi} \omega_0(\varphi)^n / V$,
(1.7.4) $\frac{\partial}{\partial t} f_{\varphi_t} = -(\Box_{\varphi_t} + 1) \dot{\varphi_t}$ for every smooth path
 $\{\varphi_t \mid a \le t \le b\}$ in H_0 .
(1.8) We shall now show that
(1.8.1) $N(\omega', \omega'') - M(\omega', \omega'') = J(\omega'', R(\omega'')) - J(\omega', R(\omega'))$ for all
 $\omega', \omega'' \in K$.

<u>Proof</u>: Choose $\varphi', \varphi'' \in H_0$ so that $\omega_0(\varphi') = \omega'$ and $\omega_0(\varphi'') = \omega''$. Let $\varphi_t := \varphi' + t(\varphi'' - \varphi') + C_t \in H_0$, $0 \le t \le 1$, and we denote each $\omega_0(\varphi_t), f_{\varphi_t}, \Box_{\varphi_t}$ respectively by $\omega^{(t)}, f_t, \Box_t$. Then by (1.7.2) ~ (1.7.4) and (1.5.1),

$$\frac{d}{dt} M(\omega_0, \omega^{(t)}) = \frac{d}{dt} \mu(\omega^{(t)}) = \frac{d}{dt} (-\int_X f_t(\omega^{(t)})^n / V)$$

$$\frac{d}{dt} N(\omega_0, \omega^{(t)}) = \int_X (-\psi_t \dot{\phi}_t) R(\phi_t)^n / V$$

$$= -\int_X (\dot{\phi}_t + \dot{f}_t) \omega_0 (\phi_t + f_t)^n / V = -\frac{d}{dt} L(0, \phi_t + f_t).$$

Since $L(0, \varphi_t) = 0$ for all t, we have

$$\frac{d}{dt} \{ N(\omega_0, \omega^{(t)}) - M(\omega_0, \omega^{(t)}) \} = \frac{d}{dt} \{ \int_X f_t(\omega^{(t)})^n / V - L(\varphi_t, \varphi_t + f_t) \} = \frac{d}{dt} J(\omega^{(t)}, R(\varphi_t)).$$

Integrating this over the interval [0,1], we obtain (1.8.1).

§ 2. Matsushima's theorem and some identities on Einstein Kähler manifolds.

Throughout this section, we assume $E \neq \phi$, and then fix an arbitrary element $\theta = \sqrt{-1} \int \theta_{\alpha \overline{\beta}} dz^{\alpha} \wedge dz^{\overline{\beta}}$ of E. By quoting the well-known theorem of Matsushima [11], we shall introduce several notations on Einstein Kähler manifolds. Some technical identities on such manifolds will also be proven for later purposes.

(2.1) Let \mathfrak{g} be the space $\operatorname{H}^{0}(X, \mathcal{O}(\operatorname{T}(X)))$ of all holomorphic vector fields on X. For each $Y \in \mathfrak{g}$, let $Y_{\mathbb{IR}}$ denote the real vector field $Y + \overline{Y}$ and we set $\mathfrak{g}_{\operatorname{real}} := \{Y_{\mathbb{IR}} \mid Y \in \mathfrak{g}\}$. Then $Y \mapsto Y_{\mathbb{IR}}$ defines the isomorphism of the complex Lie algebras $(\mathfrak{g}, \sqrt{-1}) \cong (\mathfrak{g}_{\operatorname{real}}, J)$, where J is the complex structure of X. We now consider the G-orbit **O** through θ in \overline{E} . This is written as

in terms of the isotropy subgroup K_{θ} of G at Θ . Let k_{θ} be the set of all Killing vector fields on X with respect to the Kähler metric Θ , where each Killing vector field is regarded as an element of \mathfrak{g} via the identification $\mathfrak{g} \cong \mathfrak{g}_{real}$. Then k_{θ} is the Lie subalgebra of \mathfrak{g} corresponding to K_{θ} in G. For each $\varphi \in C^{\infty}(X)_{\mathfrak{C}}$, we define the vector field Y^{φ}_{θ} on X by

 $(2.1.1) \qquad \Upsilon^{\varphi}_{\theta} := \frac{1}{2} \sum \varphi^{\alpha} \partial / \partial z^{\alpha},$

where $\varphi^{\alpha} = \sum \theta^{\overline{\beta}\alpha} \partial_{\overline{\beta}} \varphi$, $(\theta^{\overline{\beta}\alpha})$ being the inverse matrix of $(\theta_{\alpha\overline{\beta}})$. Take the 1-parameter group $y^{\varphi}_{\theta t} := \exp(t Y^{\varphi}_{\theta \mathbf{R}})$, $t \in \mathbf{R}$, on X generated by $Y^{\varphi}_{\theta \mathbf{R}}$. We now have the following theorem of Matsushima [11]:

This theorem in particular implies the following identification:

$$(2.2.3) \quad \mathbb{T}_{\theta}(\mathbf{O}) \cong \mathbb{T}_{\mathbf{e}}(\mathbf{G}/\mathbf{K}_{\theta}) = \mathfrak{P}_{\theta} \cong \mathbb{H}_{\theta}$$
$$\frac{\partial}{\partial t}((\mathbf{Y}_{\theta t}^{\boldsymbol{\phi}}) \star \theta)|_{t=0} \iff \mathbf{Y}_{\theta}^{\boldsymbol{\phi}} \iff \boldsymbol{\phi}.$$

We put $\tilde{Y} := Y^{\phi}_{\theta}$ for brevity. Then by the next computation, the left-hand side of (2.2.3) is shown to have a very nice description:

 $(2.2.3) \quad \frac{\partial}{\partial t} \left(\left(Y_{\theta t}^{\varphi} \right)^{*\theta} \right) \Big|_{t=0} = L_{\widetilde{Y}}^{\varphi} = (d \circ i_{\widetilde{Y}} + i_{\widetilde{Y}}^{\varphi} \circ d)_{\theta}$ $= \frac{\sqrt{-1}}{2} d(-\partial \varphi + \overline{\partial} \varphi) = \sqrt{-1} \partial \overline{\partial} \varphi .$

We shall now prove the following technical Lemma:

Lemma (2.3). Let $\langle , \rangle_{\theta} : \{p-forms on X\} \times \{p-forms on X\} \rightarrow C^{\infty}(X)_{\mathbb{C}}, p = 1, 2, \dots, be the natural hermitian pairings induced from the$

Kähler metric
$$\theta$$
. Then for all $\varphi, \psi \in H_{\theta}$ and all
 $\zeta \in C^{\infty}(X)_{\mathbf{R}}$,
(2.3.1) $\Box_{0} < \partial \zeta, \partial \varphi >_{0} = < \partial \overline{\partial} \zeta, \partial \overline{\partial} \varphi >_{0} + < \partial (\Box_{0} \zeta), \partial \varphi >_{0}$, and

 $(2.3.1) \quad \Box_{\theta} < \partial \zeta, \partial \phi >_{\theta} = < \partial \partial \zeta, \partial \phi >_{\theta} + < \partial (\Box_{\theta} \zeta), \partial \phi >_{\theta} , and$ $\underline{\text{in particular}} \quad (\Box_{\theta} + 1) < \partial \psi, \partial \phi >_{\theta} = < \partial \overline{\partial} \psi, \partial \overline{\partial} \phi >_{\theta} = (\Box_{\theta} + 1) < \partial \phi, \partial \psi >_{\theta};$

$$(2.3.2) \quad -\int_X \varphi < \partial \overline{\partial} \zeta, \partial \overline{\partial} \psi >_{\theta} \theta^n = \int_X (\varphi \psi - \langle \partial \varphi, \partial \psi >_{\theta}) \{ (\square_{\theta} + 1) \zeta \} \theta^n .$$

<u>Proof:</u> Fix an arbitrary point x of X, and we choose a holomorphic local coordinates $(z^1, z^2, ..., z^n)$ centered at x such that $\theta_{\alpha\beta}(x) = \delta_{\alpha\beta}$ and $(d\theta_{\alpha\overline{\beta}})(x) = 0$ for all α and β . Note that φ^{α} , $\alpha = 1, 2, ..., n$ are all holomorphic (cf. (2.1.1), (2.2.2)). Therefore, at the point x,

$$\begin{aligned} \Box_{\theta} < \partial \zeta, \partial \phi >_{\theta} &= \Box_{\theta} \left(\sum \zeta_{\alpha} \phi^{\alpha} \right) = \sum_{\alpha, \beta} \zeta_{\alpha} \overline{\beta} \phi^{\alpha} \beta + \sum_{\alpha, \beta} \zeta_{\alpha} \beta \overline{\beta} \phi^{\alpha} \\ &= \langle \partial \overline{\partial} \zeta, \partial \overline{\partial} \phi >_{\theta} + \langle \partial (\Box_{\theta} \zeta), \partial \phi >_{\theta} \end{aligned}$$

which proves (2.3.1). For (2.3.2), let ξ := $(\texttt{p}_{\theta}+1)\zeta$. Then

$$\begin{split} &\int_{X} (\varphi \psi - \langle \partial \varphi, \partial \psi \rangle_{\theta}) \xi \theta^{n} \\ &= -\sqrt{-1} \int_{X} (\varphi \partial \overline{\partial} \psi + \partial \varphi \wedge \overline{\partial} \psi) \xi \wedge n \theta^{n-1} \qquad (\text{because } \psi = -\pi_{\theta} \psi) \\ &= -\sqrt{-1} \int_{X} \xi \partial (\varphi \overline{\partial} \psi) \wedge n \theta^{n-1} = \sqrt{-1} \int_{X} \varphi \partial \xi \wedge \overline{\partial} \psi \wedge n \theta^{n-1} = \int_{X} \varphi \langle \partial \xi, \partial \psi \rangle_{\theta} \theta^{n} \\ &= \int_{X} \varphi \{\pi_{\theta} \langle \partial \zeta, \partial \psi \rangle_{\theta} - \langle \partial \overline{\partial} \zeta, \partial \overline{\partial} \psi \rangle_{\theta} \} \theta^{n} + \int_{X} \varphi \langle \partial \zeta, \partial \psi \rangle_{\theta} \theta^{n} \quad (\text{cf.}(2.3.1)) \\ &= -\int_{X} \varphi \langle \partial \overline{\partial} \zeta, \partial \overline{\partial} \psi \rangle_{\theta} \theta^{n} \qquad (\text{because } \varphi = \pi_{\theta} \varphi). \end{split}$$

§ 3 Lower bounds of Green functions of Laplacians.

In this section, using the isoperimetric inequality of Gallot [6], we shall construct some lower bound of the Green's function of the Laplacian on a compact Riemannian manifold. This bound applies to our compact Kähler situation and allows us to obtain an interesting inequality which becomes crucial in out later investigation.

(3.1) Let (M,g) be an m-dimensional compact Riemannian manifold. The corresponding Ricci tensor, volume, volume form and diameter are denoted respectively by r_g, V_g, dM_g and D_g . We then set

$$a_g := D_g^2 Inf \{r_g(r,r)/(n-1); ||r||_g = 1\}$$
,

where the infinum is taken over all unit targent vectors in T(M). Let Δ_g be the Laplacian of (M,g) (we choose Δ_g so that it always has nonpositive eigen values), and $G_g \in C^{\infty}(M \times M - (diagonal))$ be the corresponding Green function (with the well-known prescribed singularity along the diagonal) characterized by the following properties:

(i)
$$\varphi(\mathbf{x}) = V_{g}^{-1} \int_{M} \varphi(\mathbf{y}) dM_{g}(\mathbf{y}) + \int_{M} G_{g}(\mathbf{x}, \mathbf{y}) (-\Delta_{g} \varphi) (\mathbf{y}) dM_{g}(\mathbf{y}) ,$$

(ii) $\int_{M} G_{g}(\mathbf{x}, \mathbf{y}) dM_{g}(\mathbf{y}) = 0 ,$

for all $x \in M$ and $\varphi \in C^{\infty}(X)_{\mathbf{R}}$.

<u>Theorem</u> (3.2). Let $(m, \alpha) \in \mathbf{Z} \times \mathbf{R}$ be an arbitrary pair <u>satisfying</u> $m \ge 2$ and $\alpha \ge 0$. Then there exists a positive <u>constant</u> $\gamma = \gamma(m, \alpha)$ <u>depending only on</u> m and α <u>such</u> <u>that</u>, for every m-dimensional compact connected Riemannian <u>manifold</u> (M,g) with $a_g \ge -\alpha^2$,

$$G_{g}(x, y) \ge -\gamma(m, \alpha) D_{g}^{2}/V_{g}$$

for all $x, y \in M$ with $x \neq y$.

<u>Remark</u> (3.3). If $\alpha = 0$, the number $\gamma(m, \alpha)$ is easily computed. For instance, $\gamma(2, 0) = 24$.

(3.4) <u>Proof of</u> (3.2): Let W_0 be the space of the functions f in $L_1^2(X)$ which satisfy $\int_M f dM_g = 0$. Then a combination of Theorems of Gallot [6;(1,3),(2,7)] shows that, there exists a positive constant $\kappa(m,\alpha)$ depending only on m and α such that, for every $f \in W_0$, the number $C := \kappa(m,\alpha) V_g^{1/m} D_g^{-1}$ satisfies

Let H(x,y,t) be the heat kernel of (M,g), and we set

$$H_0(x,y,t) := H(x,y,t) - V_g^{-1}$$

The proof is now divided into two cases:

(Case 1) m ≥ 3. A result of Cheng and Li [5;(2.9)] says that (3.4.1) implies

$$0 < H_0(x,x,t) \le 4(2t c^2/m)^{-m/2}$$

where $(x,t) \in M \times \mathbb{R}$ with t > 0. Hence for all x and y,

$$|H_0(x,y,t)| \le H_0(x,x,t)^{1/2} H_0(y,y,t)^{1/2} \le 4(2tC^2/m)^{-m/2}.$$

,

Together with $H_0(x,y,t) \ge -V_g^{-1}$, we obtain

$$G_{g}(x,y) = \int_{0}^{\infty} H_{0}(x,y,t) dt \ge -\int_{0}^{\tau} V_{g}^{-1} dt - \int_{\tau}^{\infty} 4 (2tC^{2}/m)^{-m/2} dt$$

for each $\tau > 0$. If we set $\tau := 2^{4/m} m \kappa(m,\alpha)^{-2} D_g^{-2/2}$, the right-hand side of this inequality is written as $-\gamma(m,\alpha) D_g^{-2/2} V_g$ for some constant $\gamma(m,\alpha)$ depending only on m and α , as required.

(Case 2)^{*)} m = 2. For each
$$(x,t) \in M \times \mathbb{R}$$
, we put
 $M_{x,t} := \{x\} \times M \times \{t\}$, which is a submanifold ($\cong M$) of $M \times M \times \mathbb{R}$.
For C^{∞} functions $\varphi(x,y,t)$ defined on an open subset of
 $M \times M \times \mathbb{R}$, we denote by $d_y : \varphi(x,y,t) \mapsto d_y \varphi(x,y,t)$
the d-operator coming only from the second factor. Then the
same argument as in Cheng and Li [5;(2,7)] together with
(3.4.2) yields

$$(3.4.3) \quad (\partial H_0/\partial t) (x, x, t) = - \| d_y H_0(x, y, t/2) \|_{L^2}^2 (M_{x, t/2, g})$$

$$\leq -C^2 v_g^{-1/2} \| H_0(x, y, t/2) \|_{L^4}^2 (M_{x, t/2, g})$$

for each $(x,t) \in M \times IR$ with t > 0. On the other hand, in view of

$$\int_{M} |H_{0}(x,y,t)| dM_{g}(y) \leq \int_{M} (H(x,y,t) + V_{g}^{-1}) dM_{g}(y) = 2$$

*) G_{g} (=G_g,M) is written in terms of the Green's function $G_{g\times g,M\times M}$, which provides us with a very simple proof of this case by reduction to m = 4. However, the estimate thus obtained is not so sharp (for instance, $\gamma(2,0)$ would exceed 24). and the Hölder inequality

$$\| \| \|_{0}(x,y,t/2) \|_{L^{3/2}(M_{x,t/2,g})}^{2/3} \| \| \|_{0}(x,y,t/2) \|_{L^{3}(M_{x,t/2,g})}^{4/3} \| \|_{L^{3}(M_{x,t/2,g})}^{3}$$

$$\geq \int_{M} \| \|_{0}(x,y,t/2) \|_{2}^{2} dM_{g}(y) (= \| \|_{0}(x,x,t)),$$

we obtain $||H_0(x,y,t/2)||^2_{L^4(M_{x,t/2,g})} \ge H_0(x,x,t)^{3/2}/2$.

This combinded with (3.4.3) shows that

$$(\partial H_0/\partial t)(x,x,t) \leq -(C^2/2) V_g^{-1/2} H_0(x,x,t)^{3/2}.$$

Then the same argument as in Cheng and Li [5;(2.9)] again applies. Thus,

$$H_0(x,x,t) \leq (tC^2/4)^{-2}V_g.$$

Finally, similar to Case 1 above, it follows that

$$G_{g}(x,y) = \int_{0}^{\infty} H_{0}(x,y,t) dt \ge -\int_{0}^{\tau} V_{g}^{-1} dt -\int_{\tau}^{\infty} (tC^{2}/4)^{-2} V_{g} dt$$
$$\ge -8/C^{2} = -8 \kappa (2,\alpha)^{-2} D_{g}^{2}/V_{g}$$

by setting $\tau = 4v_g/C^2$.

(3.5) We now return to our original compact Kähler situation. In terms of the notation in (3.2) above, let $\beta(n) := \gamma(2n,0)$, which is a constant depending on n alone. Furthermore, for each $\varphi \in H$, let Δ_{φ} (resp. Δ_{0}) denote the real Laplacian $2 \sigma_{\varphi}$ (resp. $2 \sigma_{\omega_{0}}$) of the compact Kähler manifold $(X, \omega_{0}(\varphi))$ (resp. (X, ω_{0})), and G_{ω} (resp. G_{0}) be its corresponding Green's function as is defined in (3.1). We denote by $-K_{\varphi}$ (resp. $-K_0$) the infimum of G_{φ} (resp. G_0) on $X \times X$ - (diagonal). Finally, for each $t \in \mathbf{R}$, we put

$$H^{(t)} := \{ \varphi \in H \mid R(\varphi) - t\omega_0(\varphi) \text{ is positive semi-definite} \}.$$

Theorem (3.2) now has the following important implication.

<u>Proposition</u> (3.6). Let t > 0 be arbitrary. Then for every $\varphi \in H^{(t)}$, its oscillation $Osc \varphi := Max \varphi - Min \varphi$ satisfies $X \qquad X$ $Osc \varphi \leq I(0,\varphi) + 2n (K_0V_0 + (n-1)\beta(n)\pi^2t^{-1})$

<u>Proof:</u> We observe, by virtue of the identity $\omega_0(\varphi) = \omega_0 + \sqrt{-1}\partial \overline{\partial}\varphi$, that the following inequalities hold:

$$-\Delta_0 \varphi \leq 2n$$
 and $-\Delta_{\varphi} \varphi \geq -2n$.

Hence we have

$$\varphi(\mathbf{x}) = V_0^{-1} \int_X \varphi \omega_0^n n! + \int_X (G_0(\mathbf{x}, \mathbf{y}) + K_0) (-\Delta_0 \varphi) (\mathbf{y}) \omega_0^n (\mathbf{y}) / n!$$

$$\leq V_0^{-1} \int_X \varphi \omega_0^n / n! + 2nK_0 V_0$$

and

$$\begin{split} \varphi(\mathbf{x}) &= V_0^{-1} \int_X \varphi \omega_0 (\varphi)^n / n! + \int_X (G_{\varphi}(\mathbf{x}, \mathbf{y}) + K_{\varphi}) (-\Delta_{\varphi} \varphi) (\mathbf{y}) \omega_0 (\varphi)^n (\mathbf{y}) / n! \\ &\geq V_0^{-1} \int_X \varphi \omega_0 (\varphi)^n / n! - 2n K_{\varphi} V_0. \end{split}$$

Therefore

$$\begin{aligned} & Osc \, \phi \leq V_0^{-1} \int_X \phi (\omega_0^n - \omega_0 (\phi)^n) / n! + 2n \ (K_0 V_0 + K_{\phi} V_0) \\ & = I(0, \phi) + 2n (k_0 V_0 + K_{\phi} V_0) \ (cf. \ (1.6.1)). \end{aligned}$$

Let D_{φ} be the diameter of $(X, \omega_0(\varphi))$. Since $R(\varphi) \ge t\omega_0(\varphi)$, Meyer's theorem asserts that $D_{\varphi} \le \pi((n-1)/t)^{1/2}$. We now conclude from Theorem (3.2) that

$$K_{\phi}V_{0} \leq \beta(n) D_{\phi}^{2} \leq (n-1) \beta(n) \pi^{2} t^{-1}$$
.

§ 4. Generalized Aubin's equations.

In a recent paper [2], Aubin introduced a very interesting 1-parameter family of non-linear equations to apply the continuity method to showing the existence of Einstein Kähler metrics on some compact Kähler manifolds with $c_1 > 0$. In this section, we shall consider a slightly modified family of equations so that it fits our purpose. Elementary properties of such a family will also be given.

(4.1) We define the mapping
$$A: H \to C^{\infty}(x)_{\mathbb{R}}$$
 by

 $A(\phi) := \log (\omega_0(\phi)^n / \omega_0^n) \quad (\phi \in H) ,$

and then consider the following 1-parameter families of equations:

(4.1.1)
$$A(\phi_t) = -t\phi_t - L(0,\phi_t) + f; 0 \le t \le 1,$$

(4.1.2) $A(\psi_t) = -t\psi_t + f; 0 \le t \le 1,$

where solutions φ_t and ψ_t are both required to belong to *H*. We call the former the family of <u>generalized Aubin's</u> <u>equations</u> of the Kähler manifold (X, ω_0) , while the latter is the original family introduced by Aubin.

<u>Remark</u> (4.1.3). At each point t of (0,1] (resp. [0,1]), we put

$$j'(\phi_{t}) := \phi_{t} + t^{-1}L(0,\phi_{t}) \quad (\text{resp. } j''(\psi_{t}) := \psi_{t} - (t+1)^{-1}L(0,\psi_{t}))$$

for every φ_t (resp. ψ_t) satisfying (4.1.1) (resp. (4.1.2)). Then by substituting $j'(\varphi_t)$ (resp. $j''(\psi_t)$) for ψ_t (resp. φ_t), one can easily see that $j'(\varphi_t)$ (resp. $j''(\psi_t)$) satisfies (4.1.2) (resp. (4.1.1)). Furthermore, $j' \circ j'' = id$ and $j'' \circ j' = id$. Hence in finding solutions for $t \neq 0$, there is no difference between (4.1.1) and (4.1.2).

(4.2) Choose an arbitrary $t \in [0,1]$. Let φ_t (resp. ψ_t) be a solution of (4.1.1) (resp. (4.1.2)), and we set $\omega^{(t)} := \omega_0(\varphi_t)$ (resp. $\omega^{(t)} := \omega_0(\psi_t)$). Then $\omega^{(t)}$ satisfies

(4.2.1) $R(\omega^{(t)}) = (1 - t) \omega_0 + t\omega^{(t)}$,

(cf. Aubin [2; p. 149]), and in particular φ_t (resp. ψ_t) belongs to $H^{(t)}$ (cf. (3.5)). On the other hand, one can easily pass from the solutions of (4.2.1) to those of (4.1.1), though we don't get into details.

(4.3) We'shall next study the solutions of (4.1.1) and (4.1.2) at t = 0. Recall the following affirmative answer to Calabi's conjecture:

<u>Theorem</u> (4.3.1) (Yau [14]). If t = 0, then (4.1.2) has a solution which is unique up to an additive constant.

This in particular implies

<u>Corollary</u> (4.3.2). For t = 0, the equation (4.1.1) has a <u>unique solution</u> φ_0 . Moreover, $L(0,\varphi_0) = 0$ and $R(\omega_0(\varphi_0)) = \omega_0$. <u>Proof</u>: The existence of a solution is straightforward from (4.1.3) and (4.3.1). For uniqueness, let φ_0 be a solution of (4.1.1) at t = 0. Then

$$\int_{\mathbf{X}} \omega_0^n = \int_{\mathbf{X}} \exp\left(\mathbf{A}(\varphi_0)\right) \omega_0^{n} = \exp\left(-\mathbf{L}(0,\varphi_0)\right) \int_{\mathbf{X}} \exp\left(\mathbf{f}\right) \omega_0^n.$$

Since $\int_X \exp(f) \omega_0^n = \int_X \omega_0^n$ (cf. (1.7.1)), we obtain $L(0, \varphi_0) = 0$. Therefore, φ_0 is a solution of (4.1.2) at t = 0. The required uniqueness now follows from a combination of Theorem (4.3.1) and $L(0, \varphi_0) = 0$. $R(\omega_0(\varphi_0)) = \omega_0$ is an immediate consequence of (1.7.1) and $A(\varphi_0) = f$.

<u>Remark</u> (4.3.3). Suppose that $\{\varphi_{t+} 0 \le t \le \tau\}$ ($\tau \ge 0$) is a smooth 1-parameter family of solutions of (4.1.1). By (4.3.2) above, $L(0,\varphi_0) = 0$, and hence by setting $\varphi_t := j'(\varphi_t) = \varphi_t + t^{-1}L(0,\varphi_t)$ (cf. (4.1.3)), we see that $\{\psi_t \mid 0 \le t \le \tau\}$ forms a smooth family of solutions of (4.1.2). (For similar arguments, see Aubin [2; p. 149].)

(4.4) Let $\mathcal{H}^{k,\alpha}$ (where $2 \ge k \in \mathbb{Z}$ and $0 < \alpha < 1$) be the set of all $\varphi \in C^{k,\alpha}(X)_{\mathbb{R}}$ with positive definite $\omega_0(\varphi)$. Note that $\mathcal{H}^{k,\alpha}$ is an open subset of $C^{k,\alpha}(X)_{\mathbb{R}}$. We now conclude this § 4 by showing the following local extension property of solutions of (4.1.1) for $0 \le t < 1$.

<u>Proposition</u> (4.4.1) (cf. Aubin [2]). Let $2 \le k \in \mathbb{Z}$ and fix $\alpha \in \mathbb{R}$ with $0 < \alpha < 1$. Let $0 \le \tau < 1$. Suppose moreover that (4.1.1) has a solution φ_{τ} at $t = \tau$. Then for some $\varepsilon > 0$, φ_{τ} uniquely extends to a smooth 1-parameter family

$$\{\varphi_{\perp} \mid t \in [0,1) \cap [\tau - \varepsilon, \tau + \varepsilon]\}$$

of solutions of (4.1.1) in H, and furthermore, if $(\varphi,t) \in H^{k,\alpha} \times [0,1)$ satisfies the conditions $\|\varphi - \varphi_{\tau}\|_{C^{k,\alpha}} \leq \varepsilon$, $|t - \tau| \leq \varepsilon$ and $A(\varphi) = -t\varphi - L(0,\varphi) + f$, then φ coincides with φ_{t} .

<u>Proof:</u> Consider the mapping $\Gamma: \mathbb{H}^{k, \alpha} \times \mathbb{R} \longrightarrow \mathbb{C}^{k-2, \alpha}(X)$ defined by

$$\Gamma(\varphi,t) := A(\varphi) + t\varphi + L(0,\varphi) - f, (\varphi,t) \in H^{k,\alpha} \times \mathbb{R}$$
.

Then its Frechet derivative $D_{\varphi}\Gamma$: $C^{k,\alpha}(X) \rightarrow C^{k-2,\alpha}(X)$ (at (φ ,t)) with respect to the first factor is given by

Note that, by the well-known regularity theorem, we have $\varphi \in H$ for every $(\varphi, t) \in H^{k, \alpha} \times \mathbb{R}$, whenever $\Gamma(\varphi, t) = 0$. Since $\Gamma(\varphi_{\tau}, \tau) = 0$, an application of the implicit function theorem now reduces the proof to showing that $D_{\varphi}\Gamma$ is invertible at (φ_{τ}, τ) . The following cases are possible. (Case 1) $\tau = 0$. Then $D_{\varphi}\Gamma \mid_{(\varphi_{\Omega}, 0)}$ is the mapping

$$\psi \in \mathsf{C}^{\mathbf{k},\alpha}(\mathbf{X})_{\mathbb{R}} \longmapsto \Box_{\varphi_0} \psi + \int_{\mathbf{X}} \psi \omega_0(\varphi_0)^n / \mathbf{v} \in \mathsf{C}^{\mathbf{k}-2,\alpha}(\mathbf{X})_{\mathbb{R}}$$

which is invertible.

(Case 2) $\tau \neq 0$. Since $R(\phi_{\tau}) - \tau \omega_{0}(\phi_{\tau})$ is positive definite (cf. (4.2)), a theorem of Lichnerowicz [7] asserts that τ is less than the first (positive) eigen value of $-\sigma_{\phi_{-}}$. Hence ${}^{D} \phi^{\Gamma} | (\phi_{\tau}^{}, \tau) \quad \text{ is invertible.}$

<u>Remark</u> (4.4.2). Proposition (4.4.1) is valid even if φ_t and (4.1.1) are replaced respectively by ψ_t and (4.1.2). This is the original local extension theorem proved by Aubin [2].

§ 5. The K-energy map along the solutions of generalized Aubin's equations

Recall that (4.1.1) has a unique solution φ_0 at t = 0 (see (4.3.2)). By using an explicit description (cf. (5.1)) of the K-energy map μ along the solutions of (4.1.1), we shall show that any φ_{τ} satisfying (4.1.1) at t = τ ($\tau \neq 1$) uniquely extends to a smooth 1-parameter family { $\varphi_t \mid 0 \leq t \leq \tau$ } of solutions of (4.1.1). Note that this fact in particular shows that (4.1.1) admits at most one solution at t = τ for $0 \leq \tau < 1$ (cf. (5.3)). The same technique enables us to show that if μ is bounded from below, then φ_0 uniquely extends to a smooth 1-parameter family { $\varphi_t \mid 0 \leq t < 1$ } of solutions of (4.1.1) (cf. (5.7)).

Theorem (5.1). Let $\{\varphi_t \mid a \leq t \leq b\}$ be an arbitrary smooth 1-parameter family of solutions of (4.1.1) in H. For brevity, we put

$$\omega^{(t)} := \omega_0(\varphi_t) , \quad I_t := I(\omega_0, \omega^{(t)}) \quad (= I(0, \varphi_t)),$$

$$J_t := J(\omega_0, \omega^{(t)}) \quad (=J(0, \varphi_t)).$$

Then on [a,b],

$$\frac{d\mu(\omega^{(t)})}{dt} = -(1-t) \frac{d}{dt} (I_t - J_t) \leq 0.$$

<u>Proof:</u> By $R(\varphi_t) = (1 - t)\omega_0 + t\omega^{(t)} = \omega_0(\varphi_t) - \sqrt{-1} (1 - t)\partial \overline{\partial} \varphi_t$ (cf. (4.2.1)), we have $\sigma(\varphi_t) = n - (1 - t) \varphi_t \varphi_t$. Hence,

$$\frac{d}{dt} \mu(\omega^{(t)}) = \frac{d}{dt} M(0, \varphi_t) = \int_X (1-t) \dot{\varphi}_t^{(\Box} \varphi_t^{(\Box} \varphi_t) \omega_0^{(\varphi_t)} / v$$
$$= -(1-t) \frac{d}{dt} (I_t - J_t) \quad (cf. \ (1.6.5)).$$

On the other hand, differentiating (4.1.1) by t, we obtain

$$\mathbf{p}_{\varphi_{t}} \mathbf{\phi}_{t} + \mathbf{t} \mathbf{\phi}_{t} + \mathbf{\phi}_{t} + \mathbf{c}_{t} = 0$$

for some constant $C_t \in \mathbb{R}$ on X. Combining this with (1.6.5), we now see that

$$\frac{\mathrm{d}}{\mathrm{d}t} (\mathbf{I}_t - \mathbf{J}_t) = \int_X (\mathbf{u}_{\varphi_t} \dot{\varphi}_t + t\dot{\varphi}_t) (\mathbf{u}_{\varphi_t} \dot{\varphi}_t) \omega_0 (\varphi_t)^n / v \ge 0 ,$$

where the last inequality is a straightforward consequence of (first eigen value of $-\sigma_{\varphi_+}$) > t.

(5.2) (i) Fix $\alpha \in \mathbb{R}$ with $0 < \alpha < 1$. Let φ_{τ} be a solution of (4.1.1) at $t = \tau$ (where $\tau \neq 0,1$). A smooth family $\{\varphi_t \mid \sigma < t \leq \tau\}$ (resp. $\{\varphi_t \mid \tau \leq t < \sigma\}$) of solutions of (4.1.1) is said to be <u>maximal</u> if for any sequence $t_j \in (\sigma, \tau]$ (resp. $[\tau, \sigma)$) (j = 1,2,...) with lim $t_j = \sigma$, the corresponding sequence $\{\varphi_{t_j}\}$ in H does not converge to any point of $H^{2,\alpha}$ in $C^{2,\alpha}$ norm.

(ii) Suppose $E \neq \phi$. Then to each $\theta \in E$, we can uniquely associate a function $\lambda_{\theta} \in H$ such that $\theta = \omega_0(\lambda_{\theta})$ and that λ_{θ} satisfies (4.1.2) at t = 1, i.e., $A(\lambda_{\theta}) = -\lambda_{\theta} + f$. An element θ of E is said to be <u>minimally Einstein</u> on the Kähler manifold (X, ω_0) if for some $\varepsilon > 0$, there exists a smooth family $\{\psi_t \mid 1 - \varepsilon \leq t \leq 1\}$ of solutions of (4.1.2) such that $\psi_1 = \lambda_{\theta}$.

<u>Theorem</u> (5.3). Let $0 < \tau < 1$. Then any solution φ_{τ} of (4.1.1) at $t = \tau$ uniquely extends to a smooth family $\{\varphi_t \mid 0 \le t \le \tau\}$ of solutions of (4.1.1). In particular (4.1.1) admits at most one solution in H at $t = \tau$.

Corollary (5.4). (i) There exists at most one $\theta \in E$ which is minimally Einstein on the Kähler manifold (X, ω_0) .

(ii) Suppose that $\theta \in E$ is minimally Einstein on the Kähler manifold (X, ω_0) . Then $M(\theta, \omega_0(\phi_0)) \ge 0$.

(5.5) <u>Proof</u> of (5.3): The required uniqueness is immediate from (4.4.1), once the existence of an extension is proven. We therefore assume, for contradiction, that any such extension is impossible. Then by (4.4.1), we have a maximal smooth family $\{\varphi_t \mid \sigma < t \le \tau\}$ of solutions of (4.1.1) for some $0 \le \sigma \in \mathbb{R}$. In this proof, we always denote by $t \in \mathbb{R}$ an arbitrary number satisfying $\sigma < t \le \tau$, and by a "constant" a positive real number which doesn't depend on either t or $x \in X$. The proof is now divided into two steps.

Step 1: By (1.6.4) and Theorem (5.1),

(5.5.1) $0 \le I_{\pm} \le (n+1) (I_{\pm} - J_{\pm}) \le (n+1) (I_{\pm} - J_{\pm}).$

We put $F_t := -t\phi_t - L(0,\phi_t) + f \in C^{\infty}(X)_{\mathbb{R}}$. Then by (4.1.1),

$$\int_{X} \omega_0^{n} = \int_{X} \exp(A(\varphi_t)) \omega_0^{n} = \int_{X} \exp(F_t) \omega_0^{n}.$$

Applying the mean value theorem, we have $F_t(x_t) = 0$ for some $x_t \in X$. Therefore for every $x \in X$, $|F_t(x)| = |F_t(x) - F_t(x_t)| = |-t(\varphi_t(x) - \varphi_t(x_t)) + f(x) - f(x_t)|$ $\leq t (Osc \varphi_t) + 2||f||_{C^0}$ $\leq tI_t + 2n (tK_0V_0 + (n-1)\beta(n)\pi^2) + 2||f||_{C^0} (cf. (3.6), (4.2)).$ Hence by (5.5.1), there exists a constant K_1 such that $(5.5.2) ||F_t||_{C^0} \leq K_1.$ Since $A(\varphi_t) = F_t$, a result of Yau [14] (see also Bourguignon et al. [4; VII]) now asserts that $(5.5.3) Osc \varphi_t \leq K_2$ for some constant K_1 . Dut \widetilde{A} is a result of $V_1 = CO^{\infty}(V)$.

for some constant K_2 . Put $\widetilde{\phi}_t:=\phi_t-\phi_t(x_t)\in C^\infty(X)_{\rm I\!R}$. In view of

$$0 = F_{t}(x_{t}) = - (1+t)\varphi_{t}(x_{t}) - L(0,\tilde{\varphi}_{t}) + f(x_{t}) \quad (cf.(1.4.3)),$$

we obtain

$$(5.5.4) |\varphi_{t}(\mathbf{x}_{t})| \leq |L(0,\widetilde{\varphi}_{t})| + ||f||_{C^{0}}.$$

Since $||\widetilde{\varphi}_{t}||_{C} 0 \leq K_{2}$, it follows that

(5.5.5)
$$|L(0,\widetilde{\varphi}_{t})| = |\int_{0}^{1} (\int_{X} \widetilde{\varphi}_{t} \omega_{0} (s\widetilde{\varphi}_{t})^{n}/V) ds|$$

 $\leq \int_{0}^{1} (\int_{X} K_{2} \omega_{0} (s\widetilde{\varphi}_{t})^{n}/V) ds = K_{2}.$

Hence by (5.5.3) and (5.5.4),

(5.5.6)
$$||\varphi_{t}||_{C} 0 \leq \kappa_{3}$$

for some constant K_3 . Since $L(0, \varphi_t) = \varphi_t(X_t) + L(0, \widetilde{\varphi}_t)$, a combination of (5.5.5) and (5.5.6) now provides us with a constant K_4 such that

(5.5.7) $|| -L(0, \varphi_{t}) + f ||_{C} 3 \leq K_{4}.$

<u>Step</u> 2: Recall that $A(\phi_t) + t\phi_t = -L(0,\phi_t) + f$. By (5.5.6) and (5.5.7), we have constants K_5, K_6, K_7 such that

$$\begin{aligned} || \varphi_t ||_C^{2,\alpha'} &\leq \kappa_5 \quad \text{for all } \alpha' \quad \text{with } \alpha < \alpha' < 1, \\ \kappa_6 \omega_0 &\leq \omega_0 (\varphi_t) \leq \kappa_7 \omega_0 \quad (\text{cf. Aubin [1; p. p. 151-154]}. \end{aligned}$$

We now choose an arbitrary decreasing sequence $t_j \in (\sigma, \tau]$, $j = 1, 2, \ldots$ such that $\lim t_j = \sigma$. Then by Ascoli theorem, there exists a convergent subsequence of $\{\varphi_{t_j}\}$ in $C^{2,\alpha}$, which leads to a contradiction to the maximality of $\{\varphi_+ \mid \sigma < t \le \tau \}$.

(5.6) <u>Proof of</u> (5.4): Let $\theta \in E$ be minimally Einstein on the Kähler manifold (X, ω_0) . Then in view of Remark (4.1.3), there exists a smooth family $\{\varphi_t | 1 - \varepsilon \leq t \leq 1\}$ of solutions of (4.1.1) such that $\theta = \omega_0(\varphi_1)$. By (5.3) above, $\{\varphi_t | 1 - \varepsilon \leq t < 1\}$ uniquely extends to <u>the</u> smooth 1-parameter family $\{\varphi_t \mid 0 \le t < 1\}$ of solutions of (4.1.1). Then (i) immediately follows from

$$\varphi_{1} = \lim_{t \to 1} \varphi_{t} \quad \text{in } C^{K}(X)_{IR} \quad (k \ge 0),$$

and (ii) from $M(\theta, \omega_0(\phi_0)) = \mu(\omega_0(\phi_0)) - \mu(\omega_0(\phi_1)) \ge 0$ (cf. (5.1)).

The following theorem, which we don't need later, is of some interest from the viewpoint of understanding the K-energy map μ . We therefore give it together with a proof.

<u>Theorem</u> (5.7). Let $S := \{\omega_0(\varphi) \mid \varphi \in H \text{ satisfies } A(\varphi) = -t\varphi$ -L(0, φ) + f for some t $\in [0,1)$ }. Suppose that μ is bounded from below on S. Then φ_0 uniquely extends to a smooth 1-parameter family $\{\varphi_+ \mid 0 \leq t < 1\}$ of solutions of (4.1.1).

<u>Proof</u>: We assume, for contradiction, that there exists a maximal smooth family $\{\varphi_t \mid 0 \le t < \sigma\}$ of solutions of (4.1.1) for some $\sigma < 1$. Let K in B be the infimum of μ on S. Then whenever $0 \le t < \sigma$, $\omega^{(t)} := \omega_0(\varphi_t)$ belongs to S and in particular $\mu(\omega^{(t)}) \ge K$. For each such t, we infer from Theorem (5.1) that

$$\begin{split} \mathbf{I}_{t} &- \mathbf{J}_{t} = \int_{0}^{t} \frac{-1}{1-s} \frac{d\mu(\boldsymbol{\omega}^{(s)})}{ds} \, \mathrm{ds} \, + \, (\mathbf{I}_{0} \, - \, \mathbf{J}_{0}) \\ &\leq \int_{0}^{t} \frac{-1}{1-\sigma} \frac{d\mu(\boldsymbol{\omega}^{(s)})}{ds} \, \mathrm{ds} \, + \, (\mathbf{I}_{0} \, - \, \mathbf{J}_{0}) \leq \frac{1}{1-\sigma} \, (\mu(\boldsymbol{\omega}_{0}(\boldsymbol{\varphi}_{0})) - \mathbf{K}) + (\mathbf{I}_{0} - \mathbf{J}_{0}) \, , \end{split}$$

where we used the notation in (5.1). Thus $I_t - J_t$ ($0 \le t < \sigma$) is bounded from above. The rest of the proof is quite similar to (5.5). § 6. Lemmas for choosing a good gauge.

(6.1) Throughout this section, we use the same notation as in § 2, and fix an arbitrary G-orbit O in E, assuming $E \neq \phi$. To each $\theta \in \mathbf{O}$, we can uniquely associate a function $\lambda_{\theta} \in H$ such that $\theta = \omega_{0}(\lambda_{\theta})$ and that $A(\lambda_{\theta}) =$ = $-\lambda_{\theta}$ + f (cf. (ii) of (5.2)). Then $\theta \mapsto \lambda_{\theta}$ defines a bijection between **0** and $\overline{\mathbf{0}} := \{\lambda_{\theta} \mid \theta \in \mathbf{0}\}$. We endow $\overline{\mathbf{0}}$ with the topology naturally induced from the $C^{2,\alpha}$ norm of H. This defines a topology on $\mathbf{0}$, in terms of which, the G-action on O is clearly continuous. Hence our topology on O coincides with the natural topology of the homogeneous space G/K_{θ} via the identification $\mathbf{O} \cong G/K_{\theta}$ (cf. (2.1)). Recall that for each $\ \phi \in H_A$, we have the corresponding 1-parameter group $y_{\theta t}^{\phi} = \exp(tY_{\theta IR}^{\phi})$, $t \in IR$. For simplicity, we put $\theta(t) := (y_{\theta t}^{\phi}) * \theta$ and $\lambda(t) := \lambda_{\ell(t)}$. Then by (2.2.4) and $\omega_0(\lambda(t)) = \theta(t)$, we have $\lambda(0) = \varphi + C$ for some $C \in \mathbb{R}$. On the other hand, differentiating the identity $A(\lambda(t)) = -\lambda(t) + f$ by t at t = 0, we obtain $\Box_{\theta}(\dot{\lambda}(0)) = -\dot{\lambda}(0)$. Hence $\int_{\mathbf{X}} \dot{\lambda}(0) \theta^{n} = 0 = \int_{\mathbf{X}} \varphi \theta^{n}$ and this implies $\dot{\lambda}(0) = \varphi$. Thus we established the following identification (cf. (2.2.3)):

(6.1.1)
$$T_{\lambda_{\theta}}(\overline{O}) \cong T_{\theta}(O) \cong H_{\theta}$$

 $\dot{\lambda}(0) = \phi \leftrightarrow \dot{\theta}(0) = \sqrt{-1}\partial \overline{\partial}\phi \leftrightarrow \phi.$

The purpose of this section is to prove the following lemmas:

Lemma (6.2). The C^{∞} function ι defined by

 $\iota: \mathbf{O} \longrightarrow \mathbf{R}$ $\theta \longmapsto \iota(\theta) := \mathrm{I}(\omega_0, \theta) - \mathrm{J}(\omega_0, \theta) \ (\geq 0)$

is a proper map. In particular its minimum is always attained at some point of the orbit **O**.

Lemma (6.3). Let $\theta \in \mathbf{O}$. Then the following are equivalent.

(i) θ is a critical point of ι ;

(ii) $\int_{X} \lambda_{\theta} \varphi \theta^{n} = 0$ for all $\varphi \in H_{\theta}$;

(iii) θ is expressible as $\omega_0(\psi)$ for some function $\Psi \in H$ such that $\int_X \varphi \psi \theta^n = 0$ for all $\varphi \in H_{\theta}$.

Lemma (6.4). Let $\theta \in \mathbf{O}$ be a critical point of ι . Then the Hessian (Hess ι)_{θ} of ι at the point θ is given by

(Hess 1)_{$$\theta$$} (φ', φ'') = $\int_{X} (1 + \frac{1}{2} \Box_{\theta} \lambda_{\theta}) \varphi' \varphi'' \theta'' / V$

for all $\phi', \phi'' \in H_{\theta} (\cong T_{\theta}(0)).$

(6.5) <u>Proof of</u> (6.2): By the well-known regularity theorem applied to the equation $A(\psi) = -\psi + f$, the proof is reduced *) to showing

(6.5.1) given a real number $r \ge 0$, one can always find positive numbers $K_r, K'_r, K''_r \in \mathbb{R}$ such that

 $\frac{\|\psi\|}{C^{2},\alpha} \leq K_{r} \quad \underline{and} \quad K_{r}'\omega_{0} \leq \omega_{0}(\psi) \leq K_{r}''\omega_{0}$ <u>hold simultaneously for all</u> $\psi \in \overline{\mathbf{O}}$ <u>satisfying</u> $|\iota_{\theta}(\omega_{0}(\psi))| \leq r$.

^{*)} This reduction is easily obtained from the following standard fact: O is a connected component of E (see Calabi's article "Extremal Kähler matrics II"; in "Differential Geometry and Complex Analysis" dedicated to H.E. Rauch, Springer-Verlag, 1985).

Fix an arbitrary element Ψ of $\overline{\mathbf{O}}$ as in (6.5.1). Then by (1.6.4), $I(0,\Psi) \leq (n+1)r$. In view of (4.2), we have $\Psi \in H^{(1)}$ and hence by (3.6),

(6.5.2) Osc $\psi \leq K$, where $K := (n+1)r + 2n(K_0V_0 + (n-1)\beta(n)\pi^2)$. On the other hand, from $A(\psi) = -\psi + f$, we obtain

$$\int_{X} \omega_0^n = \int_{X} \exp \left(A(\psi) \right) \omega_0^n \neq \int_{X} \exp \left(-\psi + f \right) \omega_0^n.$$

Therefore by the mean value theorem, there exists a point $x \in X$ such that $\psi'(x) = f(x)$. Together with (6.5.2), we have

$$\|\Psi\|_{C^{0}} \leq K + \|f\|_{C^{0}}$$

Then applying standard arguments (cf. Aubin [1;pp. 151 - 154]), we obtain K_r, K_r', K_r'' as required in (6.5.1).

(6.6) <u>Proof of</u> (6.3): (ii) and (iii) are clearly equivalent. To see the equivalence of (i) and (ii), we fix an arbitrary $\varphi \in H_{\theta} (\cong T_{\theta} (\mathbf{O}))$ with its corresponding 1-parameter families $\{\theta(t) \in \mathbf{O} \mid t \in \mathbb{R}\}$ and $\{\lambda(t) \in \overline{\mathbf{O}} \mid t \in \mathbb{R}\}$ as in (6.1). Then in view of (1.6.5),

$$\frac{d}{dt} \iota(\theta(t))|_{t=0} = \frac{d}{dt} (I(0,\lambda(t)) - J(0,\lambda(t)))|_{t=0}$$
$$= -\int_{X} \lambda_{\theta} u_{\theta}(\dot{\lambda}(0)) \theta^{n} / V = -\int_{X} \lambda_{\theta} (u_{\theta} \phi) \theta^{n} / V \quad (cf. (6.1.1))$$
$$= \int_{X} \lambda_{\theta} \phi \theta^{n} / V \quad (because \ \phi \in H_{\theta}).$$

The required equivalence is now straightforward.

(6.7) <u>Proof of</u> (6.4): Let $\{\lambda_{s,t} | (s,t) \in [-\varepsilon,\varepsilon] \times [-\varepsilon,\varepsilon]\} (\varepsilon > 0)$

be a smooth family of functions in \overline{O} satisfying the following conditions:

(i)
$$\lambda_{0,0} = \lambda_{\theta}$$
,
(ii) $\left(\frac{\partial}{\partial s}\lambda_{s,t}\right) | (s,t) = (0,0) = \phi'$,
(iii) $\left(\frac{\partial}{\partial t}\lambda_{s,t}\right) | (s,t) = (0,0) = \phi''$.

We shall denote $\Box_{\lambda s,t}$ (resp. $\omega_0(\lambda_{s,t})$) by $\Box_{s,t}$ (resp. $\theta_{s,t}$) for brevity. Differentiating the identity $A(\lambda_{s,t}) + \lambda_{s,t} - f = 0$ by t, we obtain

(6.7.1)
$$\Box_{s,t}(\frac{\partial}{\partial t}\lambda_{s,t}) + (\frac{\partial}{\partial t}\lambda_{s,t}) = 0.$$

Further differentiation by s yields

$$-\langle \partial \overline{\partial} \left(\frac{\partial}{\partial s} \lambda_{s,t} \right), \ \partial \overline{\partial} \left(\frac{\partial}{\partial t} \lambda_{s,t} \right) \rangle_{s,t} + \left(\Box_{s,t} + 1 \right) \left(\frac{\partial^2}{\partial s \partial t} \lambda_{s,t} \right) = 0,$$

where we denote $\langle , \rangle_{\theta}$ (cf. (2.3)) simply by $\langle , \rangle_{s,t}$. Evaluating this at (s,t) = (0,0), we obtain

$$(\Box_{\theta} + 1) \left(\frac{\partial}{\partial s \partial t} \lambda_{s,t} \right) | (0,0) = \langle \partial \bar{\partial} \phi' \rangle_{\theta}.$$

Together with (2.3.1), it then follows that

(6.7.2)
$$\left(\frac{\partial^2}{\partial s \partial t} \lambda_{s,t}\right) | (0,0) = \langle \partial \phi', \partial \phi'' \rangle_{\theta} = \langle \partial \phi', \partial \phi'' \rangle_{\theta}$$

(modulo $H_{\theta}^{(\mathbf{C})}$).

We can now finish the proof by the following computation:

$$= \int_{X} \{\varphi'\varphi'' + 2^{-1}((\Box_{\theta}\varphi')\varphi'' + (\Box_{\theta}\varphi'')\varphi' + \langle \partial\varphi', \partial\varphi'' \rangle_{\theta} + \langle \partial\varphi'', \partial\varphi'' \rangle_{\theta}\} \theta^{n}/V$$
$$= \int_{X} \{\varphi'\varphi'' + \frac{1}{2}\lambda_{\theta}\Box_{\theta}(\varphi'\varphi'')\} \theta^{n}/V = \int_{X} (1 + \frac{1}{2}\Box_{\theta}\lambda_{\theta})\varphi'\varphi'' \theta^{n}/V.$$

§ 7. Unfolding the singularity at t = 1 by a bifurcation technique.

We again assume that $E \neq \phi$ and fix an arbitrary G-orbit **O** in *E*. Using the same notation as in § 6, we fix a critical point θ of the mapping $\iota: \mathbf{O} \rightarrow \mathbf{R}$ (cf. (6.2)). The purpose of this section is to find out a good sufficient condition for θ to be minimally Einstein (cf. (ii) of (5.2)). Fixing $\alpha \in \mathbf{R}$ with $0 < \alpha < 1$, we set

$$H_{\theta,k}^{\perp} := \{ \psi \in C^{k,\alpha}(X)_{IR} \mid \int_{X} \phi \psi \theta^{n} = 0 \text{ for all } \phi \in H_{\theta} \},$$
$$k = 0, 1, 2, \dots$$

Recall that, corresponding to θ , we have the function $\lambda_{\theta} \in H$ with the following properties (cf. (6.1), (6.3)):

- (i) $\theta = \omega_0(\lambda_\theta)$,
- (ii) $A(\lambda_{\theta}) = -\lambda_{\theta} + f$,
- (iii) $\lambda_{\theta} \in H_{\theta,k}^{\perp}$.

Let $k \ge 2$, and we now consider the mapping

$$\Phi: \mathbb{R} \times C^{k,\alpha}(X)_{\mathbb{R}} \longrightarrow C^{k-2,\alpha}(X)_{\mathbb{R}} ,$$

$$(t,u) \longmapsto \Phi(t,u) := A(u) + tu - f.$$

Note that, by the well-known regularity theorem, any $v \in H^{k,\alpha}$

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(cf. (4.4)) satisfying $\Phi(t,v) = 0$ for some t is automatically in H. Let $P:C^{0,\alpha}(X)_{\mathbb{R}} \cong H_{\theta} \oplus H_{\theta,0}^{1} \longrightarrow H_{\theta}$ be the natural projection to the first factor. For each $u \in C^{k,\alpha}(X)_{\mathbb{R}}$, we write

(7.1.1)
$$u = \lambda_{\theta} + \varphi + \psi ,$$

with $\varphi := P(u - \lambda_{\theta}) \in H_{\theta}$ and $\Psi := (1 - P) (u - \lambda_{\theta}) \in H_{\theta,k}^{\perp}$. Now the equation

 $(7.1.2) \quad \Phi(t,u) = 0$

is written in the form

$$\begin{split} & \mathbb{P}\Phi\left(\mathsf{t},\lambda_{\theta}^{-}+\phi+\psi\right) = 0 \quad \text{and} \quad \Psi(\mathsf{t},\phi,\psi) = 0 \ , \\ & \text{where} \quad \Psi: \mathbf{R} \times \ \mathrm{H}_{\theta}^{-} \times \ \mathrm{H}_{\theta,k}^{\perp} \longrightarrow \ \mathrm{H}_{\theta,k-2}^{\perp} \quad \text{is the mapping defined} \\ & \text{by} \end{split}$$

$$\Psi(t, \varphi, \psi) := (1 - P) \Phi(t, \lambda_{\Theta} + \varphi + \psi) \quad ((t, \varphi, \psi) \in \mathbb{R} \times \mathbb{H}_{\Theta} \times \mathbb{H}_{\Theta}^{\perp}, k).$$

Then $\Psi(1,0,0) = 0$ and the Frechet derivative $D_{\psi}\Psi|(1,0,0)$ of Ψ with respect to Ψ at $(t,\varphi,\Psi) = (1,0,0)$ is

$$\psi' \in H^{\perp}_{\theta, k} \longmapsto (D_{\psi} \Psi) \mid (1, 0, 0) (\psi') = (\Box_{\theta} + 1) \psi' \in H^{\perp}_{\theta, k-2},$$

which is invertible. Therefore the implicit function theorem enables us to obtain a smooth mapping $(t, \varphi) \in U \mapsto \psi_{t,\varphi} \in H^{\perp}_{\theta,k}$ of a small neighbourhood U of (1,0) in $\mathbb{R} \times H_{\theta}$ to the Banach space $H^{\perp}_{\theta,k}$ such that (i) $\psi_{1,0} = 0$, (ii) $|| \psi_{t,\varphi} ||_C k, \alpha \leq \delta$ on U for some $\delta > 0$, and (iii) $\Psi(t,\varphi,\psi) = 0$ (where $||\psi||_C k, \alpha \leq \delta$) is, as an equation in $\psi \in C^{k,\alpha}(X)_{\mathbb{R}}$, uniquely solvable in the form $\psi = \psi_{t,\varphi}$ on U.

Differentiating the identity $\Psi(t,\phi,\psi_{t,\phi}) = 0$ at $(t,\phi) = (1,0)$, we obtain

$$(7.1.3) \qquad (\Box_{\theta} + 1) \left(\frac{\partial}{\partial t} \psi_{t,\phi}\right) (1,0) = -\lambda_{\theta} ,$$

(7.1.4)
$$(D_{\varphi}^{\psi}_{t,\varphi})|_{(1,0)}(\varphi') = 0$$
 for all $\varphi' \in H_{\theta}$,

where $(D_{\varphi}\psi_{t,\varphi})|_{(1,0)}: H_{\theta} \rightarrow H_{\theta,k}^{\perp}$ denotes the Frechet derivative of $\psi_{t,\varphi}$ with respect to φ at the point $(t,\varphi) = (1,0)$. Then the equation (7.1.2), on a small neighbourhood of λ_{θ} , reduces to

(7.1.5)
$$\Phi_0(t,\phi) = 0$$
 (with $u = \lambda_0 + \phi + \psi_{t,\phi}$),

where we put $\Phi_0(t, \varphi) := P\Phi(t, \lambda_{\theta} + \varphi + \psi_{t, \varphi})$ for $(t, \varphi) \in U$. Recall that $\Phi(1, u) = 0$ for all $u \in \overline{O}$. Hence $\Phi_0 = 0$ on $\{t = 1\}$ and therefore the mapping

$$(t, \varphi) \in U_{|\{t+1\}} \mapsto \Phi_1(t, \varphi) := \Phi_0(t, \varphi) / (t-1) \in H_{\theta}$$

naturally extends to a smooth map: $U \rightarrow H_{\theta}$ (denoted by the same Φ_1) of finite dimensional sets. In view of (7.1.3), we obtain

$$\Phi_1(1,0) = (\partial \Phi_0 / \partial t)(1,0) = 0.$$

Furthermore we shall later show that the Frechet derivative $D_{\phi} \phi_1 |_{(1,0)} : H_{\theta} \rightarrow H_{\theta}$ of ϕ_1 with respect to ϕ at $(t,\phi) = (1,0)$ is written in the following form:

Lemma (7.2). For all $\varphi', \varphi'' \in H_{\theta}$ (where θ is a critical point of 1),

Suppose now that (Hess 1)₀ : H₀ × H₀ \rightarrow **R** is a nondegenerate bilinear form. Then by this lemma, $D_{\phi} \phi_{1|(1,0)}$ is invertible, and the implicit function theorem shows that the equation $\phi_{1}(t,\phi) = 0$ in ϕ is uniquely solvable in a neighbourhood of (1,0) to produce a smooth curve $\{\phi(t) \mid 1 - \varepsilon \leq t \leq 1\}$ ($\varepsilon > 0$) in H₀ such that (i) $\phi(1) = 0$ and (ii) $\phi_{1}(t,\phi(t)) = 0$ ($1 - \varepsilon \leq t \leq 1$). Therefore, in view of (7.1.5), we have $\phi(t,\lambda_{\theta} + \phi(t) + \psi_{t,\phi(t)}) = 0$ ($1 - \varepsilon \leq t \leq 1$), and hence $\{\psi_{t} := \lambda_{\theta} + \phi(t) + \psi_{t,\phi(t)} \mid 1 - \varepsilon \leq t \leq 1\}$ forms a smooth 1-parameter family of solutions of (4.1.2) in H with $\psi_{1} = \lambda_{\theta}$, i.e., θ is minimally Einstein on the Kähler manifold (X, ω_{0}) (cf. (ii) of (5.2)). Thus we obtained

<u>Theorem (7.3).</u> Every critical point θ of ι with nondegenerate Hessian is minimally Einstein on the Kähler <u>manifold</u> (X, ω_0) .

We shall finally show Lemma (7.2).

(7.4) Proof of (7.2): By (7.1.4), using the notation in (2.3), we have

$$D_{\varphi} \Phi_{1} | (1,0) (\varphi') = (D_{\varphi} \frac{\partial}{\partial t} \Phi_{0}) | (1,0) (\varphi')$$
$$= \varphi' - P < \partial \overline{\partial} (\frac{\partial}{\partial t} \Psi_{t,\varphi} | (1,0)), \ \partial \overline{\partial} \varphi' >_{\theta}.$$

Hence, it follows that

$$= \int_{X} (1 + \frac{1}{2} \mathfrak{n}_{0} \lambda_{0}) \varphi' \varphi'' \theta'' n! \quad (\text{see the end of (6.7)}).$$

Since $V = n V_0$, this completes the proof (cf. (1.4), (6.4)).

<u>Remark</u> (7.4.1). If $\theta \in \mathbf{O}$ is a point where ι attains its minimum (cf. (6.2)), then (Hess ι)_{θ} is positive semidefinite. In the next section, we shall realize a critical point of ι with positive definite Hessian via a small change of our presently fixed ω_0 . § 8 . Proof of Main Theorem

(8.1) <u>Proof of (i) of Theorem A</u>: Fix an element $\tilde{\omega}$ of K^+ and a G-orbit O in E arbitrarily. In this section, we write

$$\omega_0 = \omega_0^{\varepsilon}$$

regarding ω_0 as a function of $\varepsilon \in [0,1]$. Hence the corresponding $f, \varphi_0, \iota, H, \omega_0(\varphi)$ and $A(\varphi)$ (where $\varphi \in C^{\infty}(X)_{\mathbb{R}}$) will be written respectively as $f_{\varepsilon}, \varphi_{0;\varepsilon}, \iota_{\varepsilon}, H^{\varepsilon}, \omega_0^{\varepsilon}(\varphi)$ and $A^{\varepsilon}(\varphi)$ (see (1.7.1), (4.3.2), (6.2), (1.1), (1.1.1) and (4.1)). We first consider the special case $\varepsilon = 0$ and then go to the general situation $\varepsilon > 0$.

<u>Case</u> 1: $\varepsilon = 0$. Put $\omega_0^0 := \mathbb{R}(\widetilde{\omega})$. Then $\iota_0: \mathbf{O} \to \mathbb{R}$ takes its minimum at some point θ of \mathbf{O} (cf. (6.2)). Corresponding to this θ , there uniquely exists a function $\lambda_{\theta;0} \in H^0$ such that $\theta = \omega_0^0(\lambda_{\theta;0})$ and that $\mathbb{A}^0(\lambda_{\theta;0}) = -\lambda_{\theta;0} + f_0$ (cf. (ii) of (5.2)). Recall that \mathbb{H}_{θ} is Ker $(\mathbb{D}_{\theta} + 1)$ in $\mathbb{C}^{\infty}(X)_{\mathbb{R}}$. Then by (6.3),

(8.1.1) $\int_{X} \lambda_{\theta;0} \varphi \, \theta^n = 0 \quad \text{for all} \quad \varphi \in H_{\theta},$

and the bilinear form (Hess ι_0) $_{\theta}$:H $_{\theta} \times H_{\theta} \longrightarrow \mathbb{R}$ is positive semidefinite.

<u>Case</u> 2: $\varepsilon > 0$. In this case, we set $\omega_0^{\varepsilon} := (1 - \varepsilon)\omega_0^0 + \varepsilon \theta = \omega_0^0(\varepsilon \lambda_{\theta;0})$. Again by (ii) of (5.2), one obtains a function $\lambda_{\theta;\varepsilon} \in H^{\varepsilon}$

uniquely determined by the identities
$$\theta = \omega_0^{\varepsilon}(\lambda_{\theta;\varepsilon})$$
 and
 $A^{\varepsilon}(\lambda_{\theta;\varepsilon}) = -\lambda_{\theta;\varepsilon} + f_{\varepsilon}$. Then in view of $\omega_0^0(\lambda_{\theta;0}) = \theta =$
 $= \omega_0^0(\lambda_{\theta;\varepsilon} + \varepsilon\lambda_{\theta;0})$, we have
(8.1.2) $\lambda_{\theta;\varepsilon} = (1 - \varepsilon)\lambda_{\theta;0} + C_{\varepsilon}$ for some $C_{\varepsilon} \in \mathbb{R}$.
Hence $\int_X \lambda_{\theta;\varepsilon} \varphi \theta^n = \int_X (1 - \varepsilon)\lambda_{\theta;0} \varphi \theta^n = 0$ if $\varphi \in H_{\theta}$ (see
(8.1.1)). Therefore by (6.3), θ is a critical point of
 $\lambda_{\varepsilon}: \mathbf{O} \longrightarrow \mathbf{R}$. Moreover for all $0 \neq \varphi \in H_{\theta}$,

$$(\text{Hess } \iota_{\varepsilon})_{\theta}(\phi,\phi) = \int_{X} (1 + \frac{1}{2} \Box_{\theta} \lambda_{\theta;\varepsilon}) \phi^{2} \theta^{n} / V \quad (\text{cf. (6.4)})$$
$$= (1 - \varepsilon) \int_{X} (1 + \frac{1}{2} \Box_{\theta} \lambda_{\theta;0}) \phi^{2} \theta^{n} / V + \varepsilon \int_{X} \phi^{2} \theta^{n} / V \quad (\text{cf. (8.1.2)})$$

$$= (1 - \varepsilon) (\text{Hess } \iota_0)_{\theta} (\phi, \phi) + \varepsilon \int_X \phi^2 \theta^n / v > 0.$$

Theorem (7.3) now shows that θ is minimally Einstein on the Kähler manifold (X, ω_0^{ϵ}) . In particular, by (ii) of (5.4),

$$(8.1.3) \qquad M(\theta, \omega_0^{\varepsilon}(\phi_{0;\varepsilon})) \ge 0.$$

By (4.3.2), $R(\omega_0^{\varepsilon}(\phi_{0;\varepsilon})) = \omega_0^{\varepsilon}$. We also have $R(\widetilde{\omega}) = \omega_0^{0}$. Note that $\omega_0^{\varepsilon} \rightarrow \omega_0^{0}$ in $C^{0,\alpha}$ (as $\varepsilon + 0$). Then by (1.1.3), $\omega_0^{\varepsilon}(\phi_{0;\varepsilon}) \rightarrow \widetilde{\omega}$ in $C^{2,\alpha}$. Let $\varepsilon + 0$ in (8.1.3). By the continuity of M, we have

$$M(\theta,\widetilde{\omega}) \ge 0$$
, i.e., $\mu^+(\theta) \le \mu^+(\widetilde{\omega})$.

Recall that μ^+ is a constant function on **O** (cf. [9]). Since both $\tilde{\omega} \in K^+$ and the G-orbit **O** in E are arbitrary, it now follows that μ^+ : $K^+ \rightarrow \mathbb{R}$ takes its absolute minimum C on E. Note that E is the set of all critical points of μ^+ (cf. [9]). Hence

 $\{\omega \in K^+ \mid \mu^+(\omega) = C\} = E,$

because, otherwise, at some point $\widetilde{\omega}$ of K^+ with $\widetilde{\omega} \notin E$, the function μ^+ would take its critical value C in contradiction to $\widetilde{\omega} \notin E$.

(8.2) <u>Proof of (ii) of Theorem</u> A: Let O' and O" be arbitrary G-orbits in E. Then from the argument of (8.1) applied to the orbit O', we see the following:

(8.2.1) For a suitable choice of $\omega_0^{\prime} \in K$, the function $\iota^{\prime}: \omega \in \mathbf{O}^{\prime} \mapsto \iota^{\prime}(\omega) := \mathbf{I}(\omega_0^{\prime}, \omega) - \mathbf{J}(\omega_0^{\prime}, \omega) \in \mathbb{R}$ has a critical point $\theta^{\prime} \in \mathbf{O}^{\prime}$ with positive definite Hessian.

Recall that the function $\iota^{"}: \omega \in \mathbf{O}^{"} \longmapsto \iota^{"}(\omega) := \mathbf{I}(\omega_{0}^{'}, \omega) - \mathbf{J}(\omega_{0}^{'}, \omega)$ $\in \mathbb{R}$ takes its minimum at some point $\theta^{"} \in \mathbf{O}^{"}$ (cf. (6.2)). We now put $\omega_{0}^{\varepsilon} := (1 - \varepsilon)\omega_{0}^{'} + \varepsilon \theta^{"}$, ($0 \le \varepsilon \le 1$). Again by the argument of (8.1) applied to $\mathbf{O}^{"}$, we have:

(8.2.2) θ " is minimally Einstein on the Kähler manifold (X, ω_0^{ε}) whenever $0 < \varepsilon \leq 1$.

We finally define $\iota_{\varepsilon}^{\prime}: \mathbf{O}^{\prime} \rightarrow \mathbf{R}$ by

 $\iota_{\varepsilon}^{*}(\omega) := \mathbf{I}(\omega_{0}^{\varepsilon}, \omega) - J(\omega_{0}^{\varepsilon}, \omega), \quad (\omega \in \mathbf{O}^{*}).$

Note that $\iota_{\varepsilon}^{\prime}$ converges to ι' , say in $C^{2,\alpha}$, as ε tends to zero. Fix a sufficiently small $\varepsilon > 0$. Then by (8.2.1), the function $\iota_{\varepsilon}^{\prime}$ takes its local minimum with positive definite Hessian at some point $\theta_{\varepsilon}^{\prime}$ of **O'** near θ' . In view of Theorem (7.3), one finds that $\theta_{\varepsilon}^{\prime}$ is also minimally Einstein on the Kähler manifold $(X, \omega_{0}^{\varepsilon})$. Combining this with (8.2.2), we conclude from (i) of (5.4) that $\theta_{\varepsilon}^{\prime} = \theta''$. Thus, $\mathbf{O}^{\prime} = \mathbf{O}^{\prime'}$ and the proof is now complete.

Theorem A is valid even when $M: K \times K \rightarrow \mathbb{R}$ is replaced by $N: K \times K \rightarrow \mathbb{R}$ (cf. (1.5)). We conclude this section by showing

Corollary (8.3). Under the same assumption as in Theorem A, the mapping

$$\nu^{+}: K^{+} \longrightarrow \mathbb{R}$$
$$\omega \longmapsto \nu^{+}(\omega) := \mathbb{N}(\omega_{1}, \omega)$$

is bounded from below and takes its absolute minimum exactly on E.

<u>Proof:</u> By (1.8.1), $v^{+}(\omega) = \mu^{+}(\omega) + J(\omega, R(\omega)) - J(\omega_{1}, R(\omega_{1}))$ for every $\omega \in K^{+}$. Since both μ^{+} and $J(\omega, R(\omega))$ ($\omega \in K^{+}$) take their minimums exactly on E (see (1.6.3) and Theorem A), so does v^{+} . § 9. Proof of Theorems B and C.

(9.1) <u>Proof of Theorem</u> B: Recall that the natural Riemannian metric on E is characterized in terms of lengths of smooth paths in E as follows (cf. [10]):

For every smooth path $\Gamma = \{\widetilde{\gamma}_t \mid a \leq t \leq b\}$ in E, let $\{\gamma_t \mid a \leq t \leq b\}$ be the corresponding smooth path in H_0 (cf. (1.4)) uniquely determined by $\omega_0(\gamma_t) = \widetilde{\gamma}_t$ (t $\in [a,b]$). Then the length $L(\Gamma)$ of the path Γ in terms of the metric of E is defined by

$$L(\Gamma) := \int_{a}^{b} (\int_{X} (\dot{\gamma}_{t})^{2} \widetilde{\gamma}_{t}^{n} / v)^{1/2} dt.$$

In view of Theorem A, the proof is reduced to showing that Aut(X) acts isometrically on E. Hence it suffices to show

$$L(q*\Gamma) = L(\Gamma)$$
 (q \in Aut(X))

for every smooth path $\Gamma = \{\widetilde{\gamma}_t \mid a \le t \le b\}$ in E. Then even if $g \notin Aut^0(X)$, the same proof as in [10] goes through as follows:

Let φ_g be the function in \mathcal{H} uniquely determined by the properties $g^*\omega_0 = \omega_0(\varphi_g)$ and $\varphi_g + g^*\gamma_a \in \mathcal{H}_0$. We put $n_t := \varphi_g + g^*\gamma_t$ (astsb). Then $g^* \widetilde{\gamma}_t = \omega_{\widetilde{U}}(n_t)$ and $n_t \in \mathcal{H}_0$ for all t. Hence

$$L(g*\Gamma) = \int_{a}^{b} (\int_{X} (n_{t})^{2} \omega_{0}(n_{t})^{n} / v)^{1/2} dt = \int_{a}^{b} (\int_{X} (g*\dot{\gamma}_{t})^{2} (g*\ddot{\gamma}_{t})^{n} / v)^{1/2} dt = L(\Gamma).$$

(9.2) <u>Proof of Theorem</u> C: By Theorem B, E is isometric to the Riemannian symmetric space G/K without compact factors. In particular, E is a simply connected Riemannian manifold with nonpositive sectional curvature. Since the compact group H acts isometrically on E, it always has a fixed point in E.

<u>Remark</u> (9.3). If X admits no nonzero holomorphic vector fields, then Theorems A and C assert the following:

Einstein Kähler metrics on X are, if any, unique up to constant multiple. Moreover, they are invariant under the action of Aut(X).

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