

**Trace Norm Estimates for
Products of Integral Operators
and Diffusion Semigroups**

**Michael Demuth
Peter Stollmann
Günter Stolz
Jan van Casteren**

Michael Demuth
Max-Planck-Arbeitsgruppe
Fachbereich Mathematik
Universität Potsdam
14415 Potsdam, Germany

Max-Planck-Institut für Mathematik
Gottfried-Claren-Straße 26
53225 Bonn
Germany

Peter Stollmann and Günter Stolz
Fachbereich Mathematik
Universität Frankfurt
60054 Frankfurt, Germany

Jan van Casteren
Dept. of Math. and Comp. Science
University of Antwerp
Universiteitsplein 1
2610 Antwerp-Wilrijk, Belgium

**Trace Norm Estimates for
Products of Integral Operators
and Diffusion Semigroups**

**Michael Demuth
Peter Stollmann
Günter Stolz
Jan van Casteren**

Michael Demuth
Max-Planck-Arbeitsgruppe
Fachbereich Mathematik
Universität Potsdam
14415 Potsdam, Germany

Max-Planck-Institut für Mathematik
Gottfried-Claren-Straße 26
53225 Bonn
Germany

Peter Stollmann and Günter Stolz
Fachbereich Mathematik
Universität Frankfurt
60054 Frankfurt, Germany

Jan van Casteren
Dept. of Math. and Comp. Science
University of Antwerp
Universiteitsplein 1
2610 Antwerp-Wilrijk, Belgium

TRACE NORM ESTIMATES FOR PRODUCTS OF INTEGRAL OPERATORS AND DIFFUSION SEMIGROUPS

Michael Demuth
Peter Stollmann
Günter Stolz
Jan van Casteren

Abstract

We give trace norm estimates for products of integral operators and for diffusion semigroups. These are applied to differences of heat semigroups. A natural example of an integral operator with finite trace which is not trace class is given.

INTRODUCTION

We prove two trace class criteria. The first, Theorem 1, provides an estimate for the trace norm of the product of two integral operators. The second, Theorem 3, concerns differences of diffusion semigroups. Both results are inspired by the same circle of problems, namely the search for trace estimates for differences of heat semigroups, which in turn are a powerful tool in the investigation of spectral properties of the associated Hamiltonians. The according applications are indicated in Section 3. Let us now give a little more details concerning the following sections.

Section 1 is devoted to a proof of Theorem 1 which states that

$$\|AB\|_{tr} \leq \int \|A[\cdot, x]\|_2 \|B[x, \cdot]\|_2 dm(x),$$

if A, B are operators with kernels $A[\cdot, \cdot], B[\cdot, \cdot]$ and the L_2 -norms in the integral are assumed to exist. As one immediately notices, this includes the well-known case that A, B are Hilbert–Schmidt, but it is much more general: The kernels A, B do not even have to define bounded operators in L_2 . We then relate the above estimate to Corollary 2, which is the key to the results of the second section. There we treat differences of ultracontractive diffusion semigroups. The advantage of Theorem 3 in comparison with the results of [9] is the fact that we do not have to assume the validity of a Feynman–Kac formula or even the existence of a stochastic process. This enables the easy application to Neumann boundary problems given in Corollary 5. We end the third section by giving an example which clarifies some aspects of the trace norm estimates for semigroup differences: We show that an additional Dirichlet boundary condition on a set of finite capacity can lead to a semigroup difference which is not trace class, but is a Hilbert–Schmidt operator with finite trace. This shows that a conjecture in [9] is wrong. Moreover it appears to be the first “natural” example of an operator with positive continuous kernel and finite trace which, nevertheless, is not trace class.

1. INTEGRAL OPERATORS

We assume throughout that (X, \mathfrak{A}, m) is a σ -finite measure space and we are concerned with trace class operators on $L_2 = L_2(X, \mathfrak{A}, m)$ which we denote by $\mathfrak{B}_1 = \mathfrak{B}_1(L_2)$. We use $\|\cdot\|_{tr}$ for the trace norm on \mathfrak{B}_1 and write $(\mathfrak{B}_2, (\cdot|\cdot)_{HS})$ for the Hilbert Schmidt operators, where

$$(A|B)_{HS} = \text{trace}(B^*A).$$

A measurable function $A[\cdot, \cdot] : X \times X \rightarrow \mathbb{C}$ such that

$$(Af|g) = \int \int A[x, y]f(y)g(x)dm(x)dm(y),$$

or, equivalently,

$$Af(\cdot) = \int A[\cdot, y]f(y)dm(y)$$

is said to be a kernel for the operator A .

THEOREM 1 *Let $A, B : X \times X \rightarrow \mathbb{C}$ be measurable such that $A[\cdot, x], B[x, \cdot] \in L_2$ for a.e. $x \in X$ and*

$$\int \|A[\cdot, z]\|_2 \|B[z, \cdot]\|_2 dm(z) < \infty. \quad (1)$$

Then there is a trace class operator $AB : L_2 \rightarrow L_2$ with kernel

$$AB[x, y] = \int A[x, z]B[z, y]dm(z)$$

such that

$$\|AB\|_{tr} \leq \int \|A[\cdot, z]\|_2 \|B[z, \cdot]\|_2 dm(z) \quad (2)$$

PROOF. Set $h(z) := \|A[\cdot, z]\|_2, g(z) := \|B[z, \cdot]\|_2$. With the convention $g^{-1}(x) := 0$ where $g(x) = 0$, we write $M_{g^{-1}}$ for the corresponding multiplication operator. It follows that

$$|Bf(x)| = \left| \int B[x, y]f(y)dm(y) \right| \leq g(x)\|f\|_2,$$

so that $M_{g^{-1}}B : L_2 \rightarrow L_\infty$ is defined and has norm less than 1. Similarly,

$$M_{h^{-1}}A^* : L_2 \rightarrow L_\infty, \|M_{h^{-1}}A^*\| \leq 1.$$

By assumption $hg \in L_1$, so that

$$M_{(hg)^{1/2}}M_{g^{-1}}B : L_2 \rightarrow L_2$$

is bounded. As a composition of a so-called Carleman operator with an L_2 -multiplication it is even Hilbert-Schmidt with Hilbert-Schmidt norm less than

$$\|(hg)^{1/2}\|_2 \|M_{g^{-1}}B : L_2 \rightarrow L_\infty\| \leq \|hg\|_1^{1/2}.$$

The same argument gives that

$$\|AM_{h^{-1}}M_{(hg)^{1/2}}\|_{HS} = \|M_{(hg)^{1/2}}M_{h^{-1}}A^*\|_{HS} \leq \|hg\|_1^{1/2}.$$

Hence the composition

$$AB := AM_{h^{-1}}M_{(hg)^{1/2}}M_{(hg)^{1/2}}M_{g^{-1}}B$$

is a trace class operator which satisfies (2). It requires only a straightforward calculation to show that the kernel has the asserted form. \square

The original proof of the following Corollary given in [10] is quite similar to the above proof of the Theorem (write $AB = AM_{\Phi^{1/2}}M_{\Phi^{1/2}}M_{\Phi^{-1}}B$).

COROLLARY 2 *Let $A \in \mathfrak{B}(L_1, L_2)$, $B \in \mathfrak{B}(L_2, L_1)$ and assume that there exists a Φ in L_1 such that $|Bf| \leq \Phi$ for every f in the unit ball of L_2 . Then*

$$\|AB\|_{tr} \leq \|A\| \cdot \|\Phi\|_1.$$

This lemma will be the key to the results in the following section. Apart from its applications in Section 3, it proved to be a very useful tool in the spectral theoretic investigations of [10].

2. DIFFUSION SEMIGROUPS

We call a semigroup $U = (U(t); t \geq 0)$ a *diffusion semigroup* if the following conditions are satisfied

- $U(t) \in \mathfrak{B}(L_2)$ is selfadjoint for all $t \geq 0$.
- $U(t)$ induces a bounded operator on L_p for all $t \geq 0, p \in [1, \infty)$
- U is positivity preserving, i.e. $U(t)f \geq 0$ for $f \geq 0, t \geq 0$.

If furthermore,

- $U(t)$ induces a bounded operator from L_1 to L_∞ for all $t > 0$

we speak of an *ultracontractive diffusion semigroup*. To simplify notation, we denote by $\|A\|_{p,q}$ the norm of an operator from L_p to L_q and we use

$$L_q := \{f; |f|^q \in L_1\}, \|f\|_q := \||f|^q\|_1$$

for $0 < q < 1$. There is a natural order for positivity preserving semigroups which comes from the order of functions, namely

$$V \leq U : \iff \forall t \geq 0, f \geq 0 : V(t)f \leq U(t)f.$$

The main result of this section deals with differences of semigroups which obey this order relation.

THEOREM 3 *Assume that U, V are ultracontractive diffusion semigroups satisfying $V \leq U$ and set $D(t) := U(t) - V(t)$ for $t \geq 0$. If $D(t)1 \in L_{1/2}$ for some $t > 0$ then*

$$\|U(2t) - V(2t)\|_{tr} \leq \|D(t)1\|_{1/2} \|D(t)\|_{1,\infty}^{1/2} (\|U(t)\|_{1,2} + \|V(t)\|_{1,2}).$$

We single out one step in the proof of Theorem 3 which can be thought of as a Cauchy-Schwarz inequality for positivity preserving operators. For integral operators it can easily be deduced from the usual Cauchy-Schwarz inequality. In the proof below we make essential use of the existence of a *lifting* for σ -finite measure spaces (see [6, 4] for background information).

LEMMA 4 *Assume that $A : L_2 \rightarrow L_2$ is positivity preserving and induces a bounded operator from L_p to L_q for all $p, q \in [1, \infty], p \leq q$. Then, for $f \in L_2$:*

$$|Af| \leq (A1)^{1/2} \cdot (A(|f|^2))^{1/2}.$$

PROOF. Denote by \mathfrak{L}_∞ the essentially bounded measurable functions (not equivalence classes!). Since m is σ -finite there exists a lifting Λ , by which we understand a linear multiplicative (hence order preserving) mapping

$$\Lambda : L_\infty \longrightarrow \mathfrak{L}_\infty,$$

such that Λf is a function in the equivalence class f . For fixed $x \in X$ set

$$q_x : L_2 \times L_2 \longrightarrow \mathbb{C}, q_x(f, g) := \Lambda(A(f\bar{g}))(x).$$

As Λ is linear and positive, q_x is a positive sesquilinear form. The Cauchy-Schwarz inequality implies

$$|\Lambda(A(f\bar{g}))(x)| \leq (\Lambda(A|f|^2)(x))^{1/2} \cdot (\Lambda(A|g|^2)(x))^{1/2}$$

for all $x \in X$. Since Λf is a representative of f , we may take $g \in L_2, 0 \leq g \leq 1$ in the last inequality and obtain

$$|A(fg)| \leq (A|f|^2)^{1/2} \cdot (A1)^{1/2},$$

since $Ag \leq A1$. Approximating the constant function 1 from below by a sequence g_n such that $0 \leq g_n \leq 1, g_n \in L_2$ and taking the limit $n \rightarrow \infty$ gives the desired inequality. \square

PROOF of Theorem 3. First note that, by the semigroup property of U and V ,

$$D(2t) = U(t)D(t) + D(t)V(t).$$

By Lemma 4, for $\|f\|_2 \leq 1$,

$$\begin{aligned} |D(t)f(x)| &\leq (D(t)1(x))^{1/2} \cdot (D(t)|f|^2(x))^{1/2} \\ &\leq (D(t)1(x))^{1/2} \|D(t)\|_{1,\infty}^{1/2} =: \Phi(x). \end{aligned}$$

Hence we can apply Corollary 2 and obtain

$$\begin{aligned} \|U(t)D(t)\|_{tr} &\leq \|\Phi\|_1 \cdot \|U(t)\|_{1,2} \\ &= \|D(t)1\|_{1/2} \|D(t)\|_{1,\infty}^{1/2} \|U(t)\|_{1,2}. \end{aligned}$$

By the same arguments

$$\begin{aligned} \|D(t)V(t)\|_{tr} &= \|V(t)D(t)\|_{tr} \\ &\leq \|D(t)1\|_{1/2} \|D(t)\|_{1,\infty}^{1/2} \|V(t)\|_{1,2}, \end{aligned}$$

so that the asserted estimate follows. \square

3. APPLICATIONS AND EXAMPLES

In this section we want to illustrate the above theorems by some applications. Although we are interested in more general Hamiltonians (see [2]) we restrict ourselves to the Laplacian on \mathbb{R}^d in order to keep preliminary definitions and technicalities at a minimum. We denote the heat semigroup by $U(t) := e^{1/2\Delta t}$ and write $\Delta_\Sigma^D, \Delta_\Sigma^N$ for the Dirichlet, respectively Neumann Laplacian on an open set $\Sigma \subset \mathbb{R}^d$. The latter are selfadjoint operators on $L_2(\Sigma)$, and we extend the semigroups they generate in the obvious way to all of $L_2(\mathbb{R}^d) = L_2(\Sigma) \oplus L_2(\Sigma^c)$ by setting $U_\Sigma^D := e^{1/2\Delta_\Sigma^D t} \oplus 0$. With the analogous notation for the Neumann operator we note in passing that

$$U_\Sigma^D \leq U, U_\Sigma^D \leq U_\Sigma^N,$$

while $U_\Sigma^N \not\leq U$ apart from trivial cases. While U, U_Σ^D are always ultracontractive (see [1], Section 2.1, especially Example 2.1.8), this is not the case for U_Σ^N (the Neumann Laplacian need not even have compact resolvent). By \mathbb{P}^x we denote the Wiener measure for particles starting in x and get:

COROLLARY 5 *Let $\phi_{\Sigma,t} := \mathbb{P}^x\{X_s \in \Sigma^c \text{ for some } s \leq t\}$ for any open $\Sigma \subset \mathbb{R}^d$.*

- (1) $\|U(2t) - U_\Sigma^D(2t)\|_{tr} \leq c(t) \int \phi_{\Sigma,t}(x)^{1/2} dx.$
- (2) *If U_Σ^N is ultracontractive, then*

$$\|U(2t) - U_\Sigma^N(2t)\|_{tr} \leq c(t) \int \phi_{\Sigma,t}(x)^{1/2} dx.$$

PROOF. By the Feynman–Kac formula ([3]),

$$U_\Sigma^D(t)1(x) = \mathbb{P}^x\{X_s \in \Sigma \text{ for all } s \leq t\} \leq \chi_\Sigma.$$

Consequently,

$$\begin{aligned} (U(t) - U_\Sigma^D(t))1(x) &= 1 - \mathbb{P}^x\{\dots\} \\ &= \mathbb{P}^x\{X_s \in \Sigma^c \text{ for some } s \leq t\}. \end{aligned}$$

Theorem 3 implies

$$\|U(2t) - U_\Sigma^D(2t)\|_{tr} \leq c(t) \|\phi_{\Sigma,t}\|_{1/2},$$

proving (1). If, furthermore, U_Σ^N is ultracontractive, we can apply Theorem 3 to the difference $U_\Sigma^N - U_\Sigma^D$, since $U_\Sigma^N 1 = \chi_\Sigma$ and therefore

$$(U_\Sigma^N(t) - U_\Sigma^D(t))1 = \phi_{\Sigma,t} \chi_\Sigma.$$

This yields (2). □

We would like to mention that the Neumann heat semigroup is ultracontractive if Σ has the extension property (see [1], Theorem 2.4.4, p. 77). Another way to prove part (2) of the above Corollary would be to apply the analysis of [9] to the Dirichlet form generated by the Neumann Laplacian. In order to do so, one faces technical problems related with the

existence of an associated process.

In the situation of Corollary 5(1) it would be desirable to weaken the assumption on $\phi_{\Sigma,t}$ to the requirement $\phi_{\Sigma,t} \in L_1$, since the latter is fulfilled for all sets Σ satisfying $\text{cap}(\Sigma^c) < \infty$ (see the proof of the following lemma), which in turn is a quite natural condition. In [9] the corresponding statement was formulated as a conjecture. The following lemma and the subsequent example show, however, that $\text{cap}(\Sigma^c) < \infty$ does not imply $\|U(t) - U_{\Sigma}^D(t)\|_{t^*} < \infty$. There is one more reason why we find this example quite interesting: From the results of [9] it is clear that the semigroup difference in question is a Hilbert–Schmidt operator with positive continuous kernel. Moreover, it is easy to see that its trace is finite. Thus, according to a remark of Simon, [7], Remark 2, p. 77 one would expect it to be trace class, which is not the case. To introduce our example we have to recall the definition of the Birman–Solomjak space

$$l_1(L_2) := \left\{ f : \mathbb{R}^d \rightarrow \mathbb{R}; \sum_{\alpha \in \mathbb{Z}^d} \left(\int_{C_\alpha} |f(x)|^2 dx \right)^{1/2} < \infty \right\},$$

where C_α denotes the unit cube centered at α ; see [7], p. 55.

LEMMA 6 *Assume that $\Gamma := \Sigma^c$ satisfies $\text{cap}(\Gamma) < \infty$ but $\chi_\Gamma \notin l_1(L_2)$. Then $\phi_{\Sigma,t} \in L_1$ but $U(t) - U_{\Sigma}^D(t) \notin \mathfrak{B}_1$ for any $t > 0$.*

PROOF. For the potential theoretic notions used in this proof we refer the reader to [5], Chapter 3. Recall that

$$\text{cap}(\Gamma) = \min \left\{ \int |\nabla f|^2 + |f|^2 dx; f \in W_0^{1,2}, \tilde{f} \geq \chi_\Gamma \right\},$$

where \tilde{f} denotes the quasi-continuous representative of f . The unique minimizing element e_Γ is called the 1-equilibrium potential of Γ and can be represented by

$$e_\Gamma(x) = \int G(x, y) d\nu_\Gamma(y),$$

where ν_Γ is a measure supported on Γ with total mass equal to the capacity of Γ , and $G(x, y)$ is the kernel of $(-\Delta + 1)^{-1}$. Since $\sup_y \int G(x, y) dx = \|(-\Delta + 1)^{-1} 1\|_\infty \leq 1$,

$$\begin{aligned} \|e_\Gamma\|_1 &= \int \int G(x, y) d\nu_\Gamma(y) dx \\ &= \int \left(\int G(x, y) dx \right) d\nu_\Gamma \\ &\leq \text{cap}(\Gamma). \end{aligned}$$

Denote $\tau(w) := \inf\{s > 0; X_s(w) \in \Sigma\}$, the first hitting time of Γ . Then

$$\phi_{\Sigma,t}(x) = \mathbb{E}^x \{\tau \leq t\} \leq e^t \mathbb{E}^x \{e^{-\tau}\} = e^t e_\Gamma(x),$$

where we used [5], Lemma 4.3.1 in the last step. This proves the first assertion. If $U(t) - U_{\Sigma}^D(t) \in \mathfrak{B}_1$ for some $t > 0$, it follows that $\chi_\Gamma(U(t) - U_{\Sigma}^D(t)) \in \mathfrak{B}_1$. Since $\chi_\Gamma(U(t) - U_{\Sigma}^D(t)) = \chi_\Gamma U(t)$, we may apply [7], Proposition 4.7, to deduce $\chi_\Gamma \in l_1(L_2)$. \square

EXAMPLE 7 If $d \geq 5$ and $\Gamma := \bigcup_n B_n$, where B_n is a ball of radius r_n centered at $(n, 0, \dots, 0)$ with $r_n \leq 1/2$ we have

$$\text{cap}(\Gamma) \leq c \sum_n r_n^{d-2}, \|\chi_\Gamma\|_{l_1(L_2)} = c' \sum_n r_n^{d/2}.$$

For $r_n = 1/2 \cdot n^{-2/d}$ it follows that $\text{cap}(\Gamma) < \infty$ and $\chi_\Gamma \notin l_1(L_2)$. Consequently, $U(t) - U_\Sigma^D(t)$ is Hilbert–Schmidt with finite trace but not trace class.

References

- [1] E.B. Davies: Heat Kernels and Spectral Theory. Cambridge, Cambridge University Press, 1989
- [2] M. Demuth: On topics in spectral and stochastic analysis for Schrödinger operators. Proceedings “Recent Developments in Quantum Mechanics”, A. Boutet de Monvel et al. (eds.), Kluwer, 1991
- [3] M. Demuth and J. van Casteren: On spectral theory of selfadjoint Feller generators. Rev. Math. Phys. 1, 325–414 (1989)
- [4] K. Floret: Maß- und Integrationstheorie. Stuttgart, B.G. Teubner, 1981
- [5] M. Fukushima: Dirichlet Forms and Markov Processes. Amsterdam, North Holland, 1980
- [6] A. and C. Ionescu–Tulcea: Topics in the theory of liftings. Berlin, Springer–Verlag, 1969
- [7] B. Simon: Trace ideals and their applications. London Mathematical Society Lecture Note Series 35, Cambridge, Cambridge University Press, 1979
- [8] B. Simon: Schrödinger semigroups. Bull. Amer. Math. Soc. 7, 447–526 (1982)
- [9] P. Stollmann: Scattering by obstacles of finite capacity. J. Funct. Anal., to appear
- [10] P. Stollmann and G. Stolz: Singular spectrum for multidimensional Schrödinger operators with potential barriers. Preprint 1993

Author’s addresses:

Michael Demuth
Max–Planck–Arbeitsgruppe
Fachbereich Mathematik
Universität Potsdam
Postfach 601553
14415 Potsdam, Germany

Peter Stollmann and Günter Stolz
Fachbereich Mathematik
Universität Frankfurt
60054 Frankfurt am Main, Germany

Jan A. van Casteren
Dept. of Math. and Comp. Science
University of Antwerp
Universiteitsplein 1
2610 Antwerp-Wilrijk, Belgium